

Project Number: 101072443

2nd NETWORK TRAINING SCHOOL

Deliverable 4.3 Dissemination level: PU - Public

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Document Information

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SEDIMARE - 101072443 - D4.3: 2nd NETWORK TRAINING SCHOOL

Version	Publication Date	Author(s)	Change
1.0	15.11.2024	Athanassios Dimas	Initial version

Date and Signature of Author(s):

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1. Program

The SEDIMARE "2nd Training School" was organized by the University of Twente and Deltares at Twente and Delft, respectively, in the Netherlands, on 5-8 November 2024. Local organizers were Prof. P. Roos and Prof. J. van der Werf. The event was attended in person by all DCs.

The SEDIMARE meeting consisted of two parts:

- 1. Training for the DCs at the University of Twente, Enschede on 5-6 November. The DCs followed a presentation skills training. They worked on the presentations to be given at Part 2 of the meeting. Draft presentations had to be ready before the start of the training.
- 2. Experimental & modeling workshop for all SEDIMARE members at Deltares, Delft on 7-8 November.

Tuesday 5 November (DCs only) - University of Twente, Horsttoren, room Z-203

- 08.45-09.00 walk-in
- 09.00 09.15 welcome by Kathelijne Wijnberg, UT professor Coastal Systems a& Nature-Based Engineering
- 09.15 12.30 training presentation skills
- 12.30 14.00 lunch
- 14.00 17.00 training presentation skills

Wednesday 6 November (DCs only) - University of Twente, Horsttoren, room Z-203

- 09.00 12.30 training presentation skills
- 12.30 13.00 lunch
- ~14.00 train to Delft

Thursday 7 November – Deltares, room Ganges Delta

Training School & Workshop "Experimental and Practical Modelling of Sediment Transport and Coastal Morphology"

- 09.00 09.30 walk-in
- 09.30 09.45 welcome by Dirk-Jan Walstra (managing director Deltares) & Jebbe van der Werf
- 09.45 10.00 Athanassios Dimas (PI) introduction of SEDIMARE
- 10.00 10.40 Keynote 1: Bas van Maren (Deltares, TUDelft) "Measurements and modeling of the erodibility of sand-mud beds" (30 min + 10 min discussion)
- 10.40-11.00 coffee/tea break

Morphodynamic modeling oriented DCs (10 min presentation, 10 min discussion)

- 11.00 11.20 Buckle Subbiah Elavazhagan (#5) Morphodynamic analysis of the upper confined and unconfined beach profiles during episodic events
- 11.20 11.40 Jowi Miranda (#7) Process-based modeling of sand mud morphodynamics
- 11.40 12.00 Nasim Soori (#4) Mixing and transport in the coastal area
- 12.00 12.20 Nishchay Tiwari (#12) Multi-models approach to scour in dynamics areas
- 12.20-13.30 lunch
- 13.30 14.10 Keynote 2: Sandra Soares-Frazão, UCL "Breaching of earthen embankments: from small-scale laboratory experiments to field measurements using photogrammetry" (30 min + 10 min discussion)
- Field/lab oriented DCs (10 min presentation, 10 min discussion)
- 14.10 14.30 Muhammed Said Parlak (#2) Nearshore wave processes by remote sensing
- 14.30-14.50 Van Thi To Nguyen (#3) Erosion and transport processes of sand-silt mixtures
- 14.50 15.10 coffee/tea break
- 15.10 15.30 Siyuan Wang (#8) Morphodynamics of breach growth and bank erosion using laboratory experiments
- 15.30 15.50 Eloah Rosas (#9) Characterization of stratification and near-bed dense layers in high-density sediment-laden flows
- 15.50 16.10 Saeed Osouli (#10) Overtopping breakwater for energy conversion
- 16.10 17.00 discussion of DC secondments (lead by Athanassios Dimas)
- *19.00 "Meeting" dinner in the evening.*

Friday 8 November – Deltares, room Ganges Delta

Training School & Workshop "Experimental and Practical Modelling of Sediment Transport and Coastal Morphology"

Detailed numerical modeling DCs (10 min presentation, 10 min discussion)

- 09.30 9.50 Ioannis Gerasimos Tsipas (#1) Large-eddy simulations of turbulent oscillatory flow and sediment transport induced by waves
- 09.50 10.10 Quan Nguyen (#11) Morphodynamic swash zone modelling
- 10.10 10.30 Evangelos Petridis (#6) Mathematical modeling and numerical simulations of water-saturated granular materials
- 10.30 11.00 Keynote 3: Mark Klein Breteler Deltares experimental facilities & measuring techniques (20 min + 10 min discussion)
- 11.00 11.15 coffee/tea break
- 11.15 12.30 guided tour of Deltares experimental facilities
- 12.30-13.30 lunch

2. Presentations

The theme of the 2nd Network Training School was "Experimental and Practical Modelling of Sediment Transport and Coastal Morphology". Therefore, the scope of the presentations by scientists was on the fundamentals and the applications of experimental and field methods, as well as on practical modelling to be used by the DCs in their research.

The corresponding presentations, as well as the presentations of all the DCs, are shown in the next pages.

"Measurements and modeling of the erodibility of sand-mud beds" (Invited Speaker Bas van Maren, Deltares & TUDelft)



Deltares

Measurements and modeling of the erodibility of (sand)-mud beds

Recent insights into the role of density

Bas van Maren

(with most of the work done by others, especially Leo van Rijn, Roy van Weerdenburg, Marcio Boechat, Ana Colina, Irene Colosimo, Thijs van Kessel, Floris van Rees)

Wadden Sea fringing flats (near Koehool)

What is the mud content? What is the dry density?

A Teres

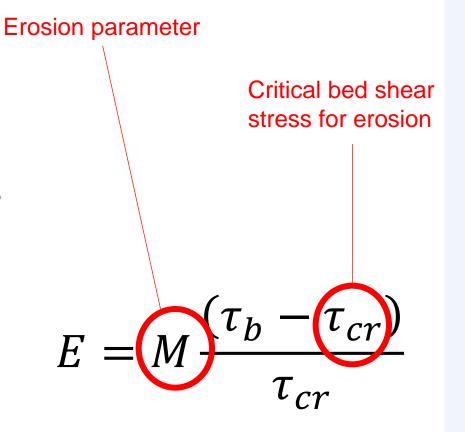
'Mud' Density ~500 kg/m³ Mud content ~50-60%

PhD work Irene Colosimo



Recent insights

- Role of density in erosion of sand-mud mixtures
 - MUSA experiments
 - Delft3D improvements
- Strength development during drying
 - Field experiments
 - Delft3D improvements



Role of density in erosion of sand-mud mixtures

- MUSA project (Improve understanding of erosion & sedimentation of sand-mud mixtures and implement it in engineering tools and numerical software)
- 2020 2023 (MUSA2 2024 2028)
- Consortium of Leo v. Rijn, WaterProof, Boskalis, DEME, Jan de Nul, RHDHV, Arcadis, DHI, HRW, Deltares (Svasek, Rijkswaterstaat, Waterbouwkundig Laboratorium Borgerhout)



Literature survey

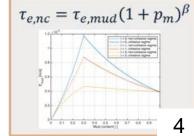


Lab experiments - Erosion (flumes)

- Properties



Field surveys - boat + frame - Flocculation camera measurements

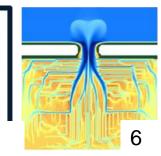


Analysis + Formulations



Toolbox - Leo van Rijn tools

- CoDeS

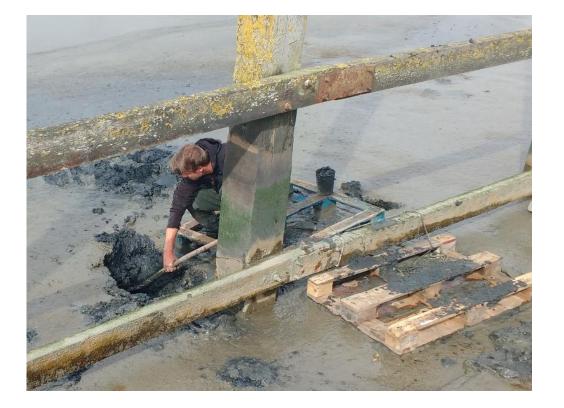


Upgrade numerical models: Delft3D + Mike

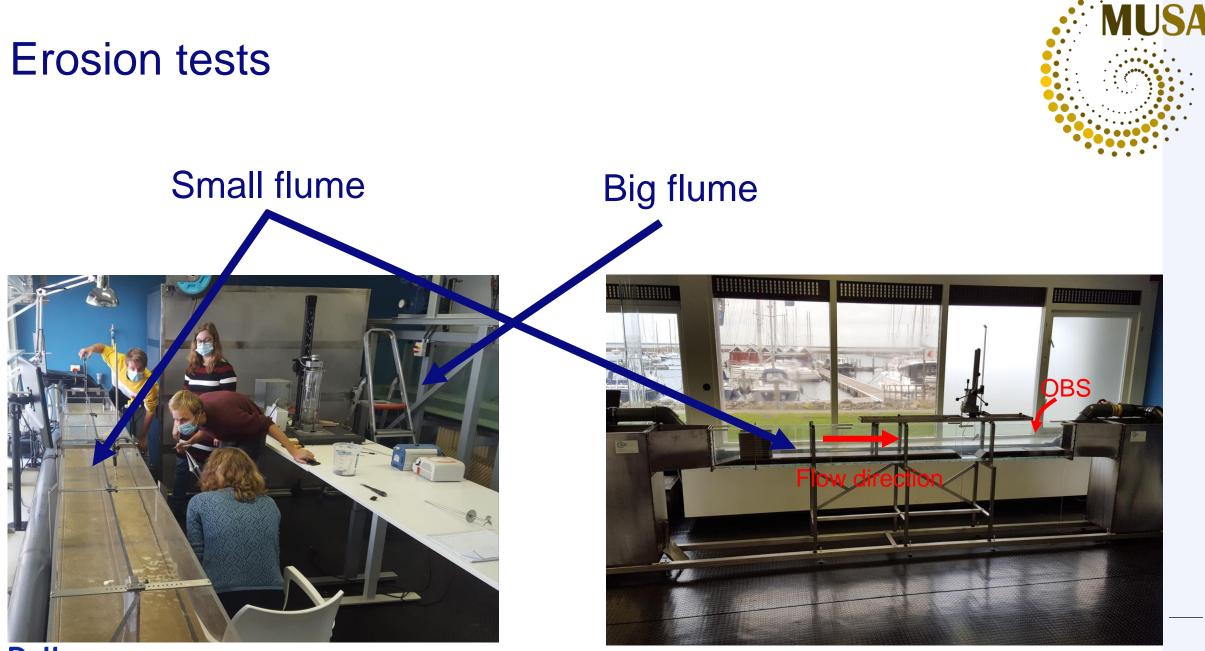
Deltares

Sample collection

- Bulk samples ٠
- Undisturbed samples ٠







Erosion tests



Deltares

MUSA

Erosion tests

Sediment beds:

- Placed beds (undisturbed samples)
- Remoulded beds
- Diluted, remoulded beds

verification

Observation method

- Visual determination of Critical Bed Shear Stress τ_{cr} (all samples)
- SSC conversion to erosion rate to determine τ_{cr} & erosion parameter (some samples)

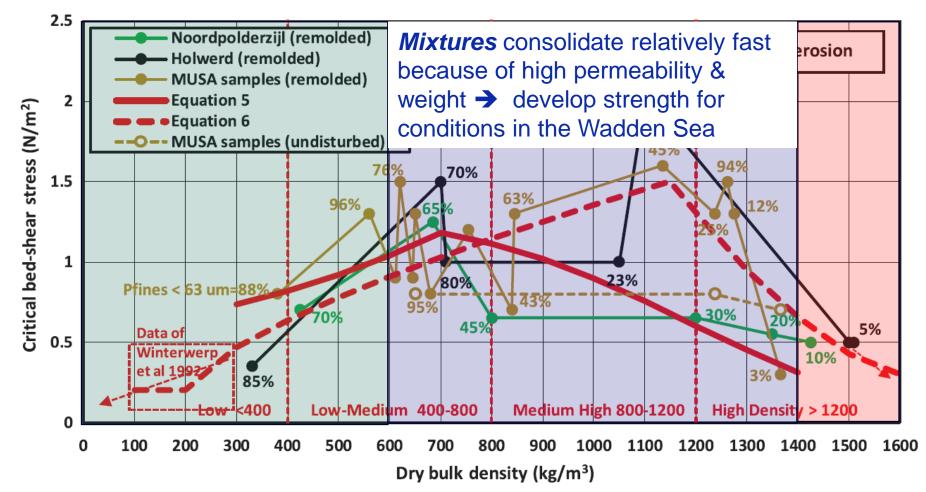


Critical Bed-Shear Stress of Mud–Sand Mixtures

L. C. van Rijn, Dr.Eng.¹; M. Boechat Albernaz, Ph.D.²; L. Perk³; A. Colina Alonso, Dr.Eng.⁴; R. J. A. van Weerdenburg⁵; and D. S. van Maren, Ph.D.⁶

Muddy sediment consolidates slowly → remains soft under Wadden Sea conditions

Sand erodes relatively easy



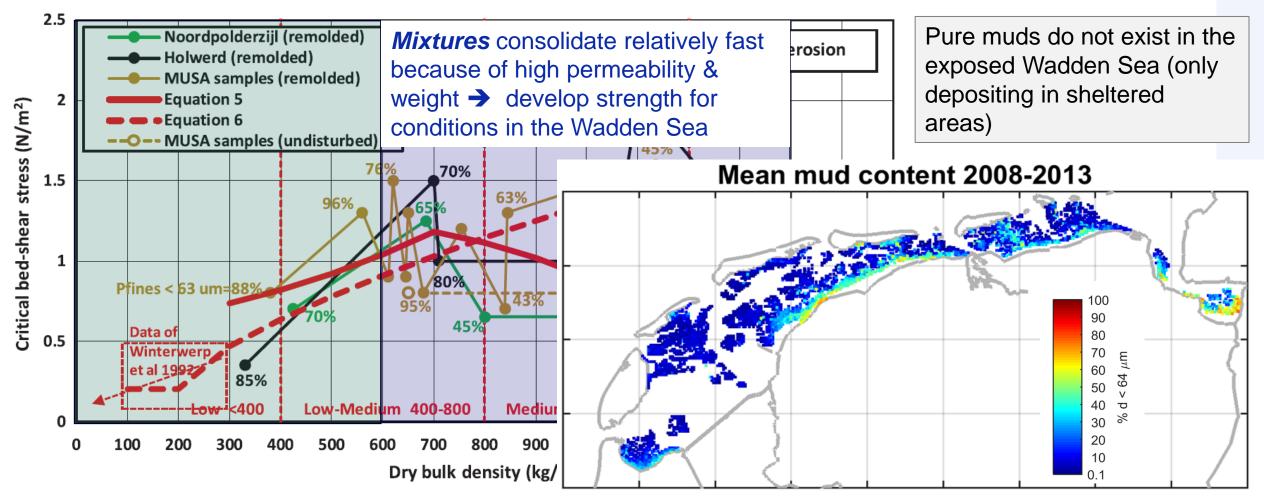
- *τ*_{cr} is highest for intermediate density
- This corresponds to mixtures with a mud content of 30-70%
- High density: sandy mixtures with a low τ_{cr} (Shields curve erosion)
- Low density: muddy mixtures with a low τ_{cr}



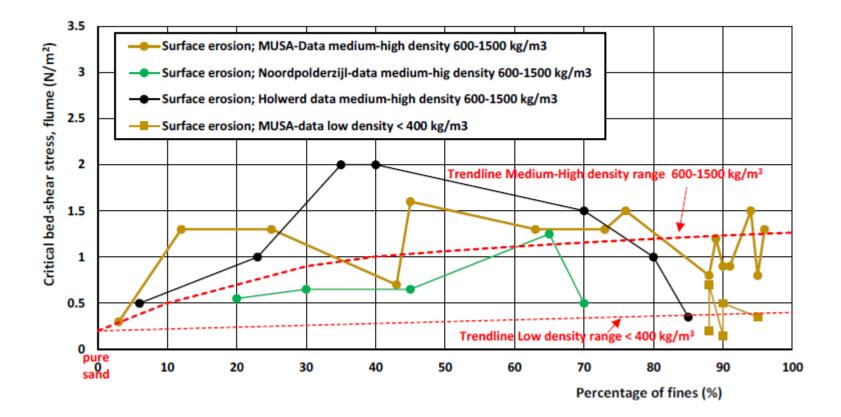
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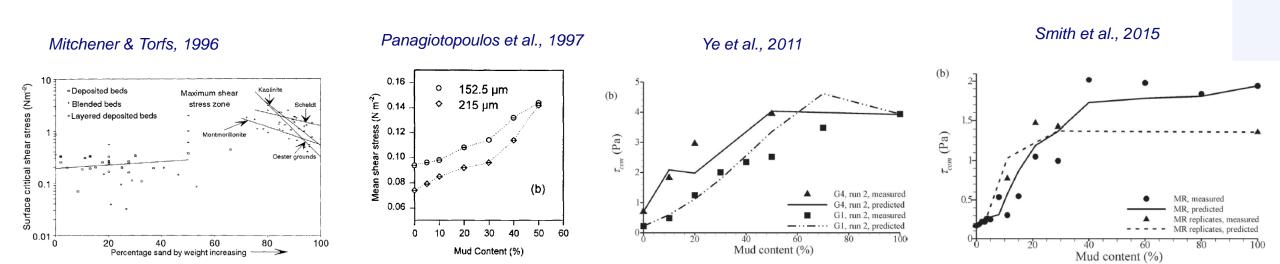
• Why not relating directly to mud content? Correlation with density is stronger.



MUSA

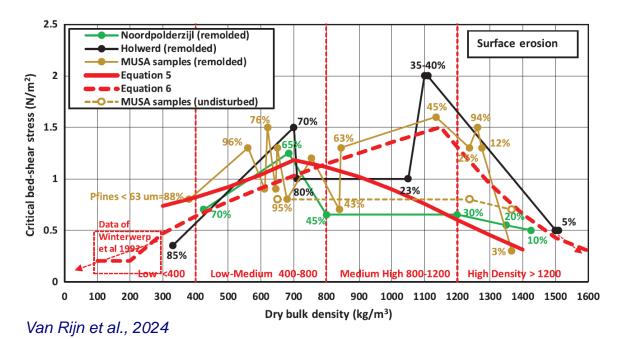
So what is new?

- Most earlier work focussed on *maximum* densities, not *actual* densities, of the sand-mud mixtures
- But in dynamic systems with regular resuspension and deposition as fluid mud, muddy mixtures do not have time to consolidate



Implementation in Delft3D

- Equation 6 implemented in Delft3D
- Density related to mud content via empirical site-specific relation $\rho_{dry} = 500 p_{fines} + 1600 p_{sand}$.



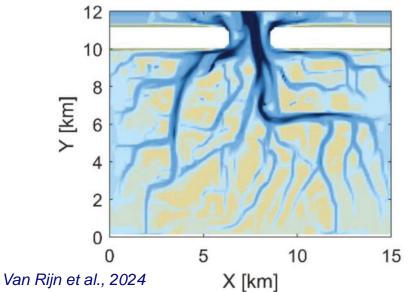
$$\tau_{cr,\text{fines}} = \tau_{cr,\text{min1}} + \left[(\rho_{\text{dry}} - \rho_{dry,\text{min}}) / (\rho_{\text{dry},*} - \rho_{dry,\text{min}}) \right]^{\alpha 1} (\tau_{cr,\text{max}} - \tau_{cr,\text{min1}}) \quad \text{for } \rho_{\text{dry}} < \rho_{\text{dry},*}$$
(6a)

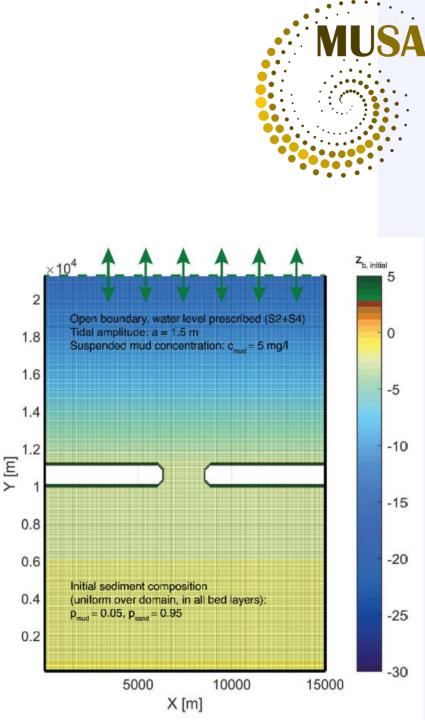
$$\tau_{cr,\text{fines}} = \tau_{cr,\text{min2}} + [(\rho_{dry,\text{max}} - \rho_{dry})/(\rho_{dry,\text{max}} - \rho_{dry,*})]^{\alpha 2} (\tau_{cr,\text{max}} - \tau_{cr,\text{min2}}) \quad \text{for } \rho_{dry} > \rho_{dry,*}$$
(6b)



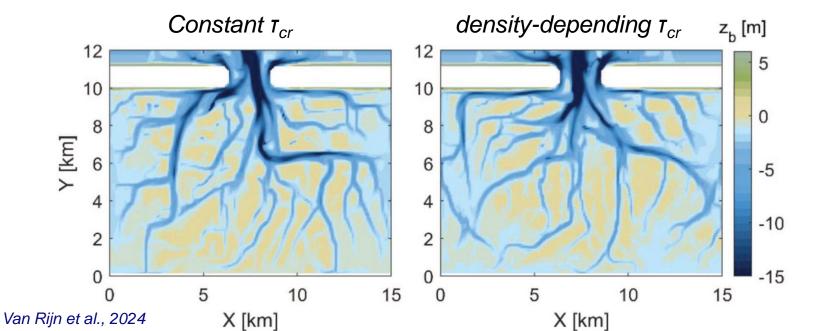
- Schematized tidal basin with uniform initial bathymetry consisting of 95% sand and 5% mud
- Mud and sand prescribed at open boundary conditions
- Model is run for tidal conditions with mild waves for 45 years to generate a typical tidal basin morphology
- Followed by 5 years with more energetic wave conditions

Constant τ_{cr}



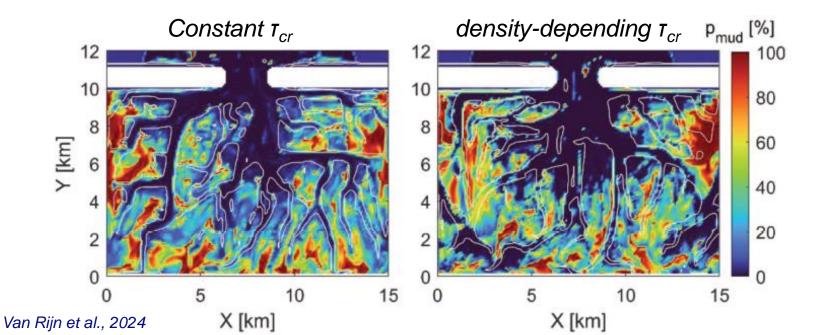


- Calm conditions (45 years):
 - marginal effect of a density-depending τ_{cr}





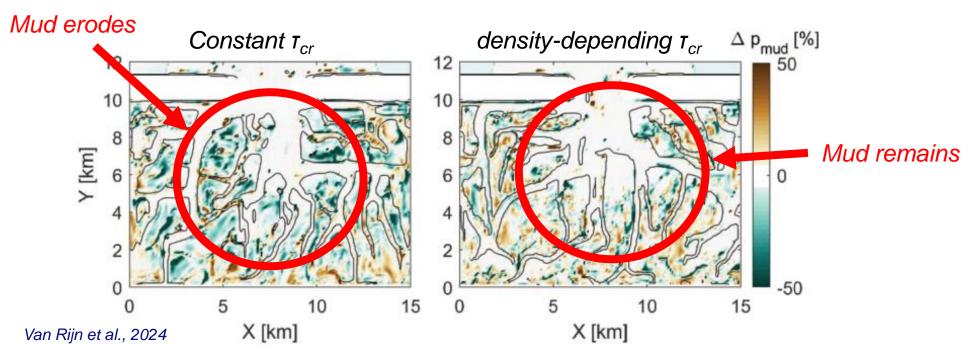
- Calm conditions (45 years):
 - marginal effect of a density-depending τ_{cr}
 - Mud deposits on the higher flats



- Calm conditions (45 years):
 - marginal effect of a density-depending τ_{cr}
 - Mud deposits on the higher flats
- Storm conditions (5 years)

A simple parameterization for a complex process:

- Tides bring sediment to the flats
- Muddy sediments consolidate and gain strength to withstand erosion
- But then... without a complex consolidation model
- And including sand-mud
- For constant τ_{cr} , a large amount of mud deposited during calm conditions is eroded
- For density-depending τ_{cr} , the mud largely remains on the flats



SE SW Colisomo et al., 2023 NW [s/m] d10 Erosion-Deposition alternation BLC [cm] 2 [ed] 1-Erosion = [0.3-1] cn = [0.5-2] Pa Fast Recover Erosion = 1.8 cm r_{peak} = [0.2-2] Pa E Pa Erosion = 1.9 cm 0 Tneek = 14.4 Pa m Erosion = 3.5 cm = [1-3] Pa .C [cm] Pa Frosion = 3.4 crB Erosion = 1.3 cm = [2.5-4] Pa No Erosion No Erosion = 10 Pa BLC [cm] Erosion = 1 cm 12/17 12/24 12/31 01/07 01/14 01/21 01/28 02/04 02/11 time (mm/dd)

KL

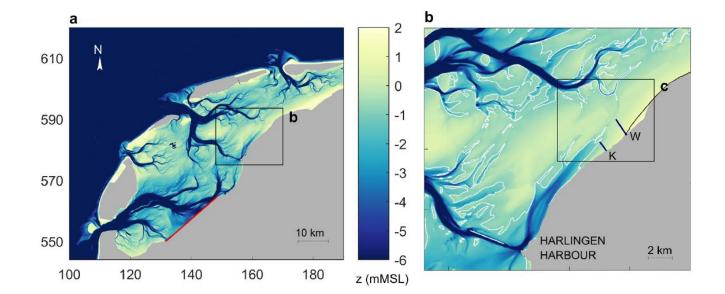
KH

WL

WH

Drying of intertidal flats

- Detailed observations of erosion during storm conditions on muddy, fringing flats along the Wadden Sea
- Four frames equiped with ADV (velocity & bed level), SSC, pressure. Converted to bed level change (BLC) & bed shear stress $\tau_{\rm b}$



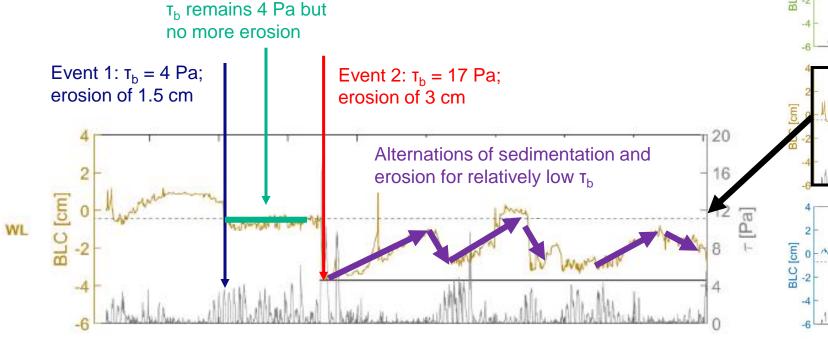
SE SW Colisomo et al., 2023 NW [s/m] 110 Erosion-Deposition alternation BLC [cm] 2 [ed] 1-Erosion = [0.3-1] cm = [0.5-2] Pa Fast Recover Erosion = 1.8 cm r_{peak} = [0.2-2] Pa C [cm] Erosion = 1.9 cm τ_{nesk} = 14.4 Pa Erosion = 3.5 cm 16 = [1-3] Pa 12 [ed] Erosion = 3.4 cr Erosion = 1.3 cm = [2.5-4] Pa No Erosion No Erosion = 10 Pa BLC [cm] Erosion = 1 cm = 18.2 Pa 12/17 12/24 12/31 01/07 01/14 01/21 01/28 02/04 02/11 time (mm/dd)

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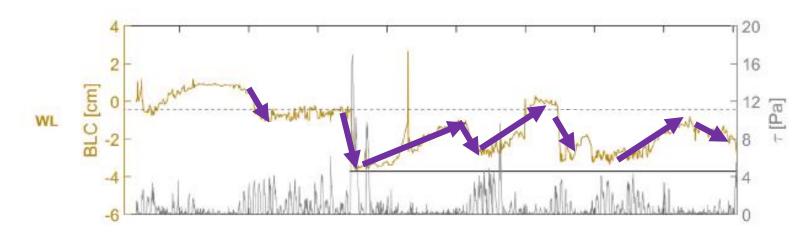
Colisomo et al., 2023 SW NW 110 Erosion-Deposition alternation BLC [cm] rosion = [0.3-1] cr = [0.5-2] Pa Fast Recover Erosion = 1.8 cm - eak = [0.2-2] Pa BLC [cm] Erosion = 1.9 cm Tnock = 14.4 Pa Erosion = 3.5 cm = [1-3] Pa BLC [cm] Erosion = 1.3 cm = [2.5-4] Pa No Erosion No Erosion BLC [cm] Erosion = 1 cm 12/17 12/24 12/31 01/07 01/14 01/21 01/28 02/04 02/11 time (mm/dd)

[m/s]

KL

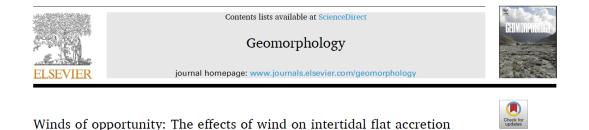
Drying of intertidal flats

- Critical bed shear stress increases with depth to values of several Pa.
- Eroded sediment is quickly replaced by new sediment
- But: this sediment has low resistance to erosion and is washed away at every new cycle with moderate wave energy (every 1-2 weeks in winter)
- What are conditions that the bed gains sufficient strength to withstand erosion? -> drying of intertidal flats

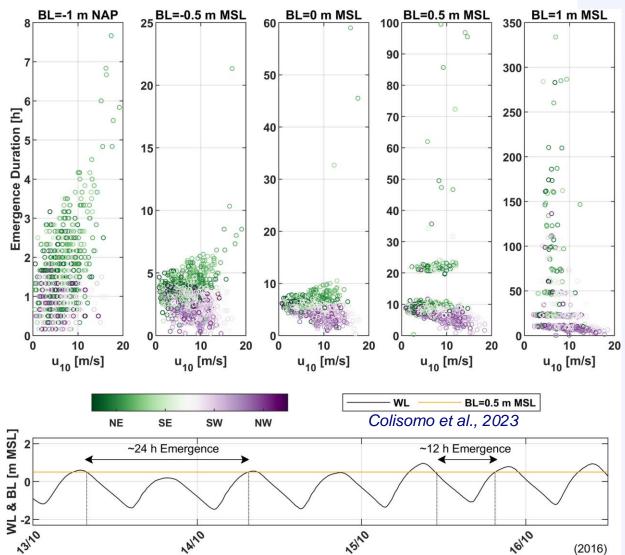


Drying of intertidal flats

- Key to gaining strength against energetic conditions is emergence of tidal flats
 - Negative pore pressures
 - Drying
- In the Wadden Sea, emergence is strongly regulated by wind effects
- The upper intertidal remains dry multiple tidal cycles during periods with wind-induced setdown



Irene Colosimo ^a, Dirk Sebastiaan van Maren ^{a, b, c,*}, Paul Lodewijk Maria de Vet ^{a, c}, Johan Christiaan Winterwerp ^a, Bram Christiaan van Prooijen ^a





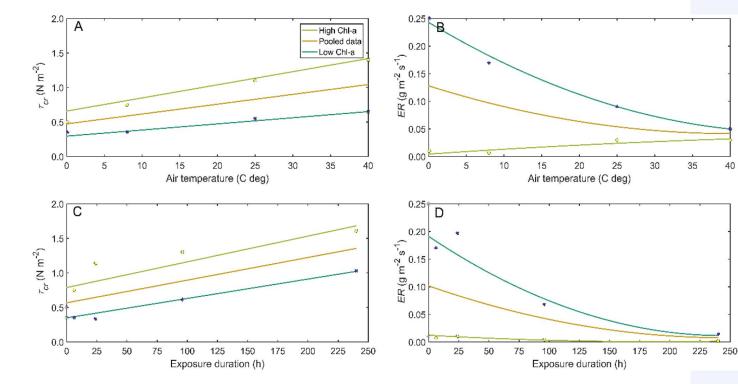
Drying of intertidal flats

Modelling effect of drying on erodibility

- 1. Empirical model of Nguyen et al. (2022)
- $\tau_{\rm cr}$ (and M) as a function of
 - Chlorofyll content
 - Exposure duration
 - Air temperature
- Implemented in the BMI environment of Delft3D-FM

$$E = M \frac{(\tau_b - \tau_{cr})}{\tau_{cr}}$$





Drying of intertidal flats

Modelling effect of drying on erodibility

- 1. Empirical model of Nguyen et al. (2022)
- $\tau_{\rm cr}$ (and M) as a function of
 - Chlorofyll content
 - Exposure duration
 - Air temperature
- Implemented in the BMI environment of Delft3D-FM
- First experiences... this method has some intrinsic instabilities

Relationships tends to enlarge

in the intertidal zone

existing topographic differences

Continental Shelf Research 245 (2022) 104802
Contents lists available at ScienceDirect
Continental Shelf Research
journal homepage: www.elsevier.com/locate/csr
Modeling the effects of aerial temperature and exposure period on
intertidal mudflat profiles

Hieu M. Nguyen^{a,*}, Karin R. Bryan^{a,b}, Zeng Zhou^{c,**}, Conrad A. Pilditch^a

High elevation \rightarrow more evaporation \rightarrow high $\tau_{cr} \rightarrow$ deposition

> Low elevation \rightarrow less evaporation \rightarrow low $\tau_{cr} \rightarrow$ easily eroded Sediment settles on nearby high point with high $\tau_{cr} \rightarrow$ erosion

Deltares

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CONTINENTAL SHELF RESEARC

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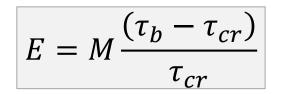
Drying of intertidal flats - modelling

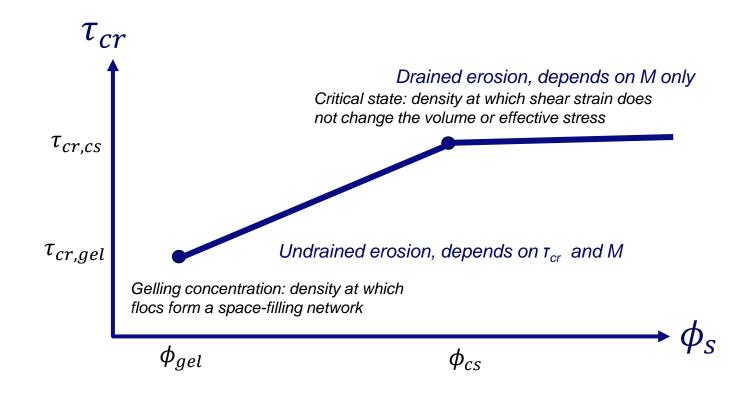
Modelling effect of drying on erodibility

- 1. Empirical model of Nguyen et al. (2022)
- 2. Darcy-based model of van Rees at al. (2024), for mixtures with mud content >30%

Effect of air exposure time on erodibility of intertidal mud flats

Floris F. van Rees 1,2,3*1 , Jill Hanssen $^{1,41},$ Stefano Gamberoni 1, Arno M. Talmon 1,4 and Thijs van Kessel 1





Effect of air exposure time on erodibility of intertidal mud flats Floris F. van Rees^{1,2,341}, Jill Hanssen^{1,41}, Stefano Gamberoni¹,

Arno M. Talmon^{1,4} and Thiis van Kessel¹

 $E = M \frac{(\tau_b - \tau_{cr})}{1 + 1}$

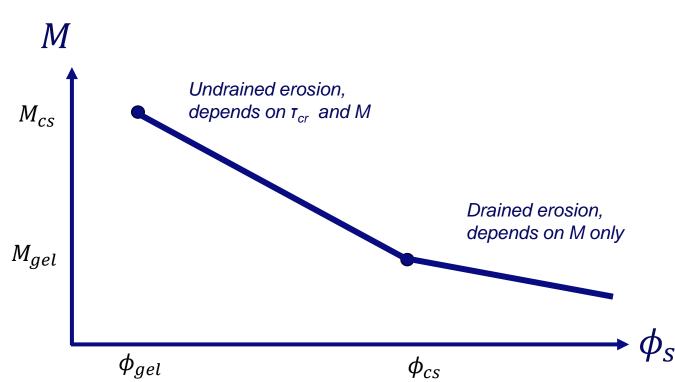
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 ϕ_{cs} computed with a Darcy model including evaporation

Input parameters for M and τ_{cr} derived from sediment properties (lab analyses)



Deltares

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Drying of intertidal flats - modelling

Effect of air exposure time on erodibility of intertidal mud flats

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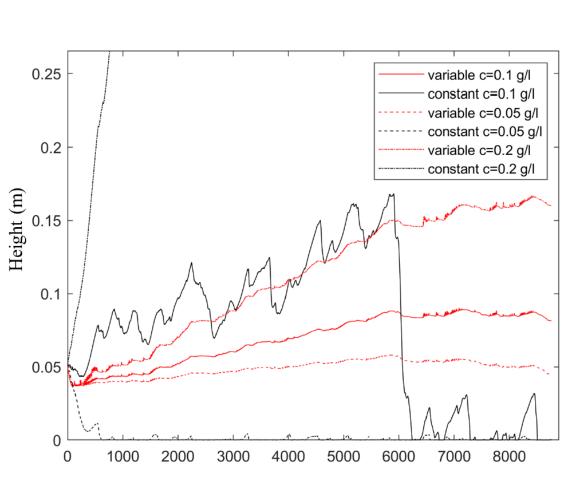
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 ϕ_{cs} computed with a Darcy model including evaporation

Input parameters for M and τ_{cr} derived from sediment properties (lab analyses)

1DV model & Delft3D-Delwaq implementation (ongoing work)



Main takeaways

- Erosion of sand-mud mixtures depends more strongly on density than on mud content
- In wave-influenced environments very muddy mixtures are easily erodible because they consolidate slowly (and therefore do not exist in the Wadden Sea)
- Sand-mud mixtures especially gain resistance against erosion during dry periods
- We have some new modelling tools accounting for these effects....

"Breaching of earthen embankments: from small-scale laboratory experiments to field measurements using photogrammetry"

(Sandra Soares-Frazão, UCLOUVAIN)

Breaching of earthen embankments From small-scale laboratory experiments to field measurements using photogrammetry

Sandra Soares-Frazão Masoumeh Ebrahimi, Nathan Delpierre, Jiangtao Yang

Outline

- Dikes, embankments, levees
- Failure modes
- What information do we need ?
- Laboratory experiments
- Field tests

Earth dams



Banasura Sagar Dam, India

Embankment dam at Diamond Valley Lake, an off-stream reservoir in Riverside County, California

River dikes



Paved bike path on top of earthen dike levee along Mississippi River in Louisiana

Sea dikes and dunes

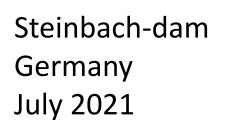


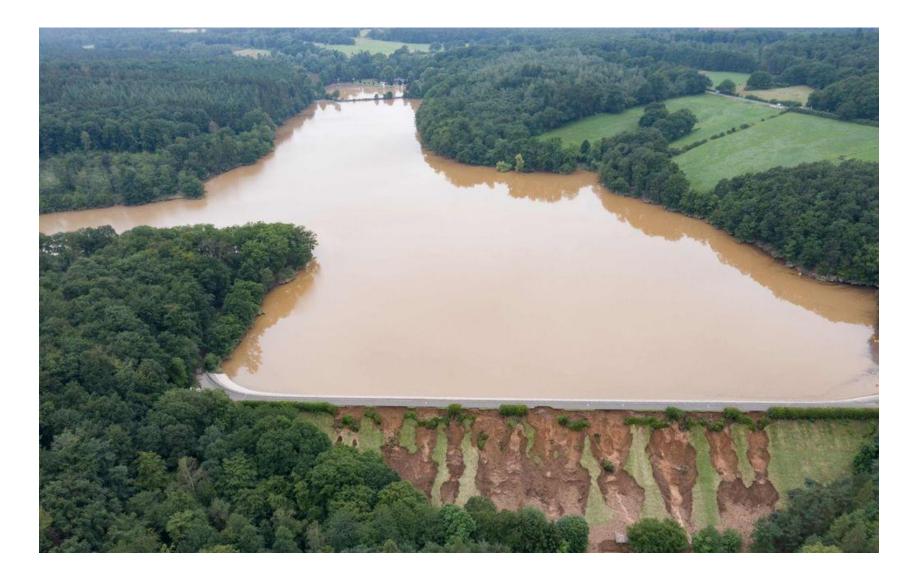
Sea-dike in Westkapelle, the Netherlands

Dunes in Belgium



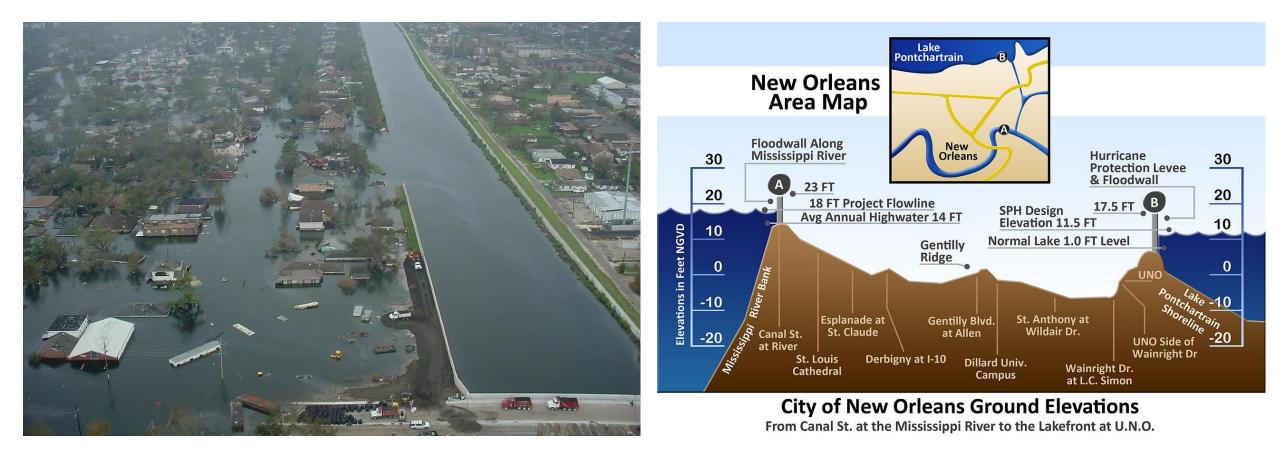
Teton dam, USA, 1976





Tous dam, Spain 1982





The 17th Street Canal, New Orleans, 2005

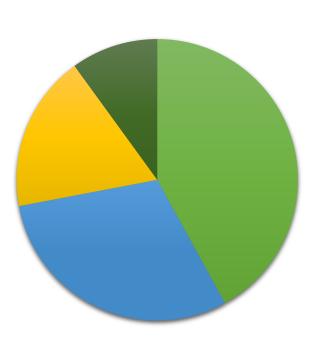


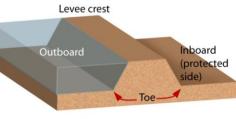




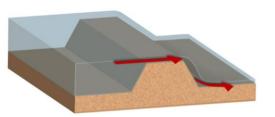
Baige landslide dam, Jinsha River, China, 2018

Failure modes

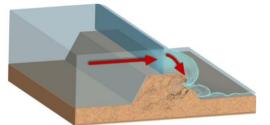




Anatomy of a levee



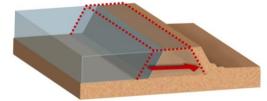
1a. Overtopping



1b. Overtopping/Jetting

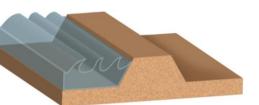


3. Surface Erosion



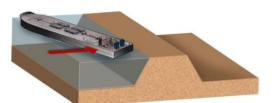
4. Sliding

Overtopping
 Seepage and piping
 Structural failure
 Others

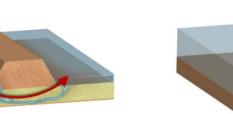


2. Internal Erosion/Piping

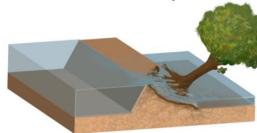
5. Wave Impacts



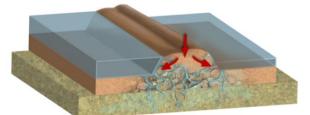
6. Structural Impacts



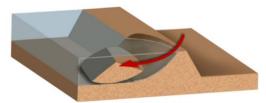
8. Piping of substratum



9. Tree damage



7. Liquefaction



10. Slope failure

Overtopping

- Non-cohesive: progressive, surface erosion
- Cohesive: mass failure, steep bank slopes





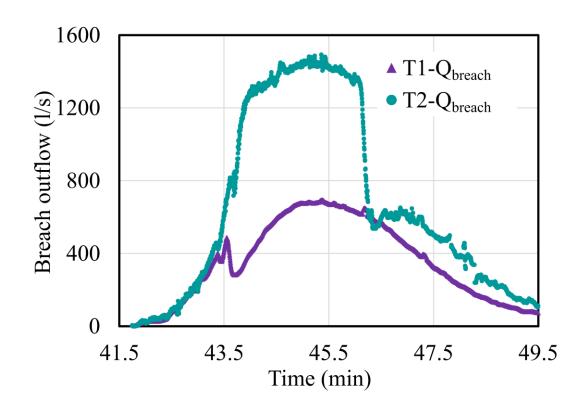
(a) Headcut erosion through a test levee [Photo: USDA-ARS-HERU Stillwater, OK. US]



Walder et al. 2015

Key variables to know

- Breach hydrograph
- Breach formation and widening
- Failure time



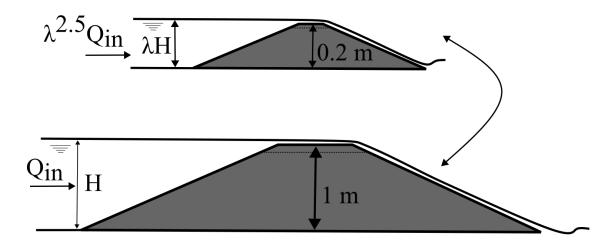


IMPACT, Mo i Rana, Norway, 2002

Laboratory experiments

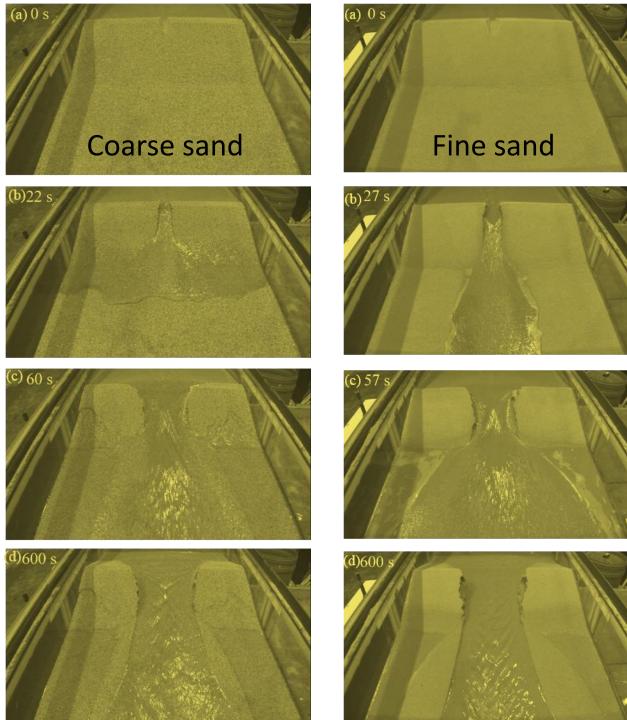
• Scale effects: Froude $\lambda_{v} = \lambda_{t} = \lambda^{1/2}$ $\lambda_{o} = \lambda^{5/2}$

Small-scale model



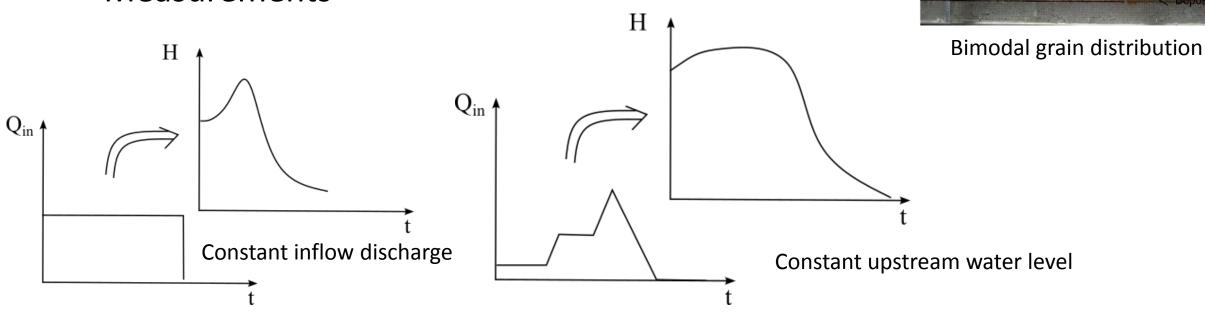
Medium-scale model

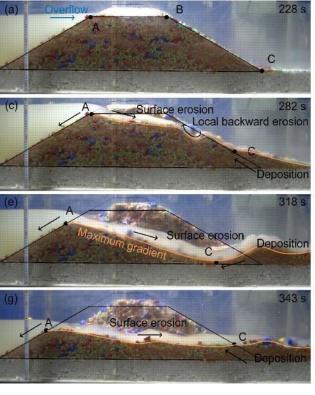
Embankment height: 0.2 m



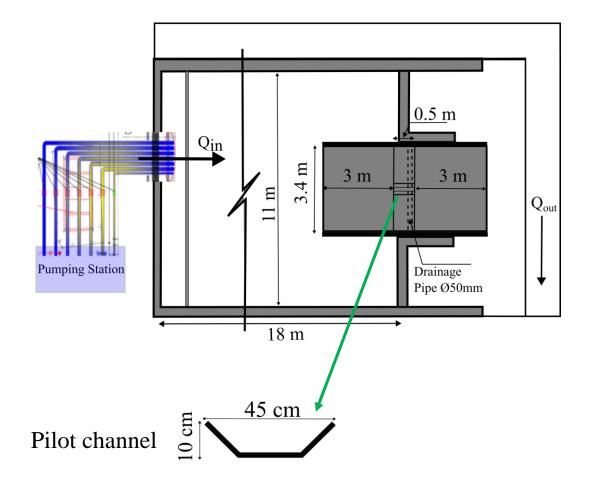
Laboratory experiments

- Dam layout
- Scale
- Dam material
- Inflow conditions
- Measurements



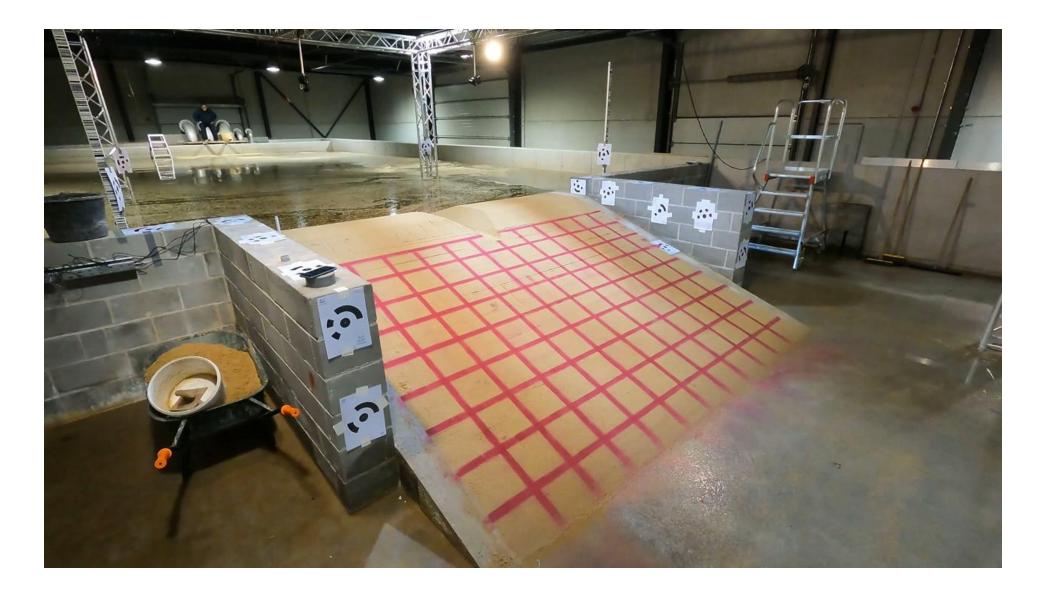


Medium-scale experiments

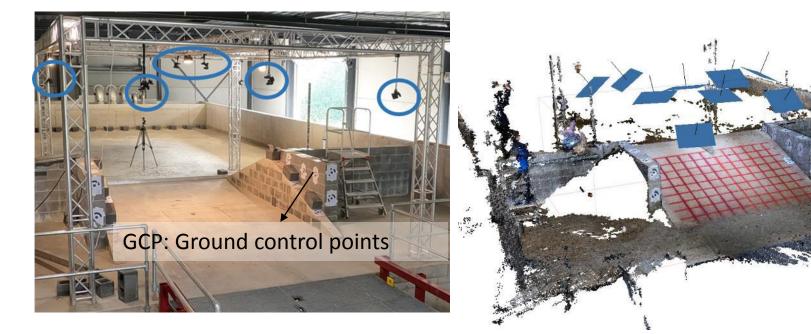


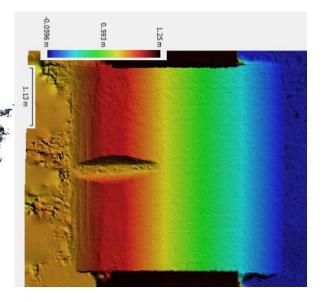


Medium-scale experiments



Photogrammetry measurements

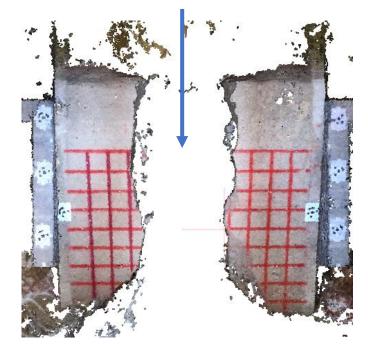




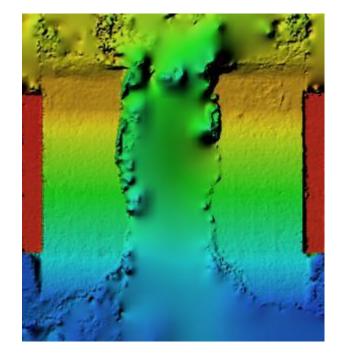
Multi-stationary synchronized cameras 10 Fixed cameras on time-lapse mode Dense Point Cloud 2,000,000 points

Elevation Map (DEM) Pixel size: 0.006 (m)

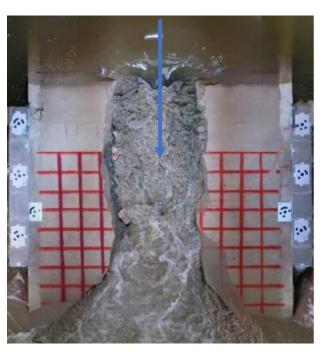
Photogrammetry measurements



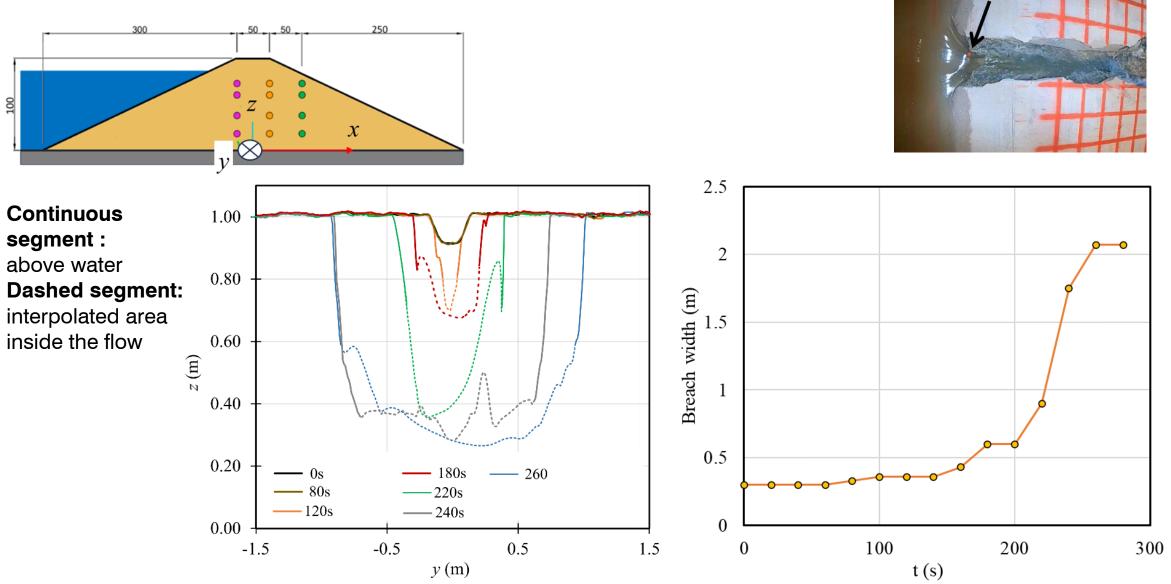
Dense Point Cloud No point or few points inside the flow



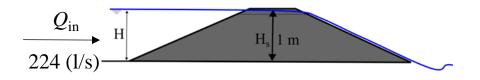
Elevation Map (DEM) Interpolation Orthophoto Resolution: 1.39 mm/pix

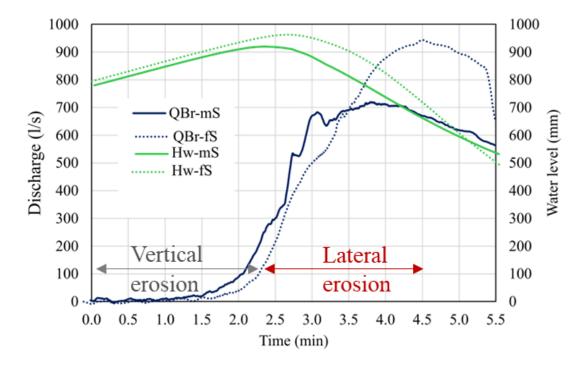


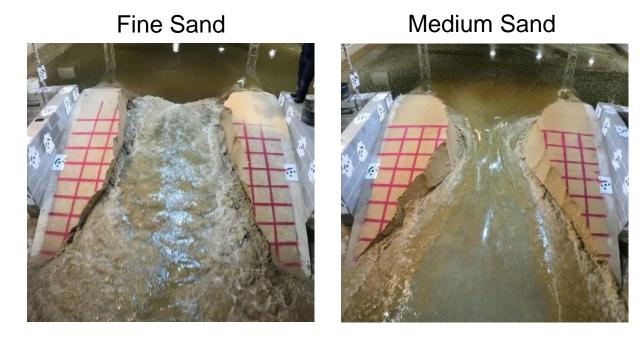
Breach evolution



Breach hydrograph





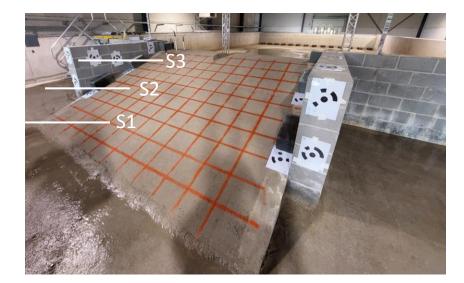


Time =3.1 min

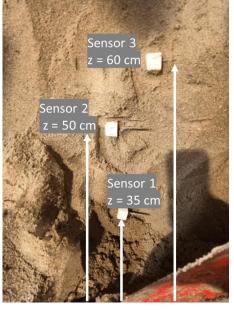
Fine Sand (fS) Medium Sand (mS)

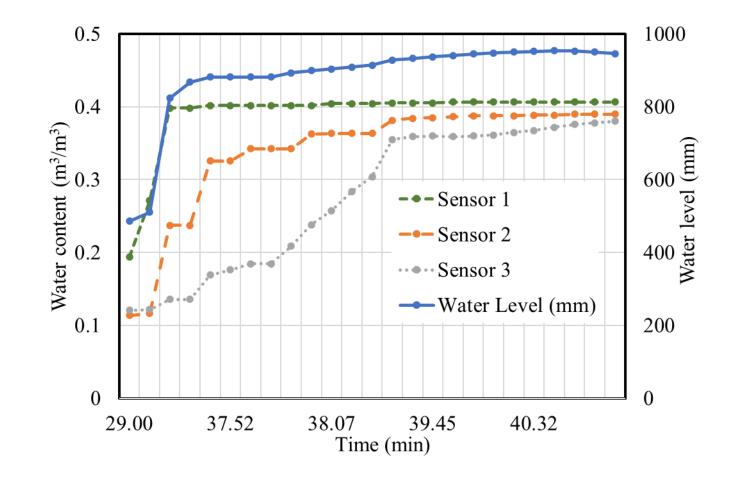
 $Q_{breach} = Q_{in} - A \frac{dz_w}{dt}$

Water content in the embankment



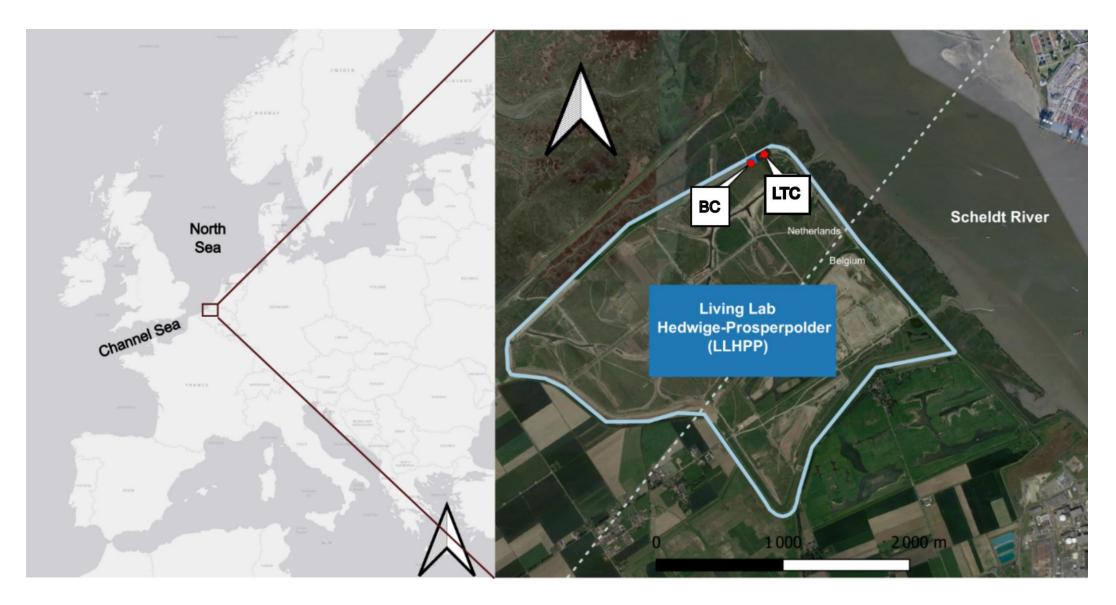
Sensors visible after dike failure







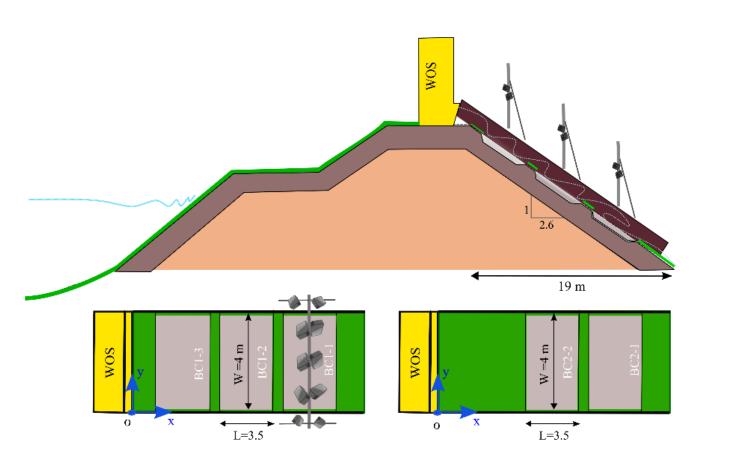
Field measurements - Polder2C's





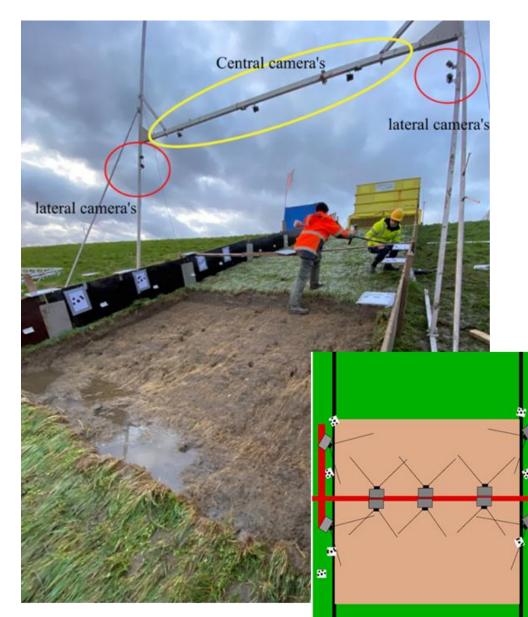
Field measurements - Polder2C's

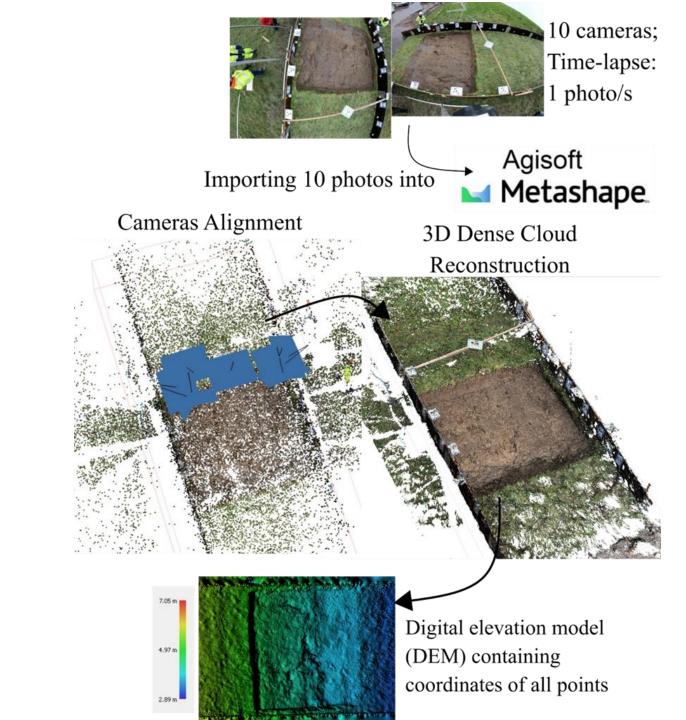
Wave overtopping tests





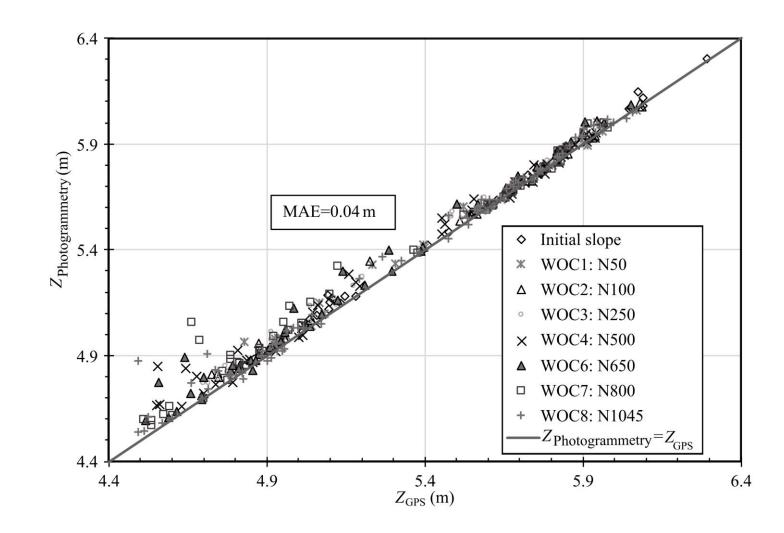
Field measurements





Field measurements - Polder2C's

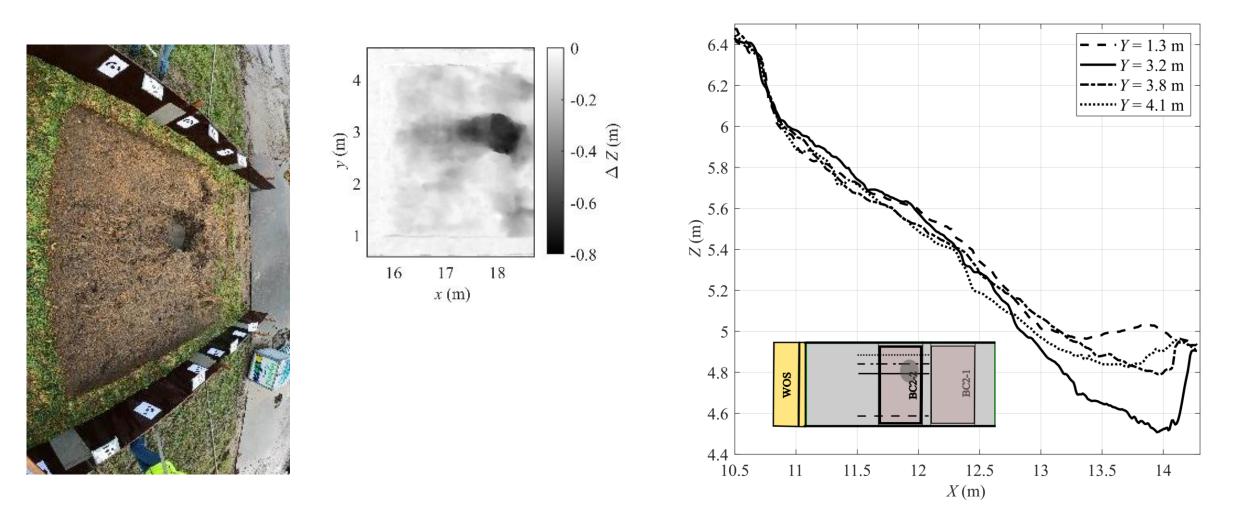




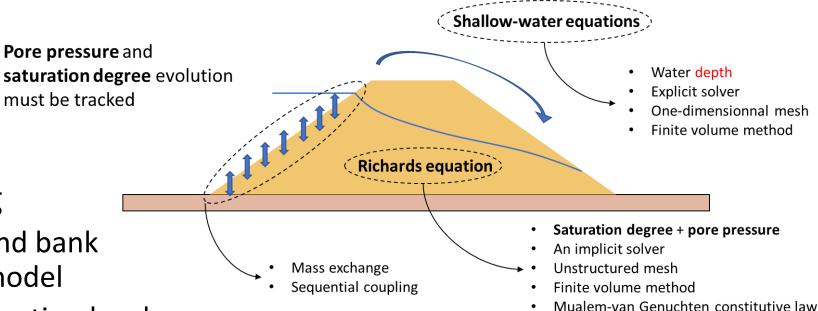


Field measurements - Polder2C's

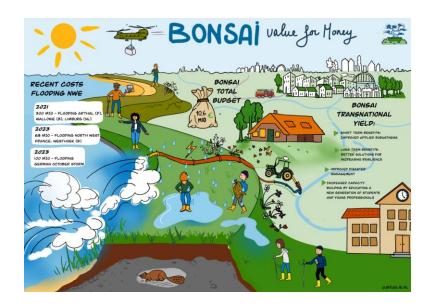
Erosion measurements



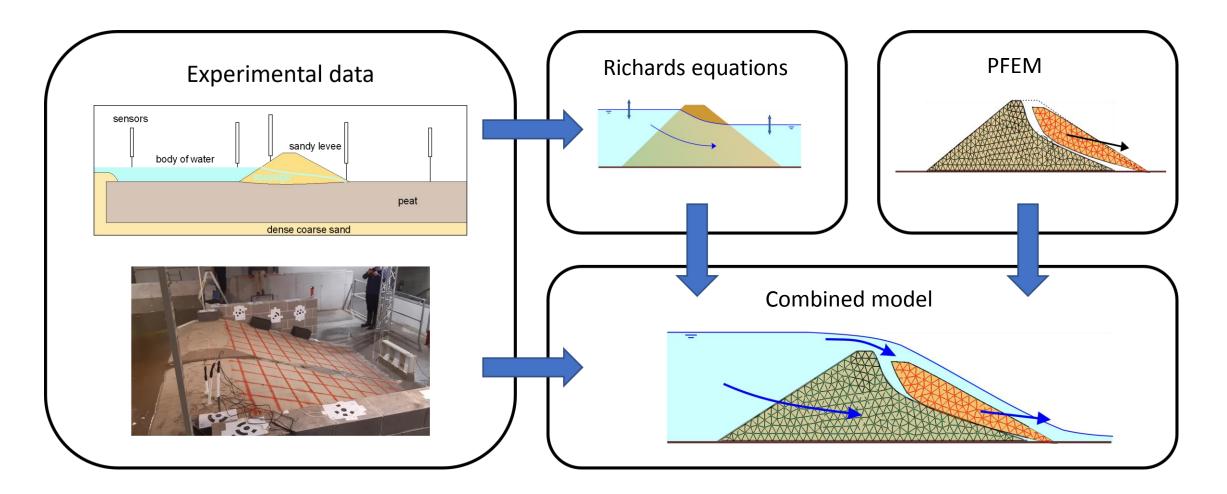
What's next ?



- Numerical modelling
 - Breaching process and bank erosion: improved model
 - Influence of soil saturation level
- Experiments
 - Small and medium scales, various materials
 - Flow though the dike
 - Velocity measurements
- Field work
 - Bonsai project

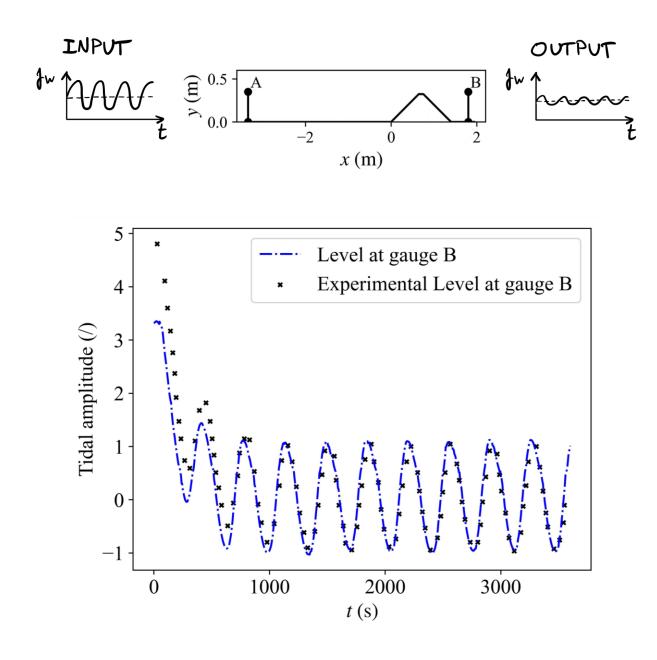


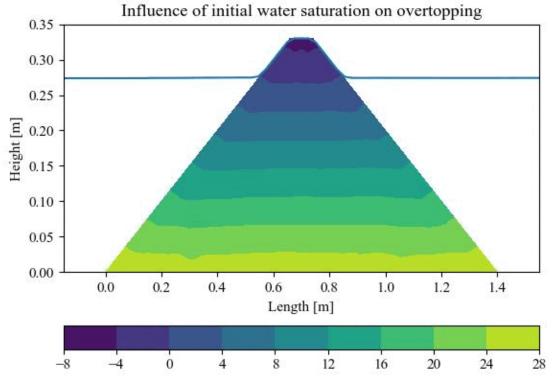
Numerical modelling



https://sites.uclouvain.be/hydraulics-group/watlab/index.html

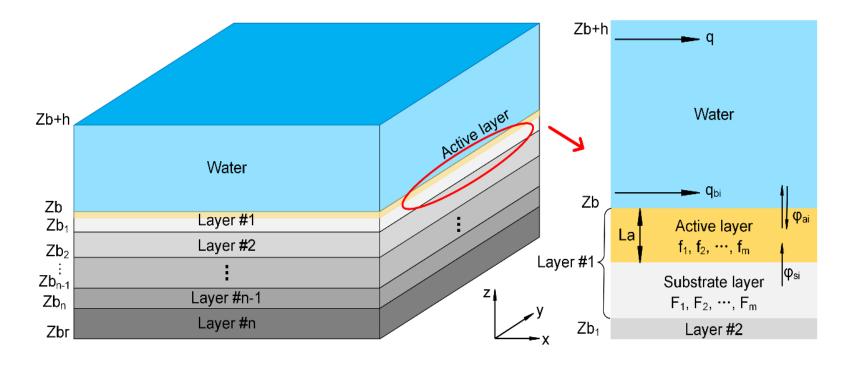
Flow through the embankment

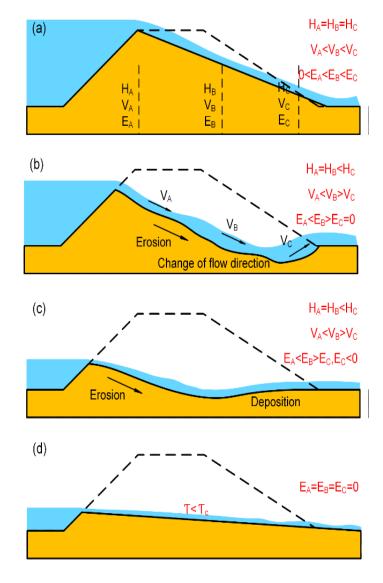




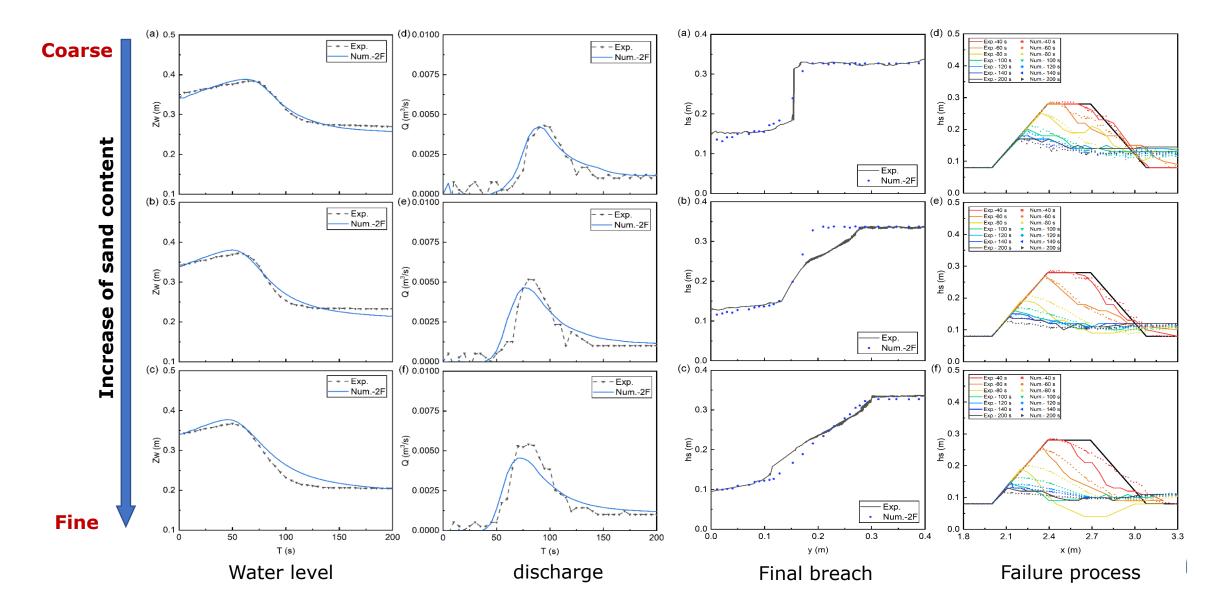
Surface erosion

2D shallow-water equations Exner equation for sediment continuity Multi-layered and non-uniform sediment transport



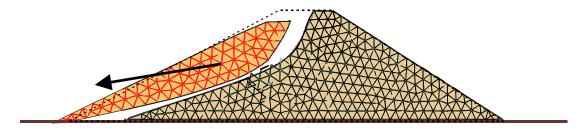


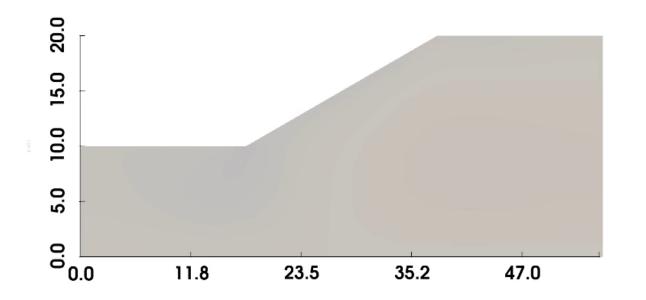
Dike erosion with different « fine » content



Mass movement using PFEM

Time (s): 0.0

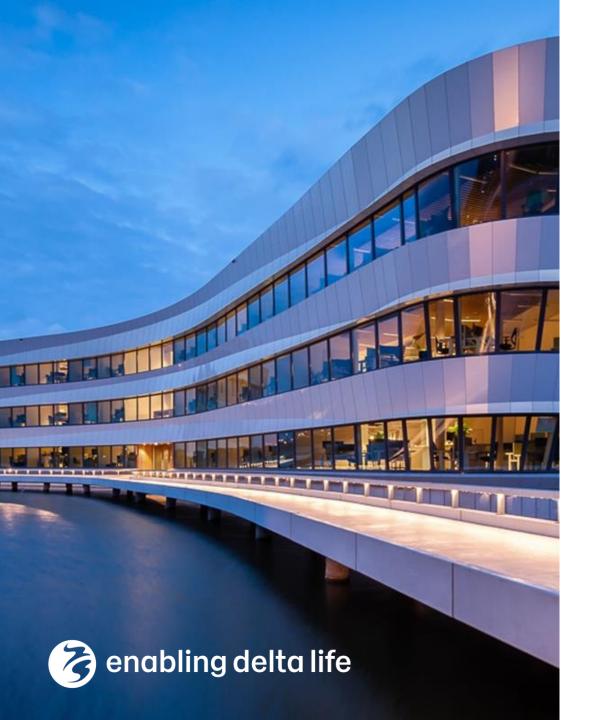






Thank you for your attention

"Deltares experimental facilities & measuring techniques" (Invited Speaker Mark Klein Breteler, Deltares)



Deltares

SEDIMARE

Experimental facilities & measurement techniques at Deltares

Mark Klein Breteler Ivo van der Werf Tijl Wijnants

Contents

- Introduction to Deltares (wave-related) facilities
 - Scheldt Flume
 - Delta Basin
 - Atlantic Basin
- Typical projects & measurements for small scale facilities
- A deeper dive into the Delta Flume

Physical modelling at Deltares

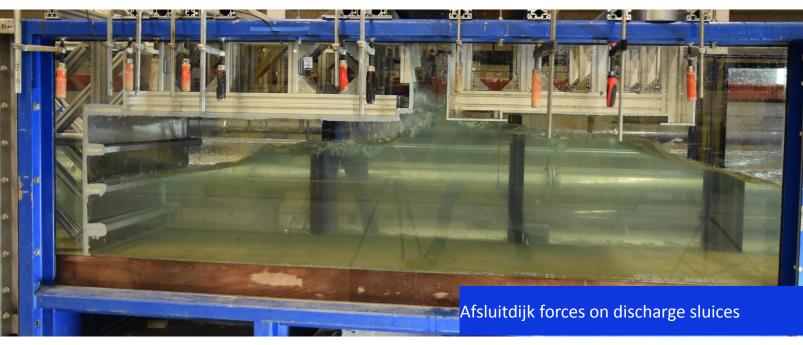
Wave flumes and wave basins:

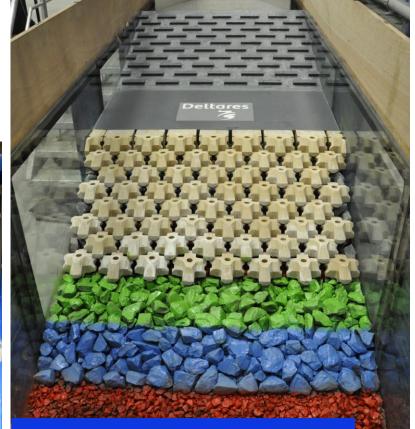
- $L(m) W(m) h(m) H_s(m)$ 2D, small scale: Scheldt Flume 110 1 ٠
- 2D & 3D, small scale: Atlantic Basin 75 8.7 1.4 0.25 ٠
- 3D, small scale: ۲
- 2D, large scale: ٠
- 300 5 9.5 2.2 Delta Flume
- 1.2 0.28 Delta Basin 50 50 1.4 0.25





Projects in the Scheldt Flume





Afsluitdijk armour stability and overtopping







Low-crested structure with dynamic cobble beach

Deltores



Projects in the Scheldt Flume







Low-crested structure with raised foreshore to lower wave overtopping

5

Projects in the Delta Basin

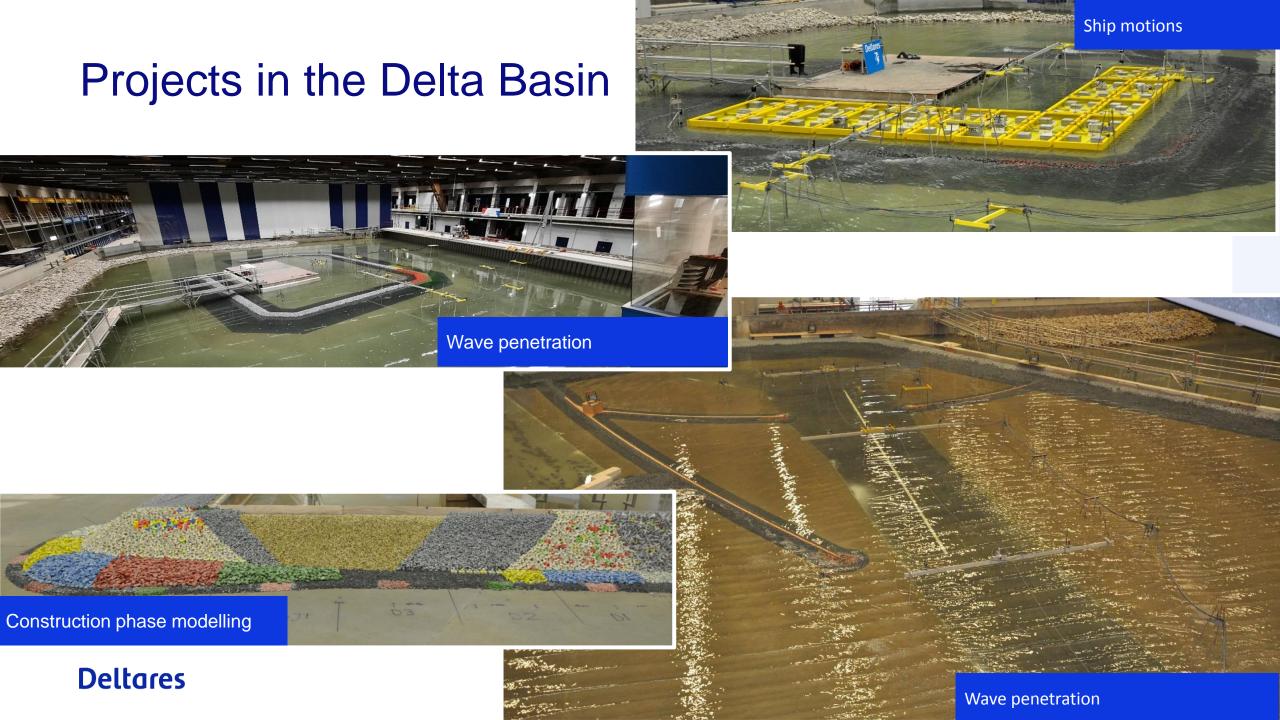
Delta Basin

Modelling of stability of groynes

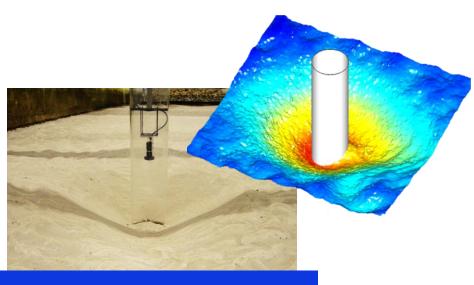
Deltares

Forces and overtopping on quay wall

Forces on deck on piles



Projects in the Atlantic Basin



Scour around monopile foundations



Deltares





Scour around jacket foundations



8

Small-scale measurement techniques

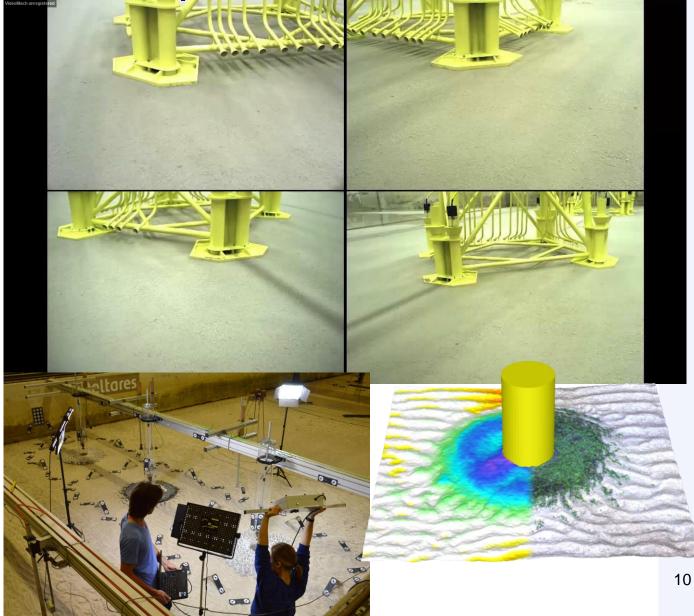
- Hydrodynamic measurements
 - Wave gauges
 - Velocity meters
- Camera measurements
 - Underwater camera
 - Internal camera
 - PIV
- 3D stereophotography
- 3D Terrestrial laser scanner
- Additional project-specific measurements
 - e.g. force sensors, run-up markers etc.





Small-scale measurement techniques

- Hydrodynamic measurements
 - Wave gauges
 - Velocity meters
- Camera measurements
 - Underwater camera
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- 3D stereophotography
- 3D Terrestrial laser scanner
- Additional project-specific measurements
 - e.g. force sensors, run-up markers etc.



Delta Flume - Specifications

- Wave height: $H_{s,max} = 2.1 \text{ m}$, $H_{max} = 4.7 \text{ m}$
- Regular waves H_{max} = 3.2 m
- Length: 300 m
- Width: 5 m
- Depth: 9.5 m
- Active reflection compensation
- Tidal water level variation during tests





Delta Flume - Examples of projects

Coastal structures:

- Wave impacts
- Vegetation (building with nature)

Energy transition: offshore wind

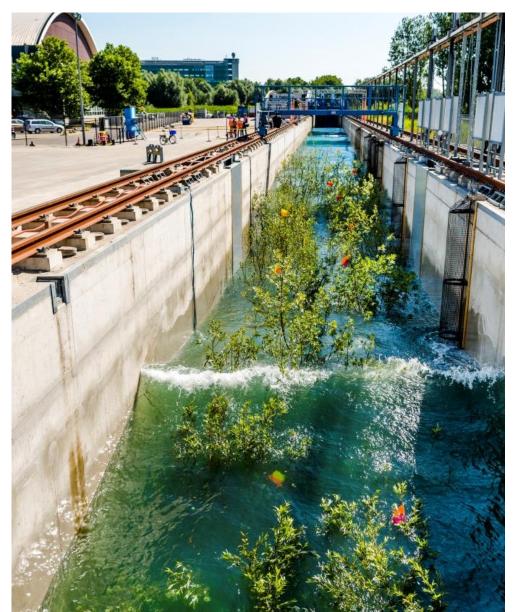
- Wave attack on offshore wind turbine substructure
- Scour protection

• Soilmechanics:

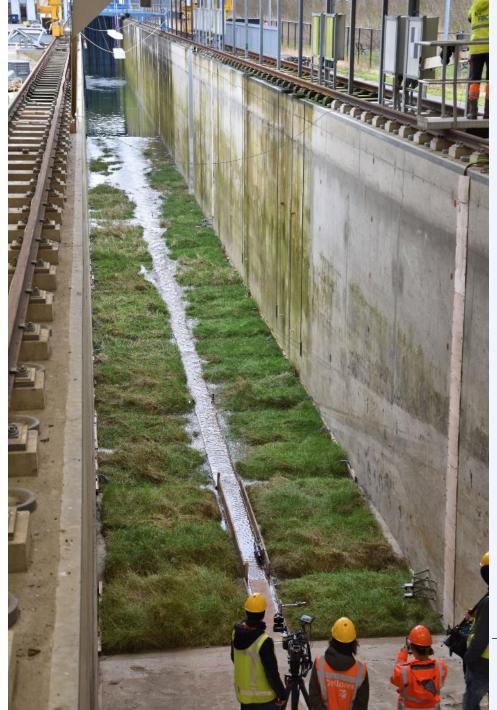
- deformations
- erosion



Building with nature







Delta Flume at Deltares

Energy transition: offshore wind

- Wave attack on offshore wind turbine • substructure
- Scour protection



Delta Flume - Examples of instruments

- Wave gauges:
 - Resistant type
 - Radac (acoustic)
 - Lidar
 - Video Al
- Flow meters
 - Electromagnetic
 - ADV and ADCP
- Pressure cells
- Accelerometers
- Faro laser scanner
- Video



Delta Flume – Example project Erosion of clay embankements by wave attack



Erosion of clay embankements by wave attack

sh-steel plate underneath



Blocks of 2x2 m², 0.85 m thick

Box without bottom,

Erosion of clay embankements by wave attack

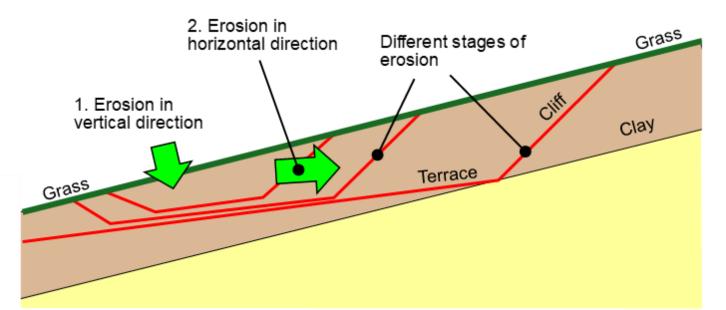


Erosion approach

Projects aimed at characterisation of erosion process

Observed erosion process:

- Phase 1: erosion grows mainly in depth
- Phase 2: terrace and cliff erosion eroded volume increases faster









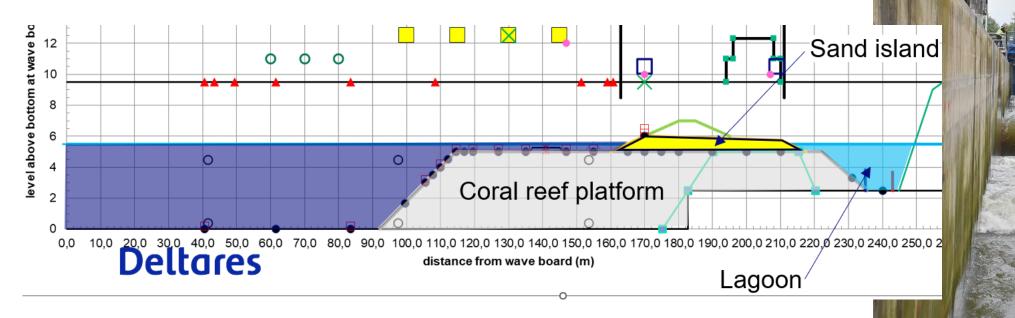
Erosion approach

Projects aimed at characterisation of erosion process

Physical modelling	Erodibility of clayErosion process	
Numerical modelling	 Extension of the Delta flume experiments Detailled information on the hydraulic loads 	
Erosion formulas	 Emprical formulas to describe the erosion rate Description of the erosion profile over time 	
Probabilistic calculations	Expected return period of erosion depthsMinimum required thickness of the clay revetment	20

Delta Flume – Example project Climate impact on atoll islands (ARISE)

- Modelling of climate impact on Atoll islands
 - Large scale because of scale effects
- Runup, overwash, erosion of island
- Validation of Xbeach numerical model



Delta Flume – Example project Climate impact on atoll islands (ARISE)

- Modelling of climate impact on Atoll islands
- Runup, overwash, erosion of island



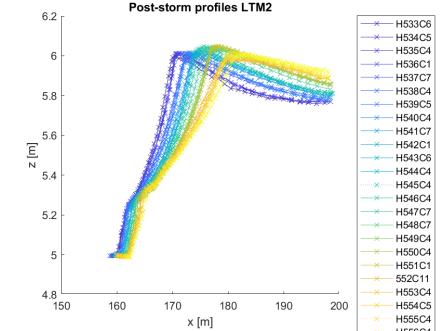
Delta Flume – Example project Climate impact on atoll islands (ARISE)

• Modelling of climate impact on Atoll islands

• Runup, overwash, erosion of island









Contact

 \times

www.deltares.nl@deltaresin linkedin.com/company/deltaresinfo@deltares.nl@deltaresf facebook.com/deltaresNL



SEDIMARE - 101072443 - D4.3: 2nd NETWORK TRAINING SCHOOL

DC Presentations



Morphodynamic Analysis of the upper confined and unconfined beach profiles during Episodic events

06/11/2024

SEDIMARE Workshop Deltares

Buckle Subbiah Elavazhagan

SEDIMARE – DC 05 IH-Cantabria, Universidad de Cantabria **Javier L. Lara** Professor, Universidad de Cantabria

María Maza Asoc. Professor, Universidad de Cantabria

Presentation on

Free long waves and its role in inaccurate numerical prediction of break bar formation





I INTRODUCTION

Shoaling zone

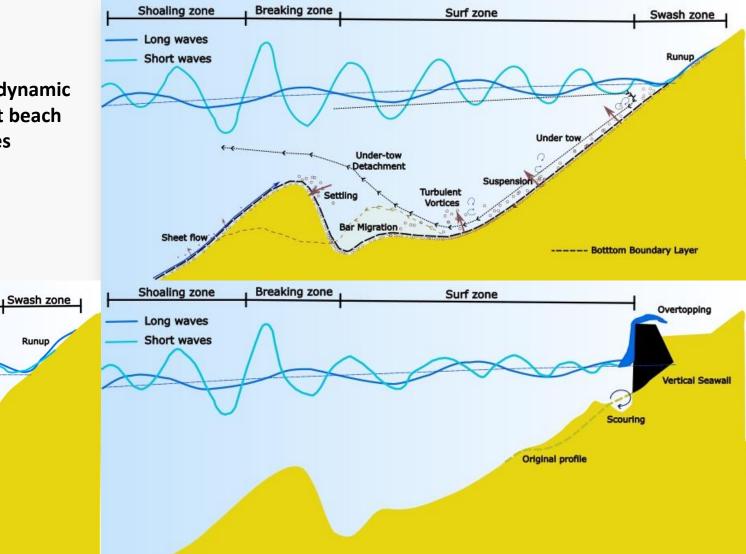
Long waves Short waves

Applying IH2VOF –SED model to model different morphodynamic processes and the governing hydrodynamics on different beach configurations and sustainable protection measures

Breaking zone

Wave and Velocity Dissipation

Surf zone



II NUMERICAL BACKGROUND

2DV RANS based solver

Turbulence is accounted using a k- ϵ closure model

Free surface reconstruction using Volume of Fluid technique

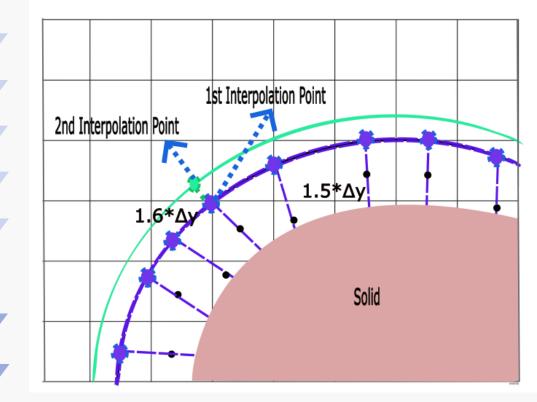
Roulond et.al 2006 Emprical formula for Bed load transport

Advective- Diffusive equation is solved for Suspended load transport

Cut-cell technique is used to model solid boundaries

Made improvements to the existing code

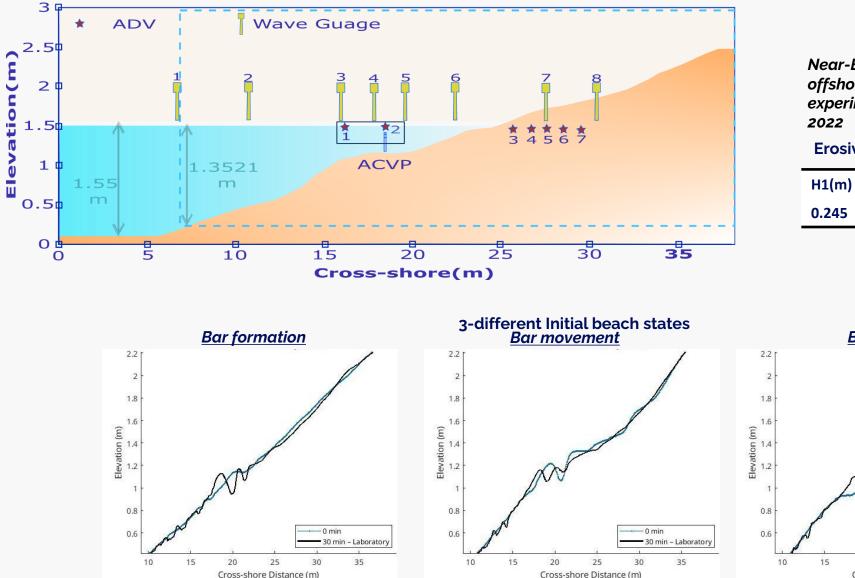
Implemented Two point Friction velocity estimation method



This approach still use log-law for estimation

But removes assumption of Roughness Length

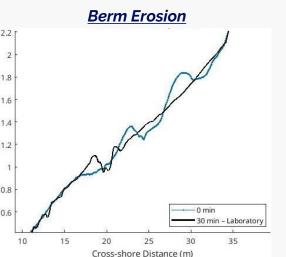
III Implementing Laboratory case



Near-Bed Sediment Transport during offshore bar migration in large-scale experiments . Florian Grossman et.al 2022

Erosive case using Bi-chromatic Waves

H1(m)	H2(m)	F1(Hz)	F2 (Hz)
0.245	0.245	0.3041	0.23657



III Implementing Laboratory case

Experimental Limitation

- 1. First Order Wave generation
- 2. Significant Presence of Free long waves
- 3. No active wave absorption

Data Limitation

- 1. Didn't get wave paddle motion signals
- 2. Certain ADV and ADCP data are lacking due to unexpected sedimentation during the experiments

J.W.M. Kranenborg et.al, 2024

Numerical Capabilities

- **1**. First and second order waves as boundary condition
- 2. Active wave absorption with possibility of disabling it

Numerical Limitation

- 1. 2D model
- 2. Difficulty in replicating entire length of the flume
- 3. High computational cost

1. Modelling entire length of the flume would be

challenging for a Sediment transport model without

coupling with a hydrodynamic model

2. No paddle signals means measured wave gauge

data at a certain point would be used as a wave

input

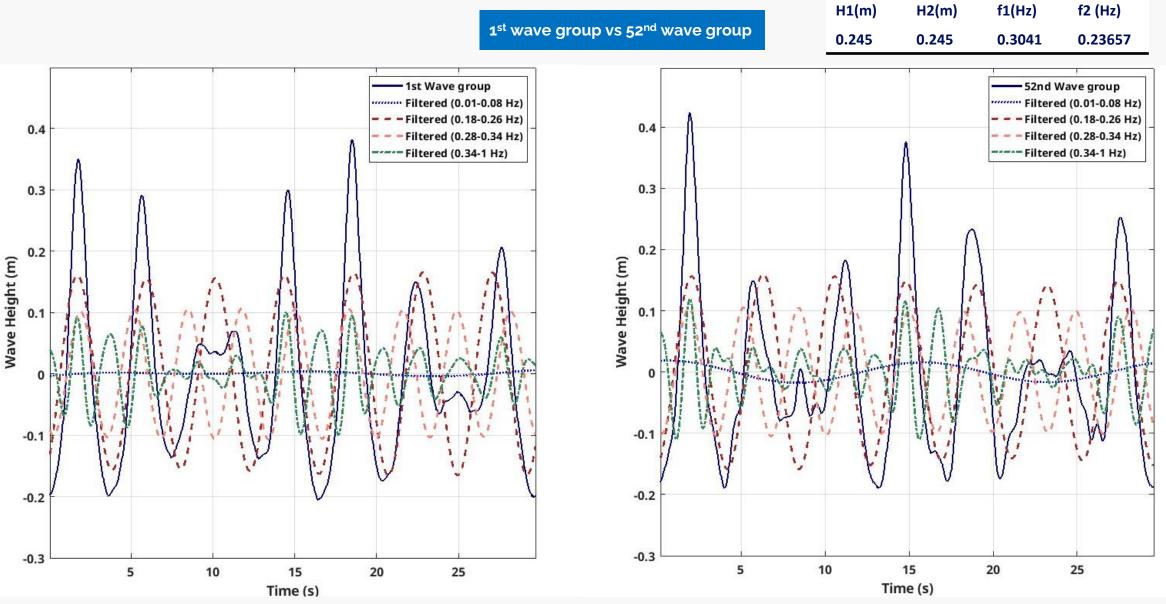
Effects of free surface modelling and wave-breaking turbulence on depth-resolved modelling of sediment transport in the swash zone, Coastal Engineering.

Joep van der Zanden et,al, 2019a

Sand transport processes and bed level changes induced by two alternating laboratory swash events, Coastal Engineering.

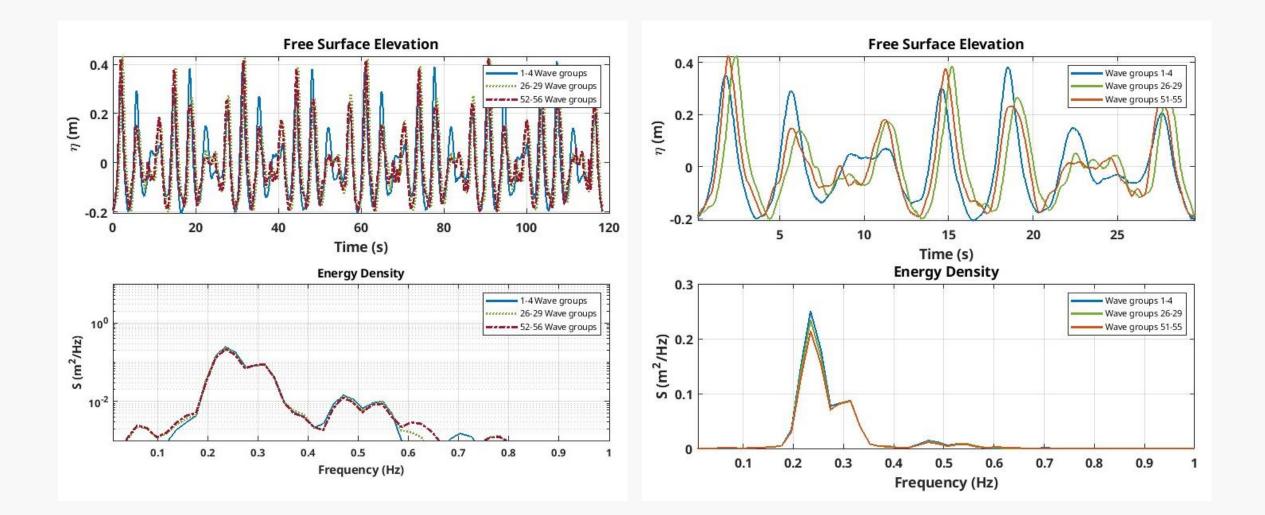
IV Free vs Bound Long waves

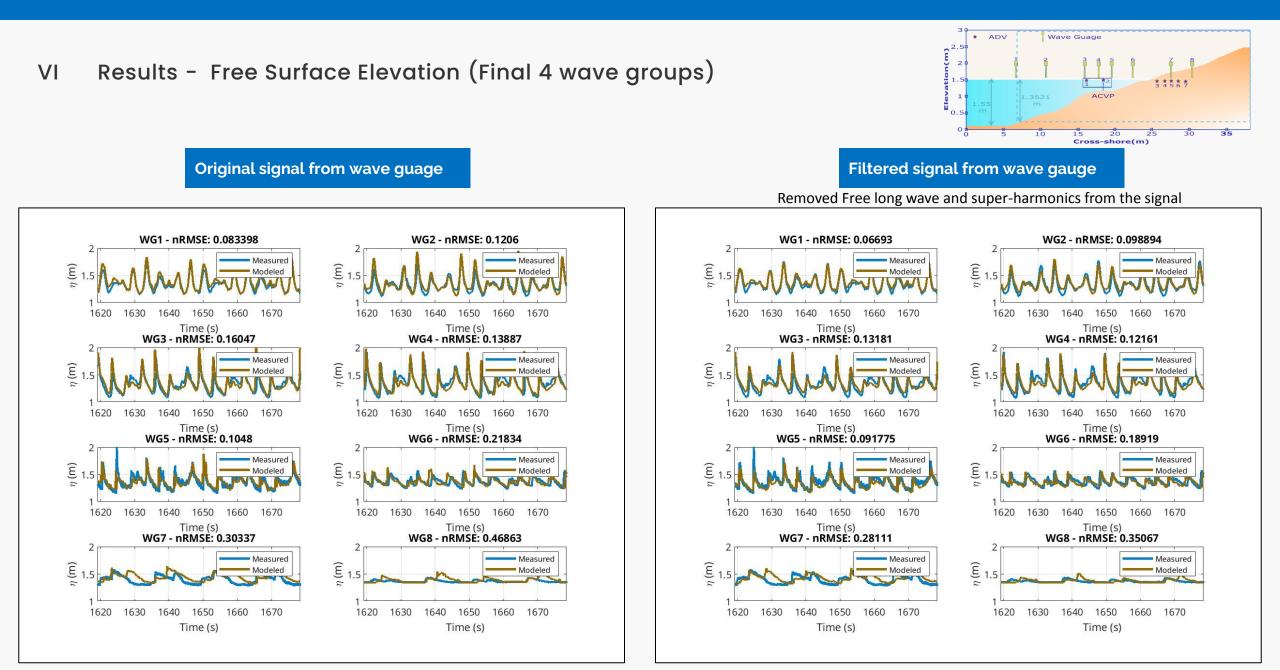
	Bound Long Waves	Free Long Waves
<u>Source</u>	Generated at the wave paddle bound to the primary wave group using 2 nd order wave generation	Generated Freely in the lab through non linear interactions of primary wave group and wave reflections
<u>Dependence</u>	Tied to primary wave group, travels at wave group celerity	Independent, self-propagating
<u>Trigger for</u> Existence	After Wave breaking becomes free	Primary wave group interaction

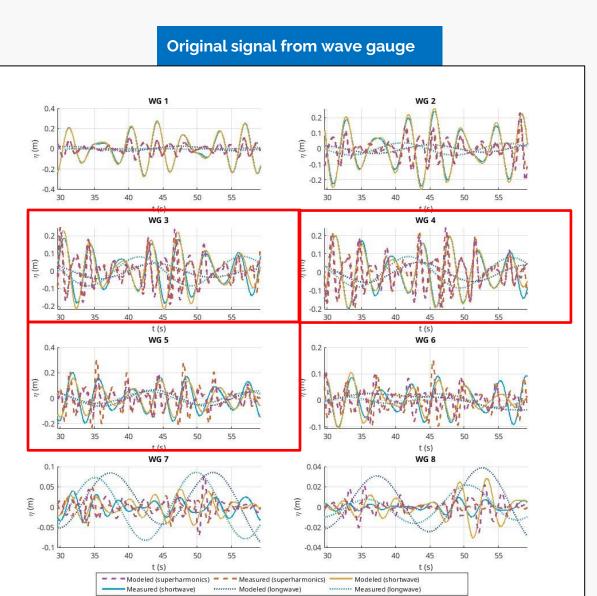


V Free Wave Generation in the Lab

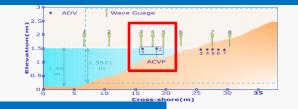
Free Wave Generation in the Lab V



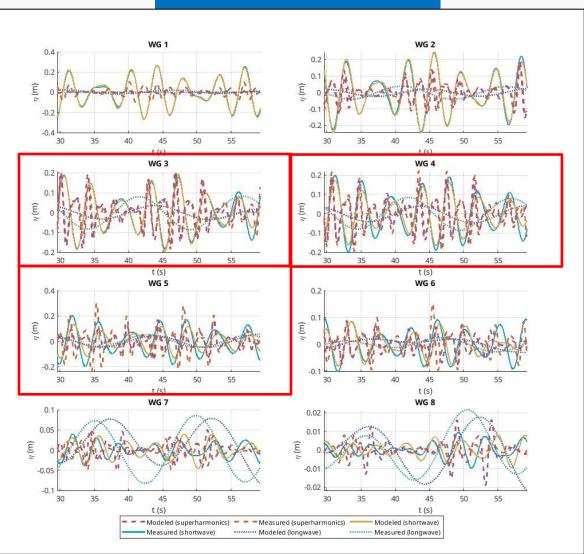




VI Results - Wave components(First 4 wave groups)



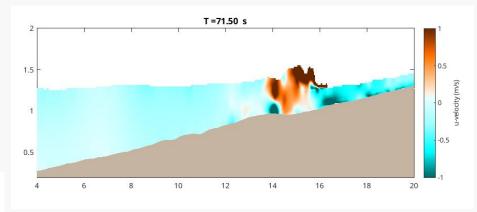
Filtered signal from wave gauge



Twin velocity concentration points at the bed potentially leading to formation of two troughs

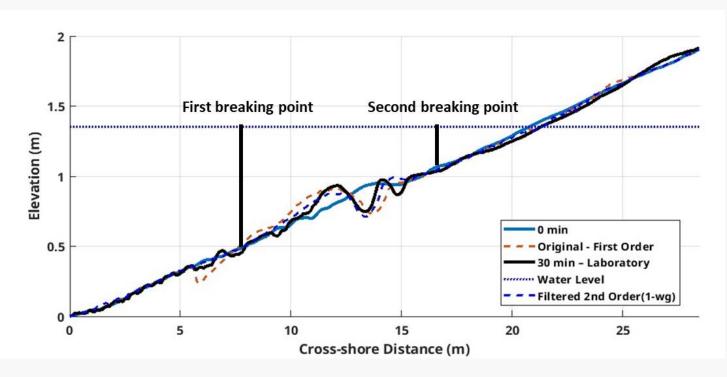
$\mathbf{r} = \mathbf{55.50 \ s}$

Single concentration point, potentially corresponding the larger of the two troughs

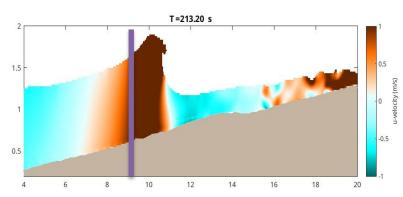


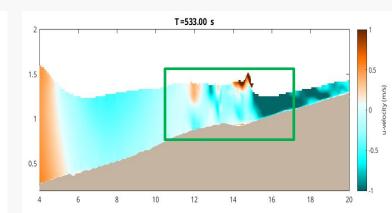
- **1.** First Breaking point is correctly modelled and thus formation of the primary bar
- 2. Second breaking influences the dynamics of swash zone and also the formation of the smaller bar
- 3. Strong off-shore directed current means, single concentration point on the bed or wave dissipates without significant impact to the bed

VI Results - Bed Level Evolution

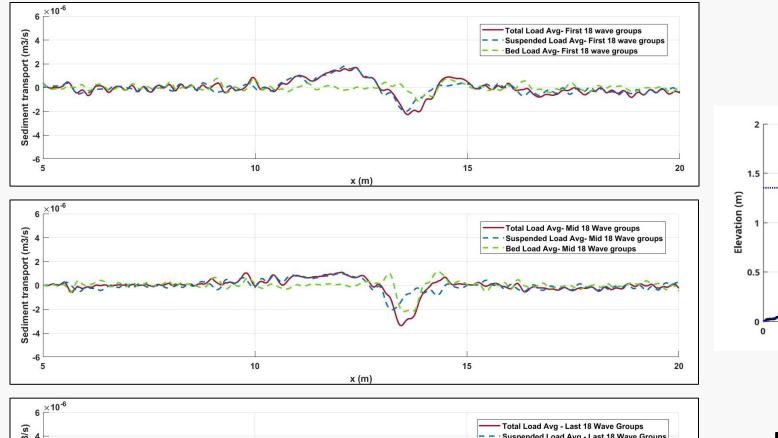


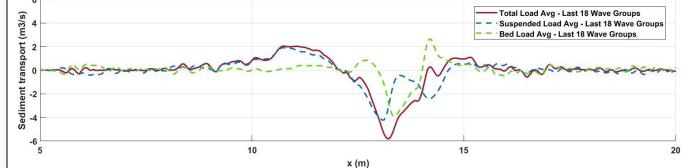
Desirable Primary breaking point



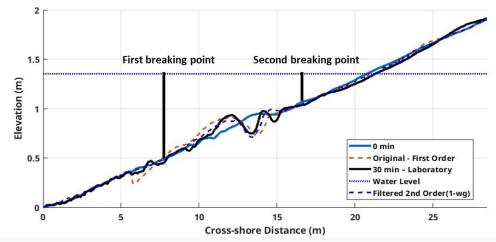


Results - Sediment Transport VI

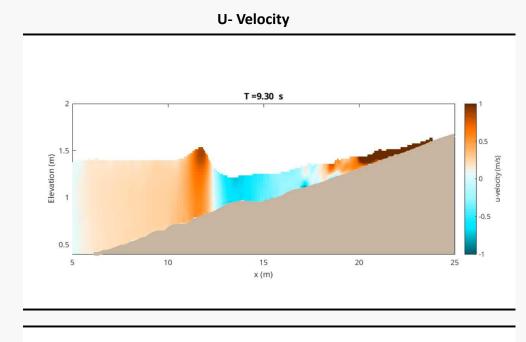


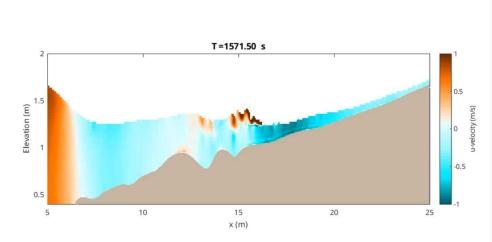


evolution Representation of of Sediment transport over 30minutes, in the Filtered + Second Order case

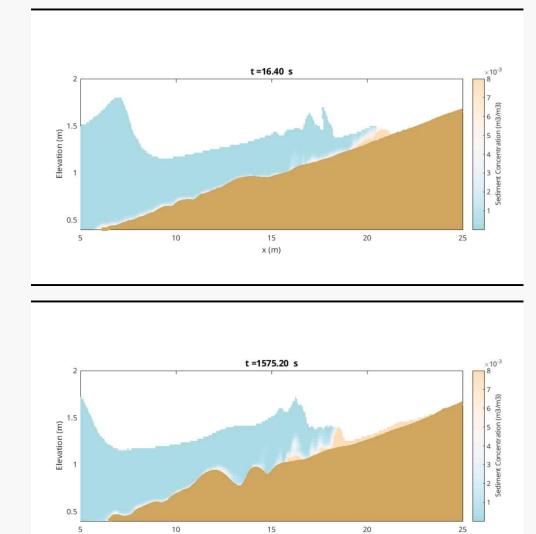


VI Results - Velocity and Sediment concentration during first 4 and final 4 wave groups



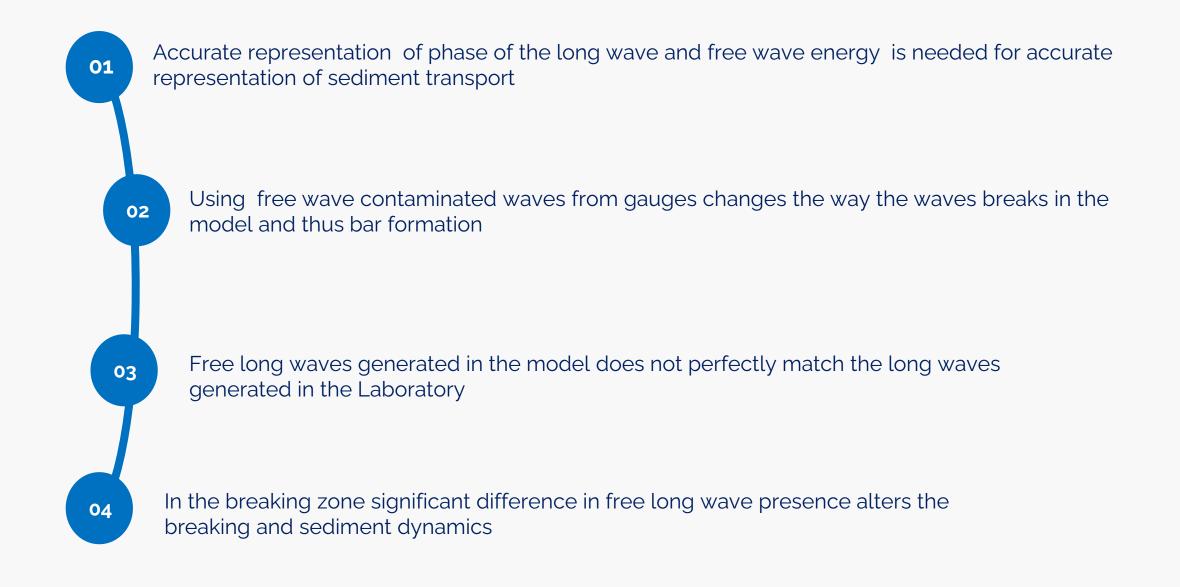


Sediment Concentration

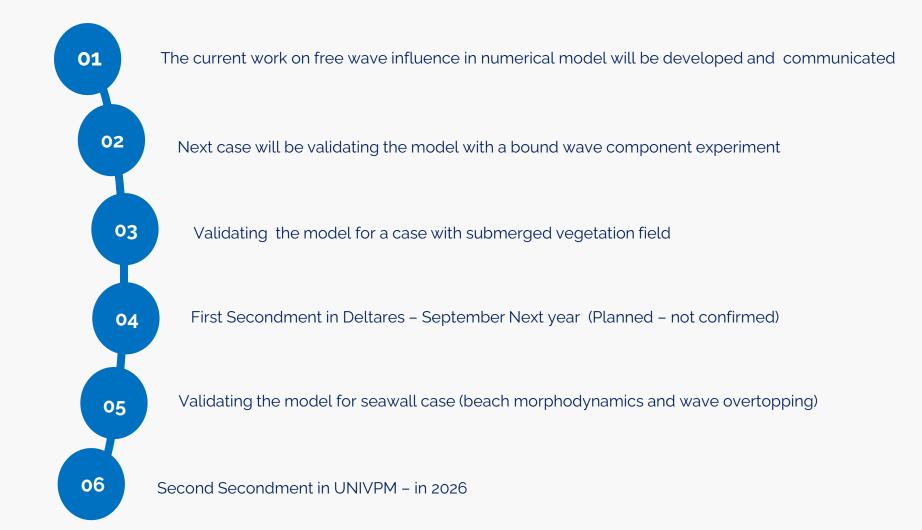


x (m)

VII Conclusion



VIII What's Next





Morphodynamic Analysis of the upper confined and unconfined beach profiles during Episodic events

06/11/2024Buckle Subbiah ElavazhaganJavier L. LaraSEDIMARE WorkshopSEDIMARE - DC 05CantabriaDeltaresIH-Cantabria, Universidad de
CantabriaMaría Maza
Asoc. Professor, Universidad de
Cantabria

Presentation on

Free long waves and its role in inaccurate numerical prediction of break bar formation

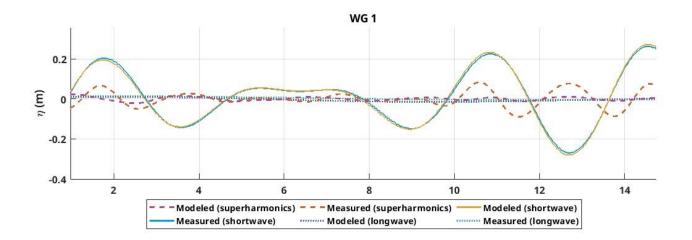
Currently in Month 9



Thank you!

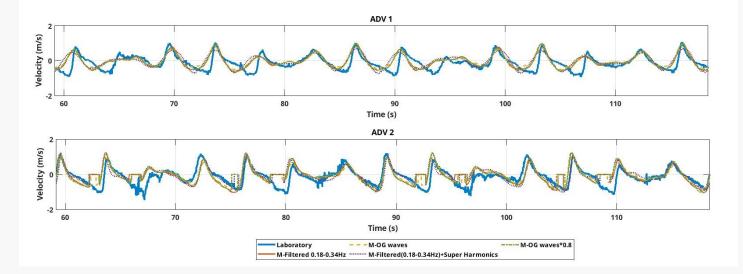
References

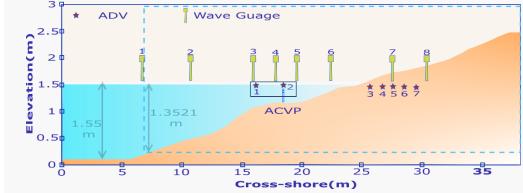
- 1. Astudillo, C., Gracia, V., Cáceres, I., Sierra, J. P., & Sánchez-Arcilla, A. (2022). Beach profile changes induced by surrogate Posidonia Oceanica: Laboratory experiments. *Coastal Engineering*, 175. https://doi.org/10.1016/j.coastaleng.2022.104144
- 2. Astudillo-Gutierrez, C., Gracia, V., Cáceres, I., Sierra, J. P., & Sánchez-Arcilla, A. (2024). Influence of seagrass meadow length on beach morphodynamics: An experimental study. *Science of the Total Environment*, *921*. https://doi.org/10.1016/j.scitotenv.2024.170888
- 3. García-Maribona, J., Lara, J. L., Maza, M., & Losada, I. J. (2021). An efficient RANS numerical model for cross-shore beach processes under erosive conditions. Coastal Engineering, 170. https://doi.org/10.1016/j.coastaleng.2021.103975
- 4. García-Maribona, J., Lara, J. L., Maza, M., & Losada, I. J. (2022). Analysis of the mechanics of breaker bar generation in cross-shore beach profiles based on numerical modelling. Coastal Engineering, 177. https://doi.org/10.1016/j.coastaleng.2022.104172
- 5. Grossmann, F., Hurther, D., Sánchez-Arcilla, A., & Alsina, J. M. (2023). Influence of the Initial Beach Profile on the Sediment Transport Processes During Post-Storm Onshore Bar Migration. Journal of Geophysical Research: Oceans, 128(4). <u>https://doi.org/10.1029/2022JC019299</u>
- 6. Lara, J. L., Losada, I. J., & Guanche, R. (2008). Wave interaction with low-mound breakwaters using a RANS model. Ocean Engineering, 35(13), 1388–1400. https://doi.org/10.1016/j.oceaneng.2008.05.006
- Lara, J. L., Ruju, A., & Losada, I. J. (2011). Reynolds averaged Navier-Stokes modelling of long waves induced by a transient wave group on a beach. Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences, 467(2129), 1215–1242. https://doi.org/10.1098/rspa.2010.0331
- 8. Maza, M., Lara, J. L., & Losada, I. J. (2013). A coupled model of submerged vegetation under oscillatory flow using Navier-Stokes equations. Coastal Engineering, 80, 16–34. <u>https://doi.org/10.1016/j.coastaleng.2013.04.009</u>
- 9. Roulund, A., Sumer, B. M., Fredsøe, J., & Michelsen, J. (2005). Numerical and experimental investigation of flow and scour around a circular pile. Journal of Fluid Mechanics, 534, 351–401. <u>https://doi.org/10.1017/S0022112005004507</u>
- Streicher, M. ;, Kortenhaus, A. ;, Altomare, C. ;, Gruwez, V. ;, Hofland, B. ;, Chen, X. ;, Marinov, K. ;, Scheres, B. ;, Schüttrumpf, H. ;, Hirt, M., Cappietti, L., Esposito, A., Saponieri, A., Valentini, N., Tripepi, G., Pasqualini, D., Di Risio, M., Aristodemo, F., Damiani, L., & Kaste, . . (2017). (Wave Loads on Walls) Large-Scale Experiments in the Delta Flume.

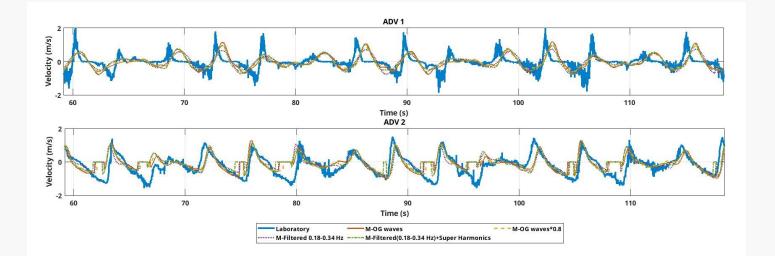


Free waves are filtered from the input wave, but there also frequencies generated in the model.

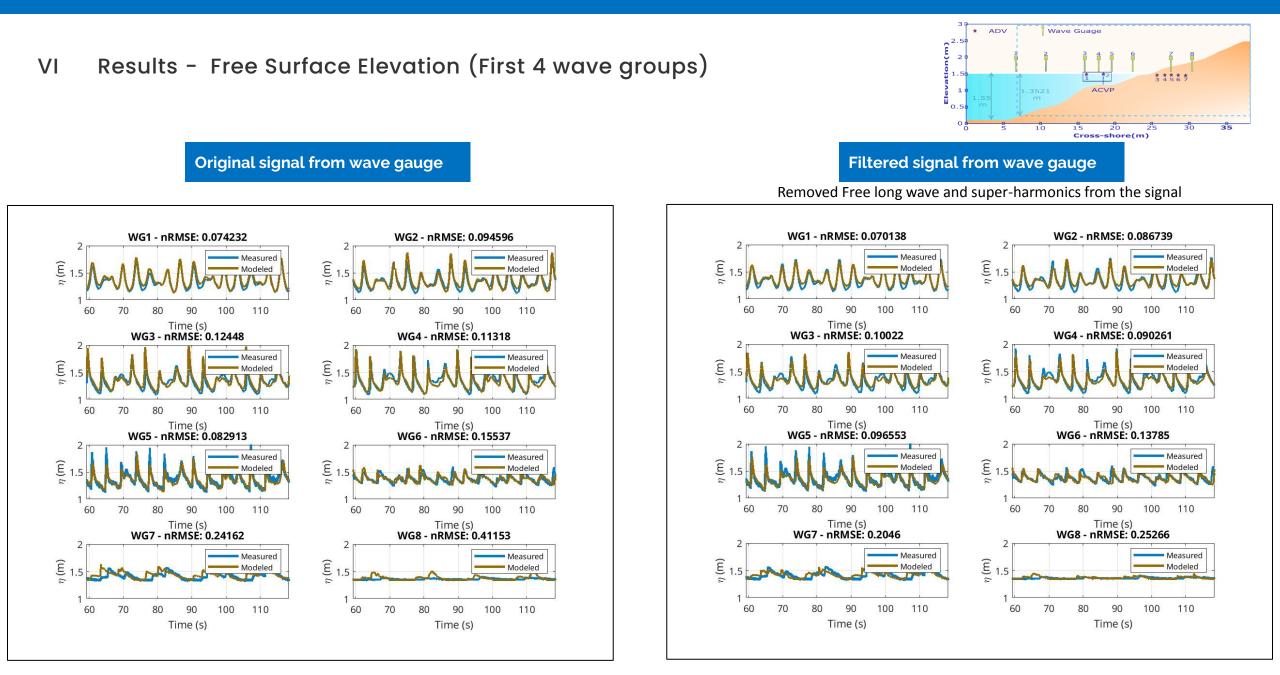
VI Results - Velocity



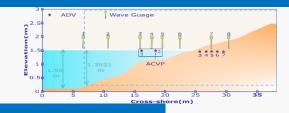




- 1. No significant changes in the velocities for the change in wave conditions
- 2. Meaning wave breaking occurs quite similarly across the different cases.
- 3. RMSE values varies between 04 to 0.6 for the different cases

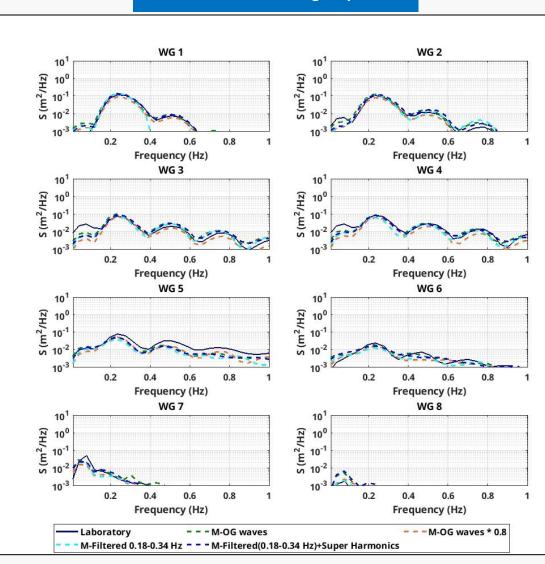


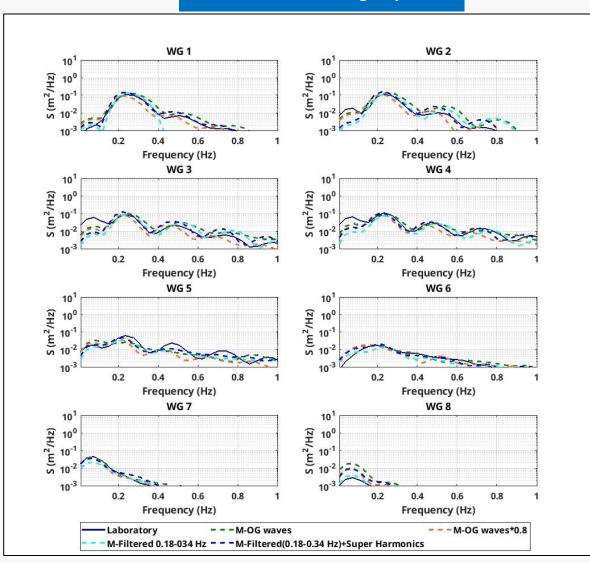
VI Results - Energy Density

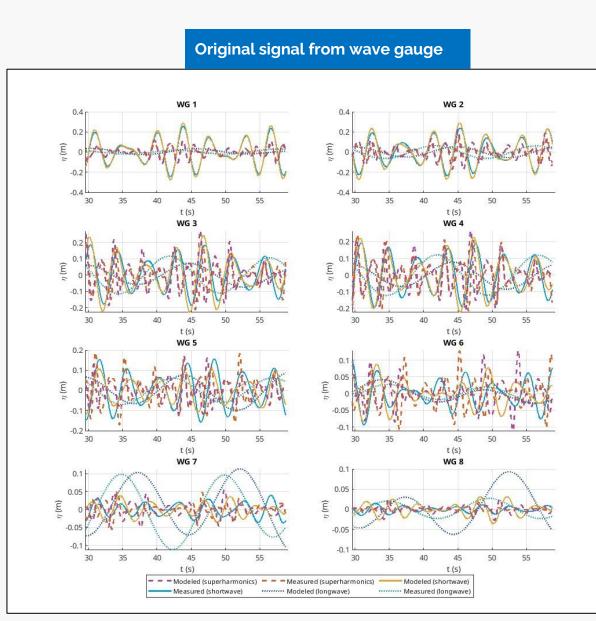


First Four Wave groups

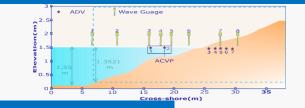
Final Four Wave groups







VI Results - Wave Components(Final 4 wave groups)

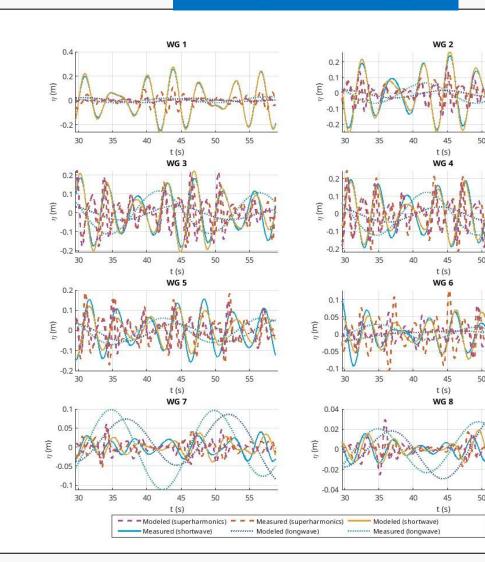


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Filtered signal from wave gauge



Initiation of motion of sand-mud beds

Paterno "Jowi" S. Miranda IV, MCE

Research update



in patrno4





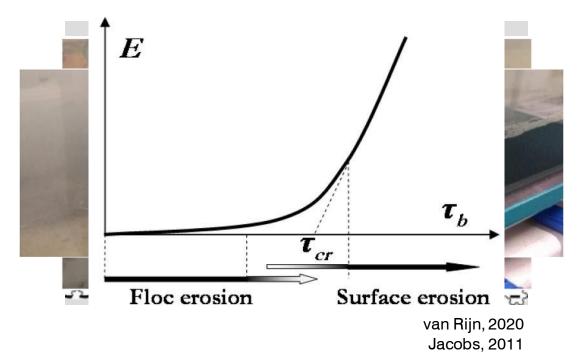
Contents

- Sand-mud erosion
- Initiation of motion of sediment classes (e.g. sand, silt, clay)
- Framework for initiation of motion of sand-mud
- Application of framework to collected data
- Next steps



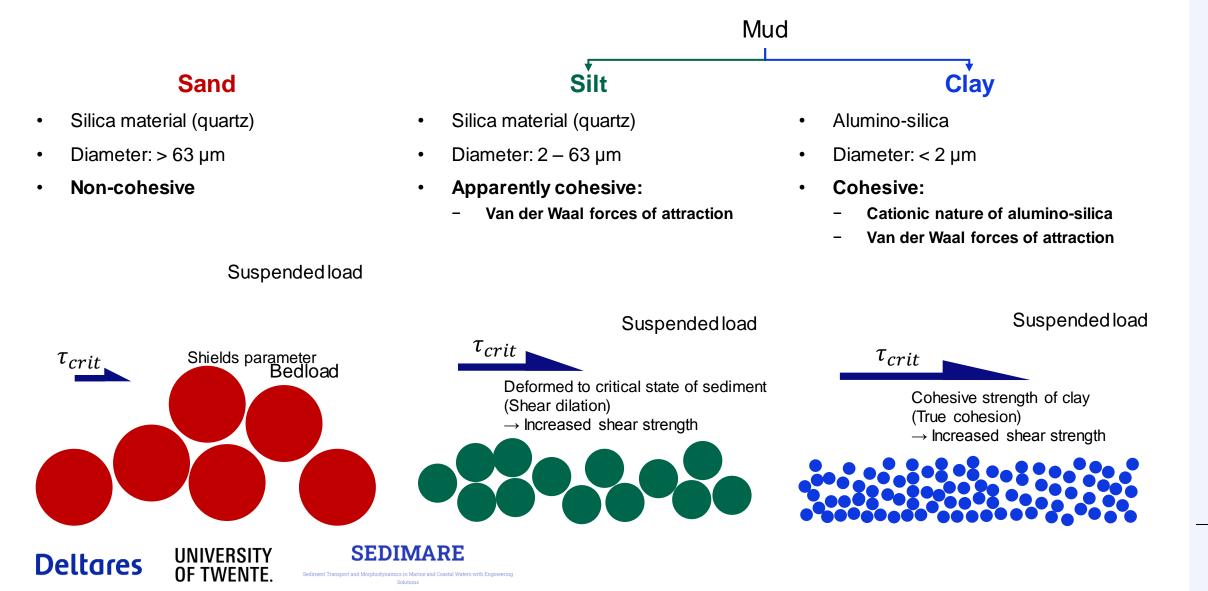
Sand-mud erosion

- Erosion of sediment bed containing both sand and mud
- Modes of erosion:
 - Particle / floc erosion
 - Surface erosion
 - Mass erosion
- Estimating erosion requires a definition for initiation of motion of both sand and mud
 - $E [kg/(m^2 s)] vs \tau_b$ (Pa)
 - Definition of τ_{crit}





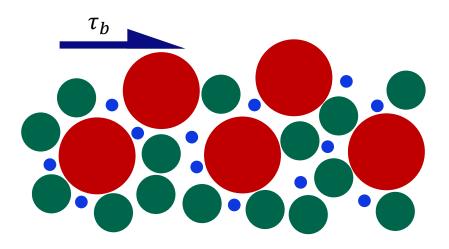
Initiation of motion for each sediment class



Initiation of motion

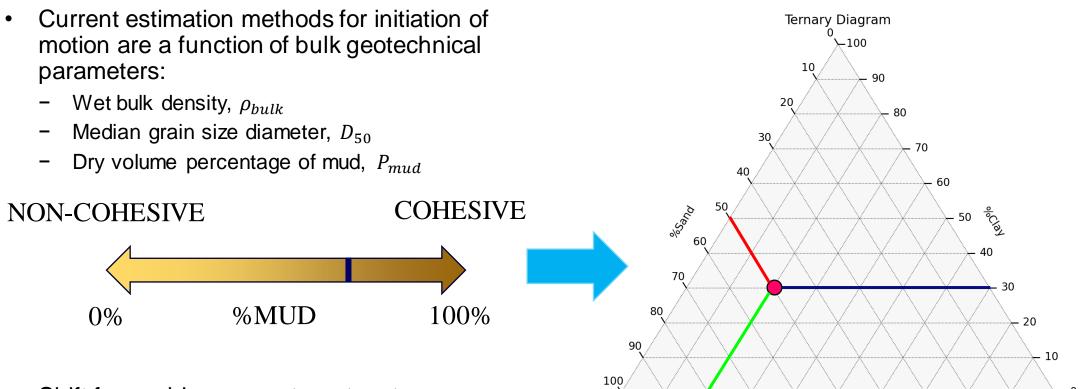
Sand-Mud (sand-silt-clay)

- What is the process for sand-mud beds?
 - Shields parameter?
 - Shear dilation?
 - True cohesion?





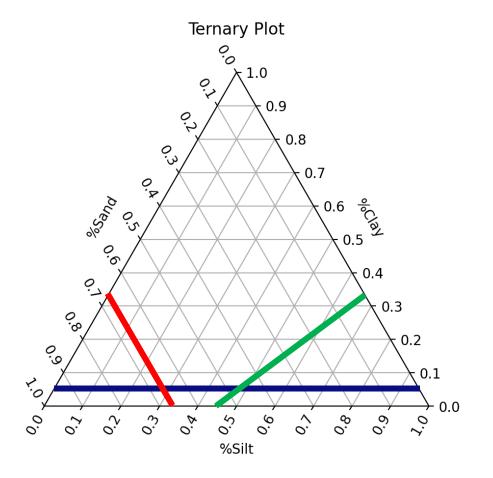




%Silt

Shift from a binary spectrum to a ternary analysis

- Van Ledden bed types
 - Based on information from ternary diagram and volume concentration of water
 - Framework characteristics:
 - 1. Cohesion
 - 2. Network structure





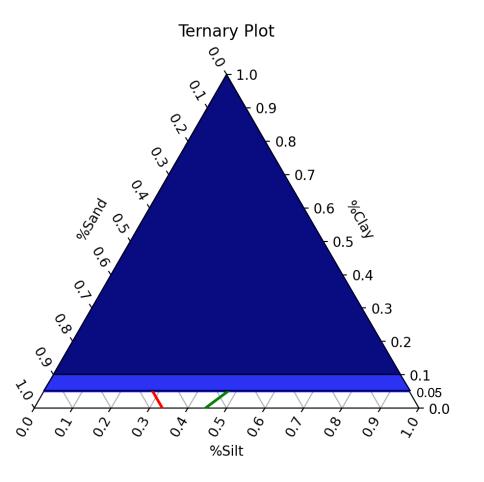
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Solution

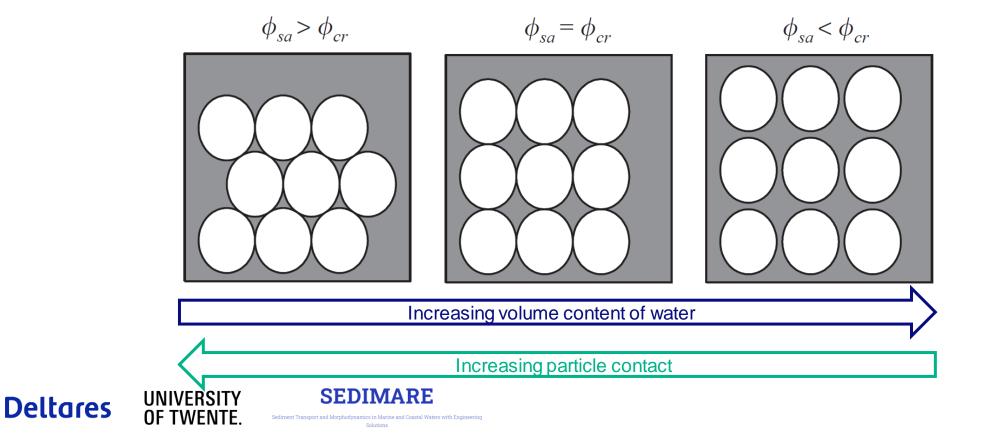
- Cohesion ٠
 - Based on dry volume percentage of clay —
 - Literature provides a lower limit of 5-10% —

 $P_{clay} = 5 - 10\%$





- Network structure
 - Interaction of particles based on the volume content of water, ϕ_{water}
 - Illustration using pure sand:

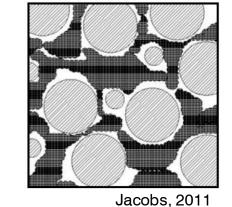


Van Ledden et al., 2003

- Network structure
 - Van Ledden et al., 2003 suggests that for sandmud, the dominant sediment can be based on wet volume percentage of sand, ϕ_{sand} , and silt, ϕ_{silt} .
 - Sand-dominated network structure: $\phi_{sand} \ge 40 50\%$
 - Silt-dominated network structure:

$$\frac{\phi_{silt}}{(1-\phi_{sand})} \ge 40 - 50\%$$

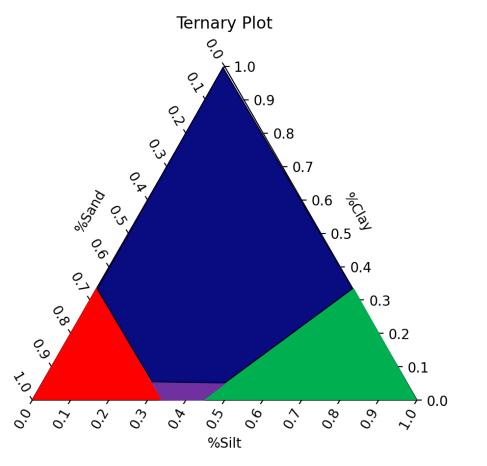
- Other network structures:
 - 1. Clay-water matrix
 - 2. Mixed structures



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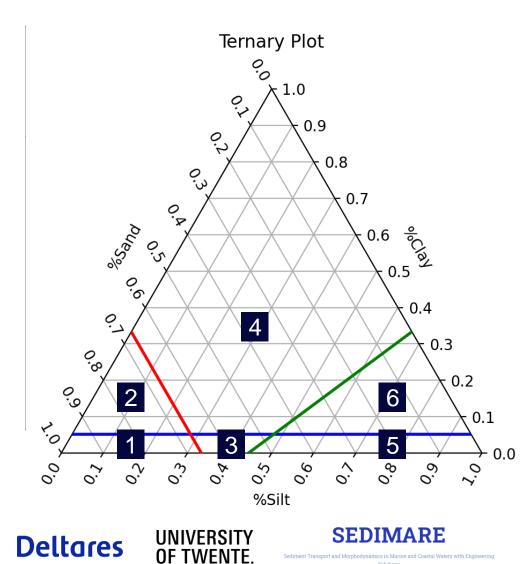
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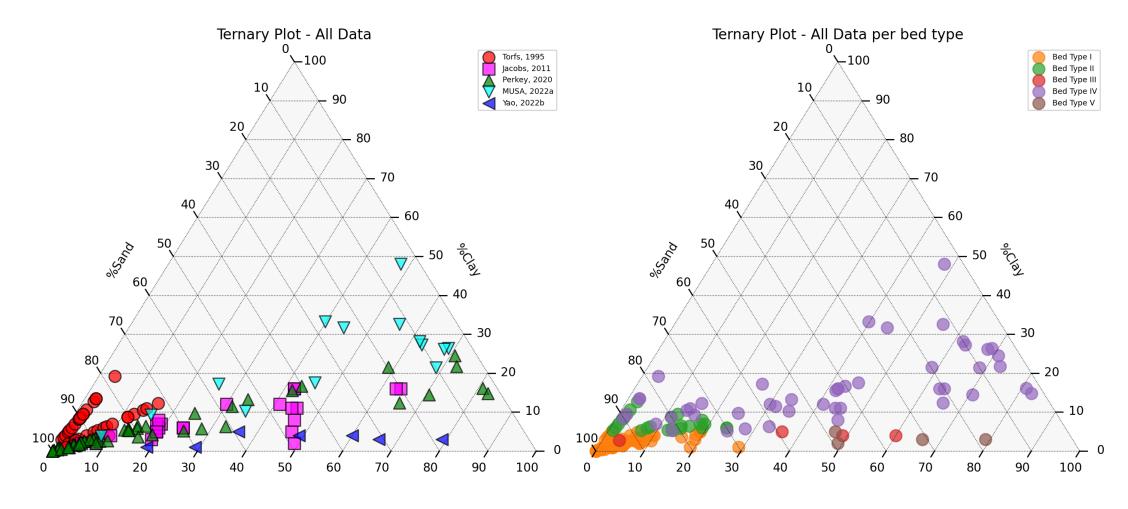
10

Solutions



Cohesion	
on-cohesive	Cohesive
$_{y} < 5 - 10\%$	$\left(P_{clay} \geq 5 - 10\%\right)$
Type 2:	
ninated, non- 1. Sand-dominated sediment bed 2. Sand-dominated	Sand-dominated, cohesive non-conesive sediment bed conesive
3. Mixed structure	n Type faitesive
natecClagrdominated	Soltestweinated, cohesive
sedisticted or enclosed r	
6. Silt-dominated of	ohesive: Type 4:
diment, non-	Clay-water matrix, cohesive
sediment bed	sediment bed

Application of framework to collected data





%Silt

Next steps

- Dominant initiation of motion process for each bed type
- Existing initiation of motion equations, τ_{crit}
 - Which equation best describes each Van Ledden bed type?
 - 1. Mehta, 1988
 - 2. Mitchener & Torfs, 1996
 - 3. Roberts et al., 1998
 - 4. Van Ledden et al., 2003
 - 5. Van Rijn, 2007
 - 6. Wu et al., 2018
 - 7. Yao et al., 2022



Initiation of motion of sand-mud beds

Paterno "Jowi" S. Miranda IV, MCE

Research update



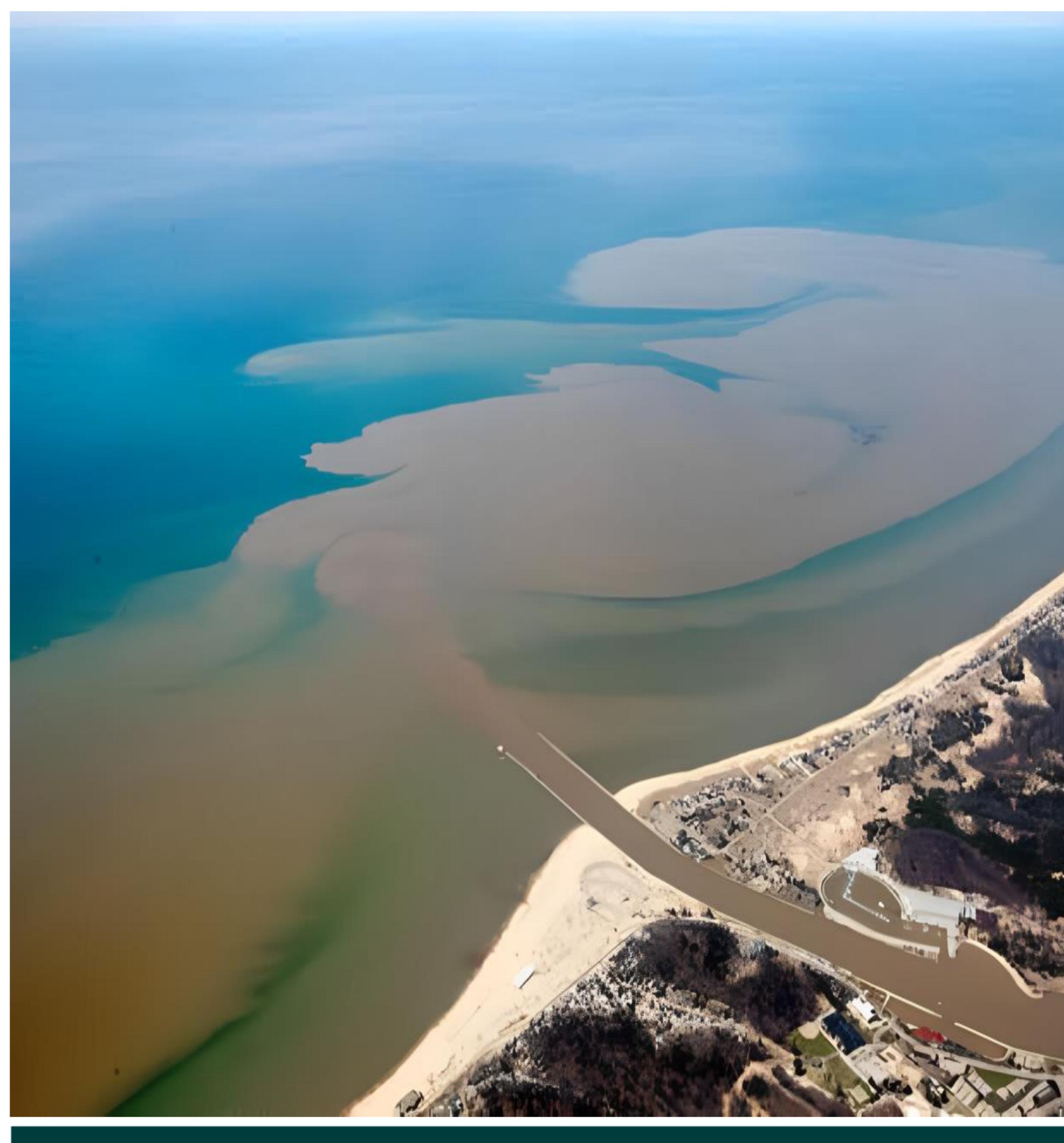
in patrno4







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DC 4: Mixing and transport in the coastal area

Subervsors. Prof. Maurizio Brocchini – UNIVPM Prof. Athanassios Dimas – UPATRAS Prof. Matteo Postacchini – UNIVPM

Nasm Soon





The nearshore area:

D Estuarine environments are dynamically complex, with hydrodynamics being triggered by many factors, e.g., nonlinear interactions between the bathymetry, the river current and many sea forcing actions.

> In our project, we are focusing on the **development of a 2D numerical solver** for the hydromorphodynamics in the nearshore area to extract the Lagrangian motion/flow field. ✓ To inspect the nearshore circulations, three current values were tested, following the study of Olabarrieta et al. (2014).

North Carolina

PhD Project- DC4: Mixing and transport in the coastal area

dynamically evolving

often densely populated

under increasing threat

Olabarrieta et al. (2014): circulation of the New River inlet,





Domain: 3 km long in both cross-shore and alongshore directions

- Boundary conditions:
- **Right** Boundary: Riemann
- Left Boundary: Riemann
- **Top** Boundary: Total Discharge
- Time Simulation: 5 hrs



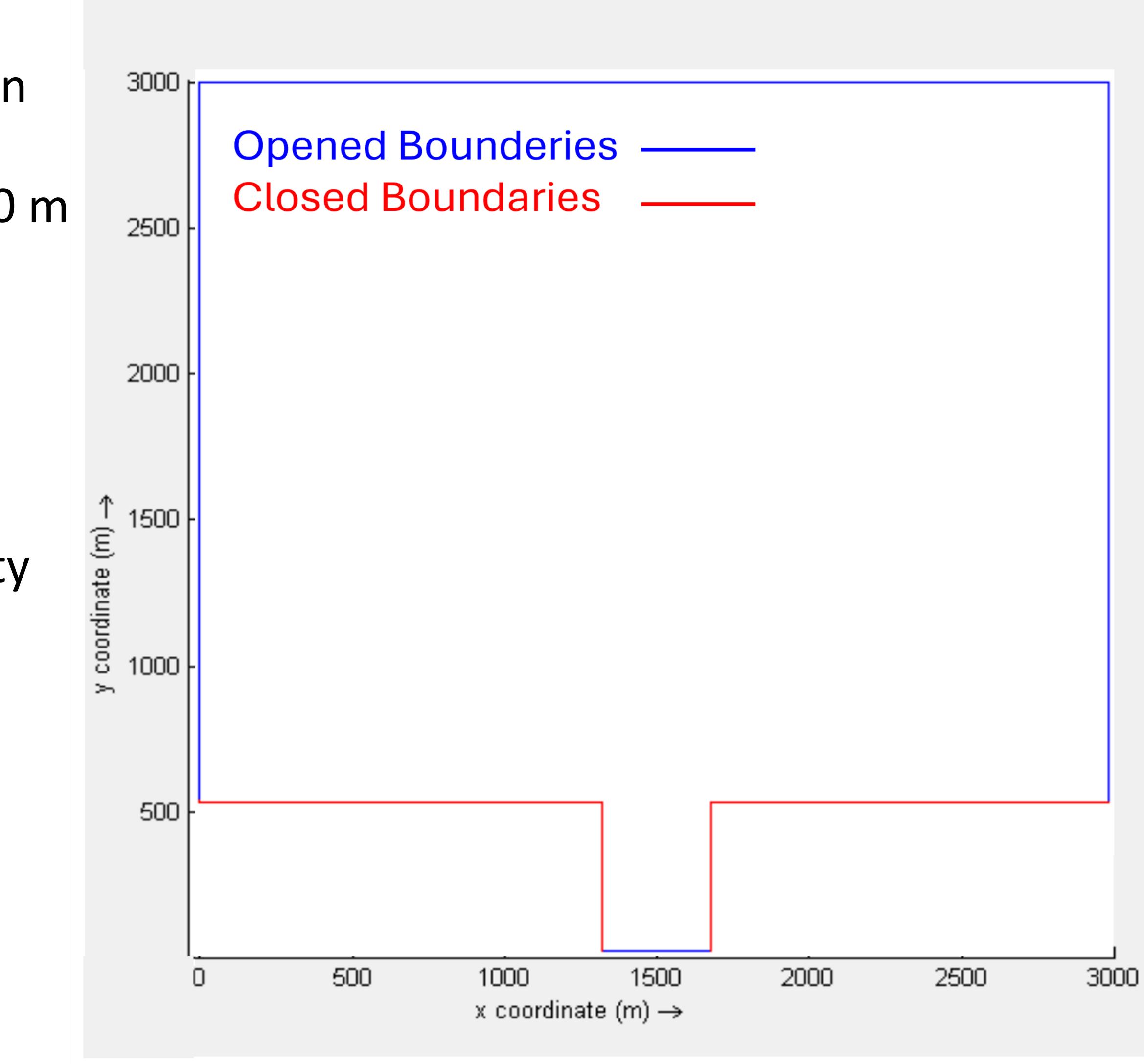
Delft-3D-Flow, to investigate the flow circulation

• Computational mesh: spacings $\Delta x = 10 \text{ m}$, $\Delta y = 20 \text{ m}$

The resolution of the grid: 151 x 300 (M x N)

- Inlet Boundary, Estuarine Window: Current velocity

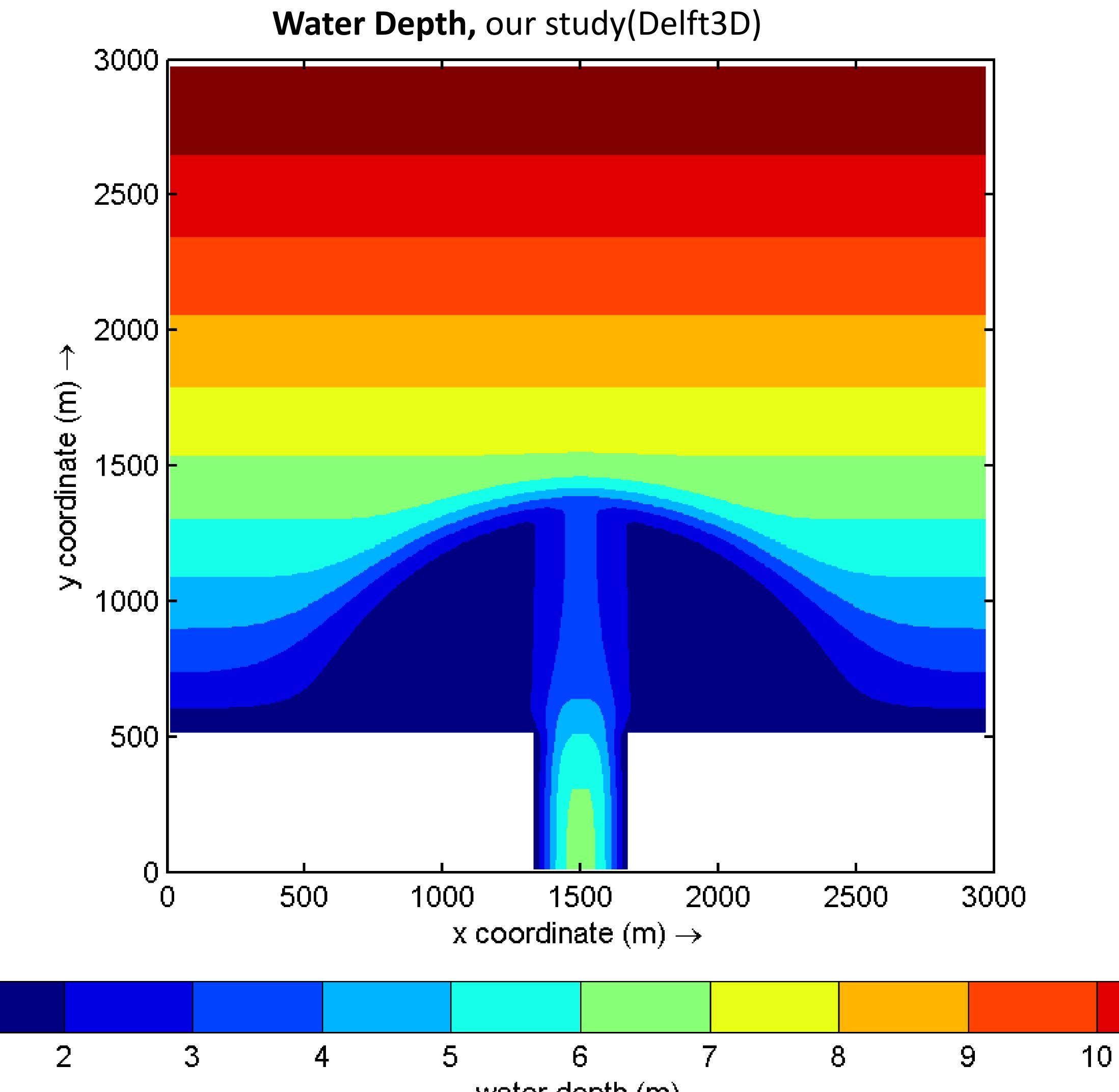
Inrtoduction



DC4: Mixing and transport in the coastal area

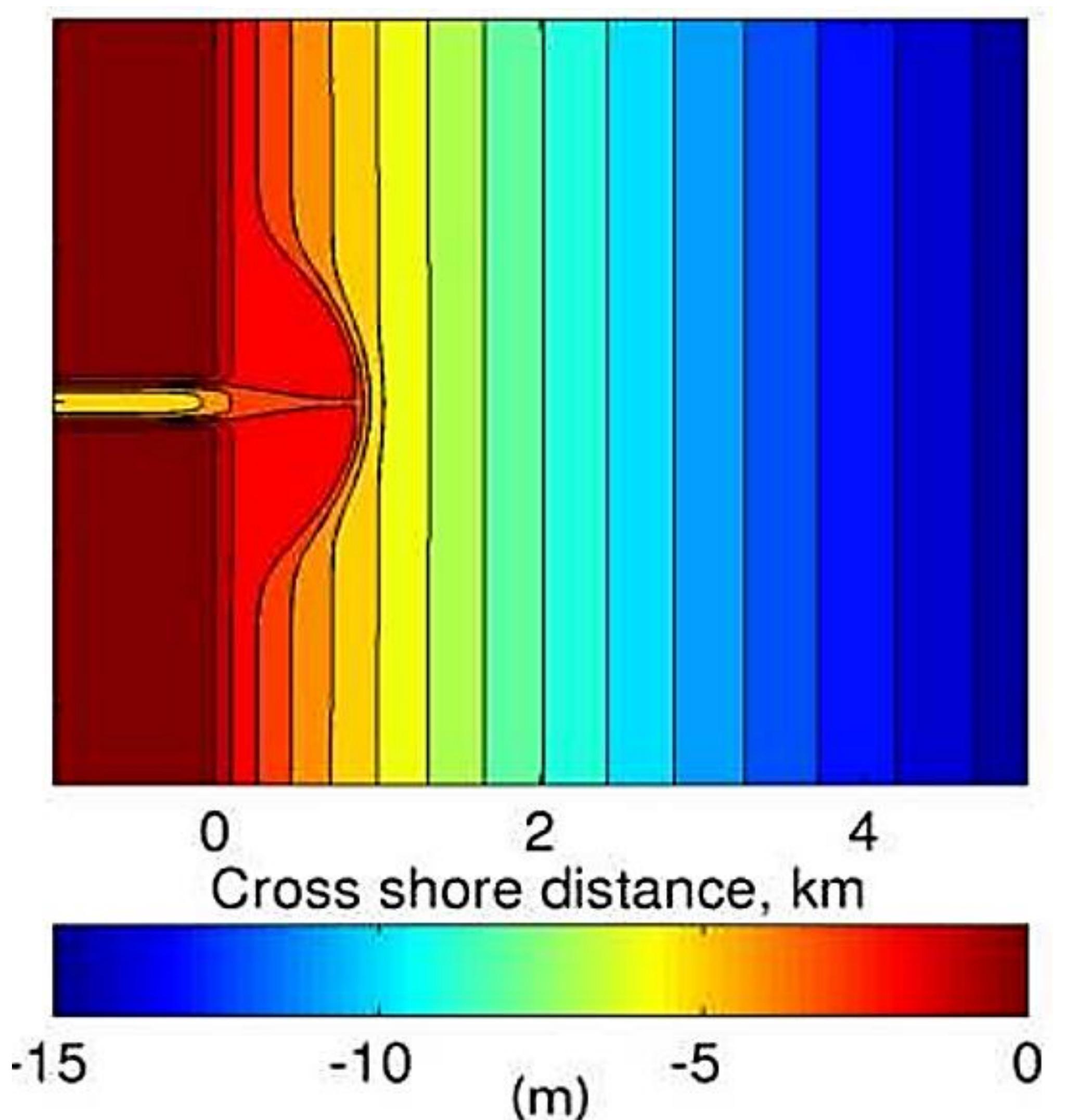


Water Depth





water depth (m)







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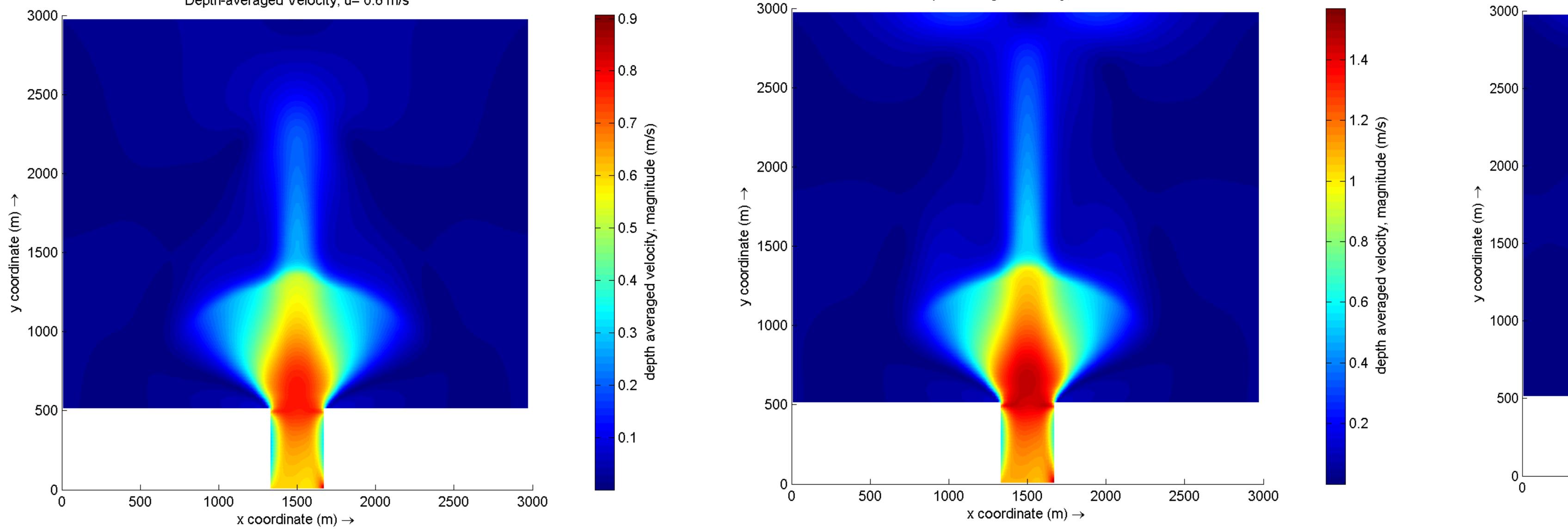
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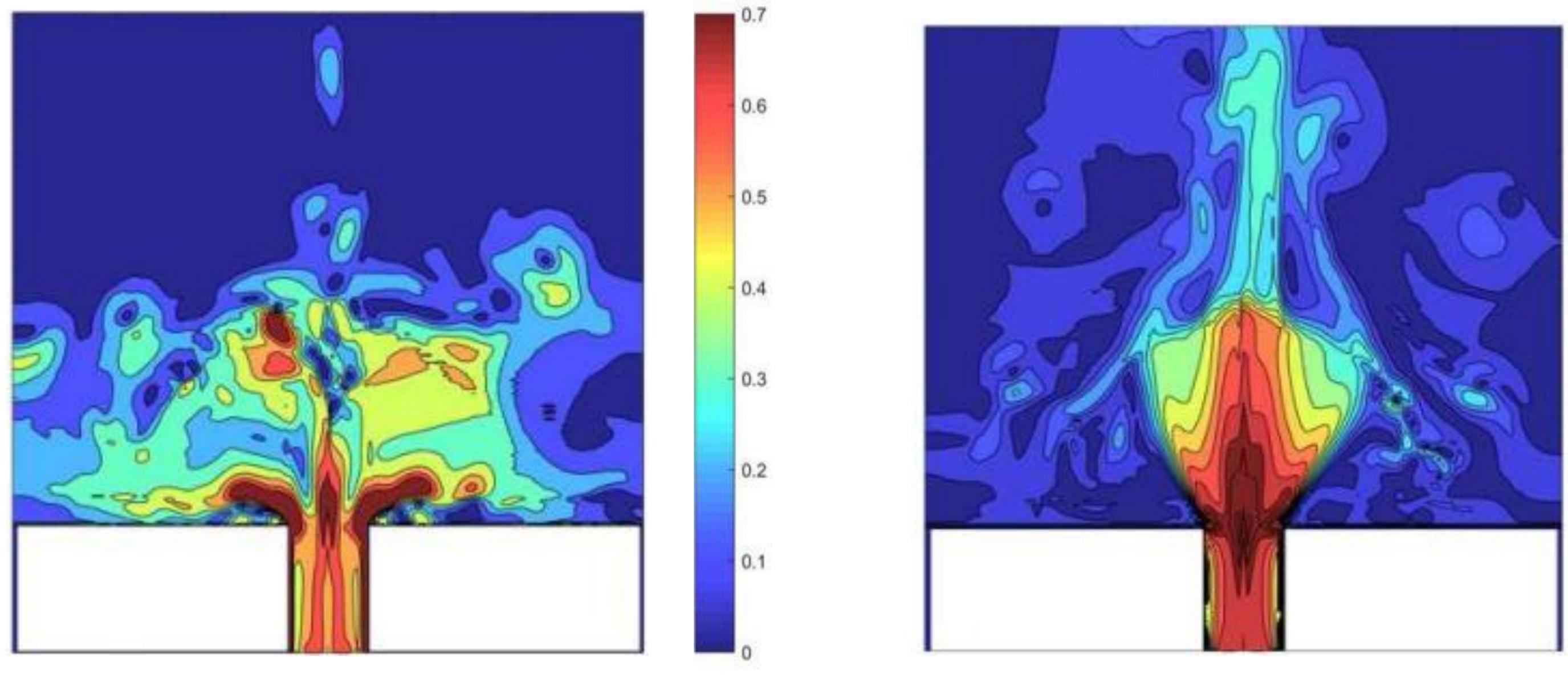
Depth-averaged Velocity- Delft-3D: our study

U=0.6 m/s

Depth-averaged Velocity, u= 0.6 m/s



U=0.6 m/s



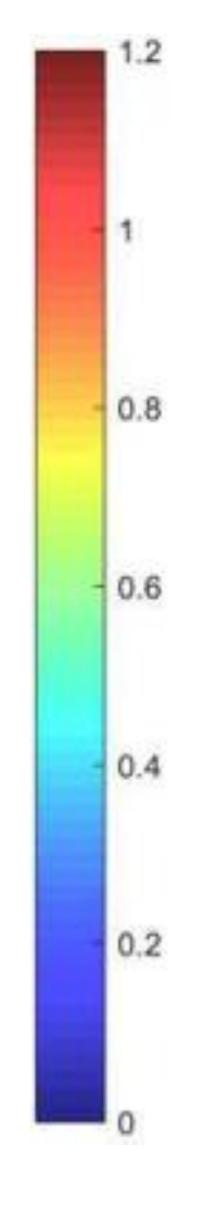


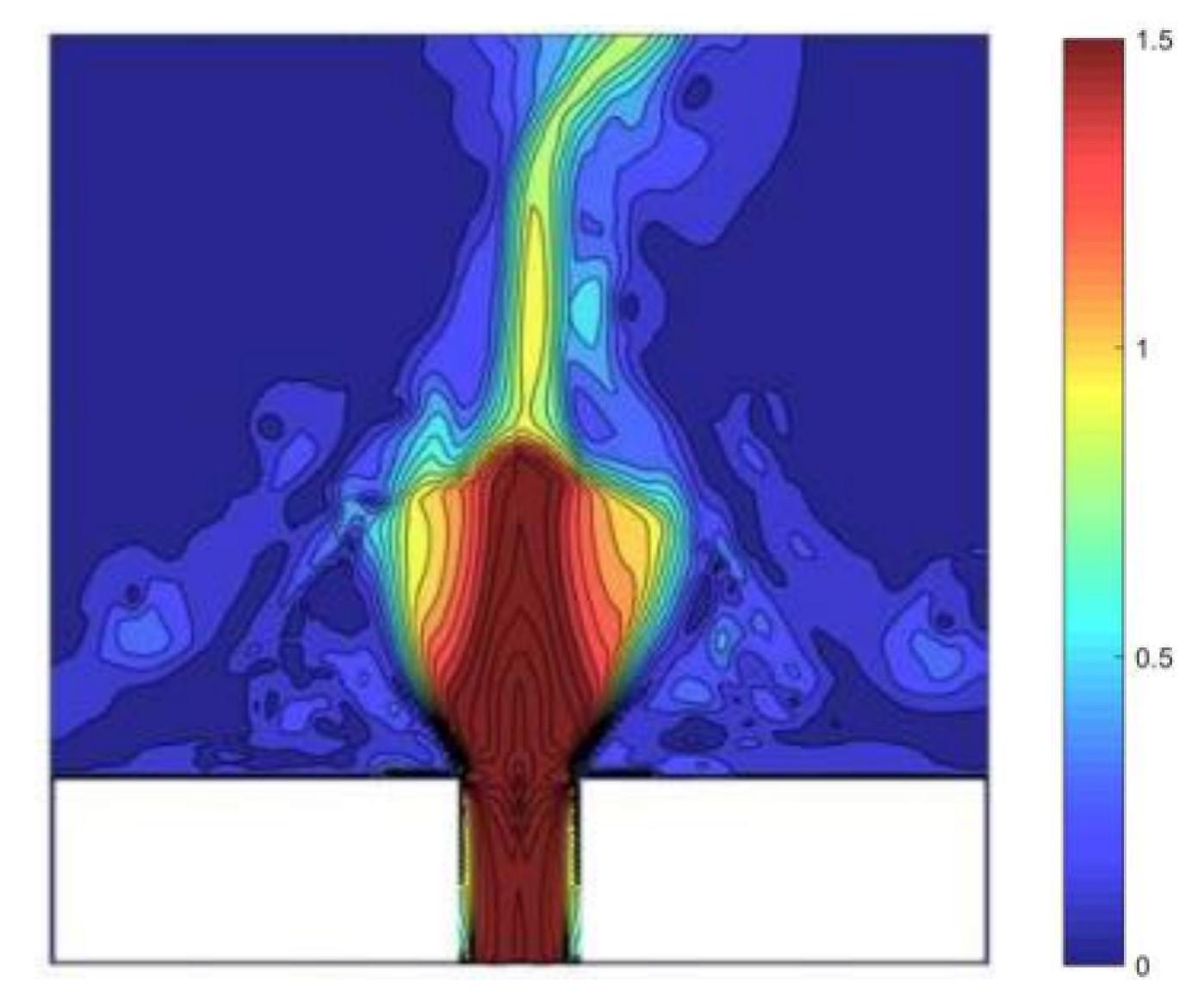
U=1.1 m/s

Depth-averaged Velocity, u= 1.1 m/s

Mean Velocity- NSWE Solver: Melito et al. (2019)

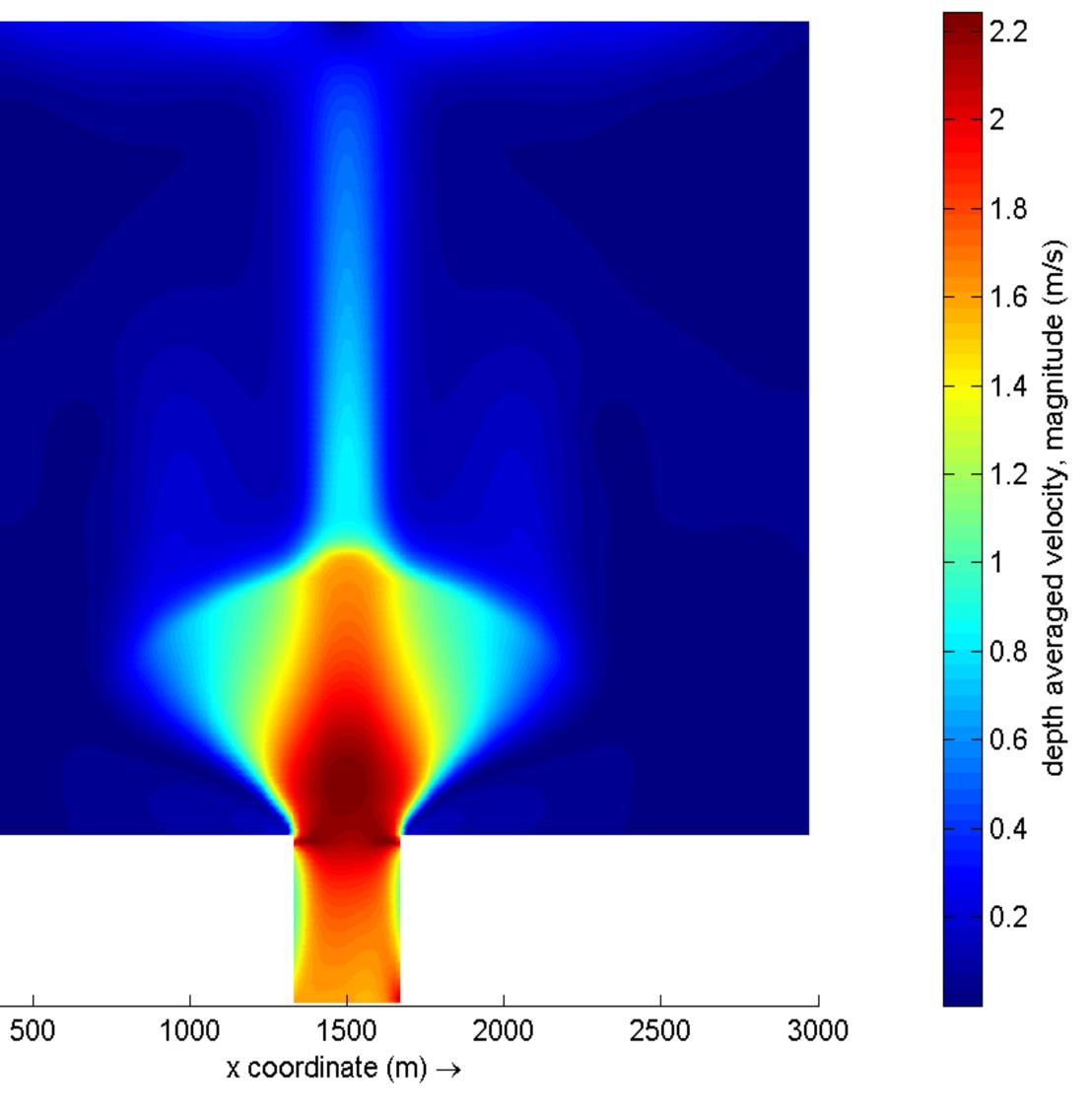
U=1.1 m/s





U=1.6 m/s

Depth-averaged Velocity, u= 1.6 m/s



U=1.6 m/s

DC4: Mixing and transport in the coastal area



D Particle Tracking:

Changes in sediment transport — coastal erosion, adverse effects on habitats, infrastructure, and communities located along the coast.

\checkmark

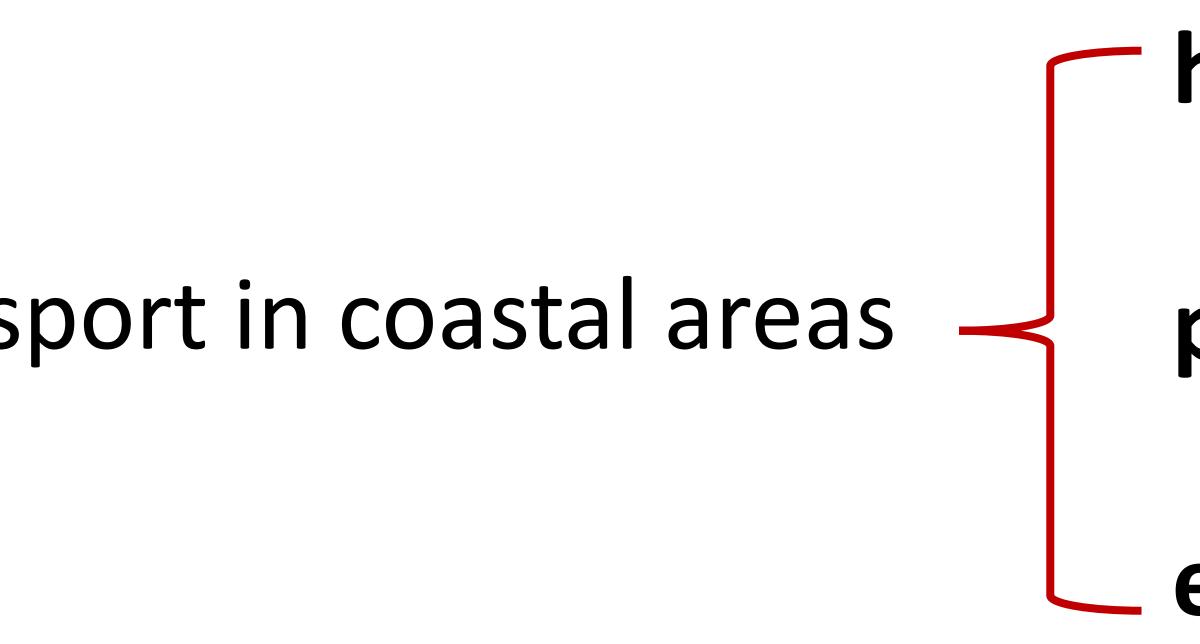
Horizontal dispersion due to turbulence.

Numerical part: Delft-3D

 \succ To effectively understand and model particle transport in coastal areas -

D-WAQ PART is a random walk particle tracking model

Pollution transport *b* detrimental effects on marine life and habitats.



DC4: Mixing and transport in the coastal area

hydrodynamic forces

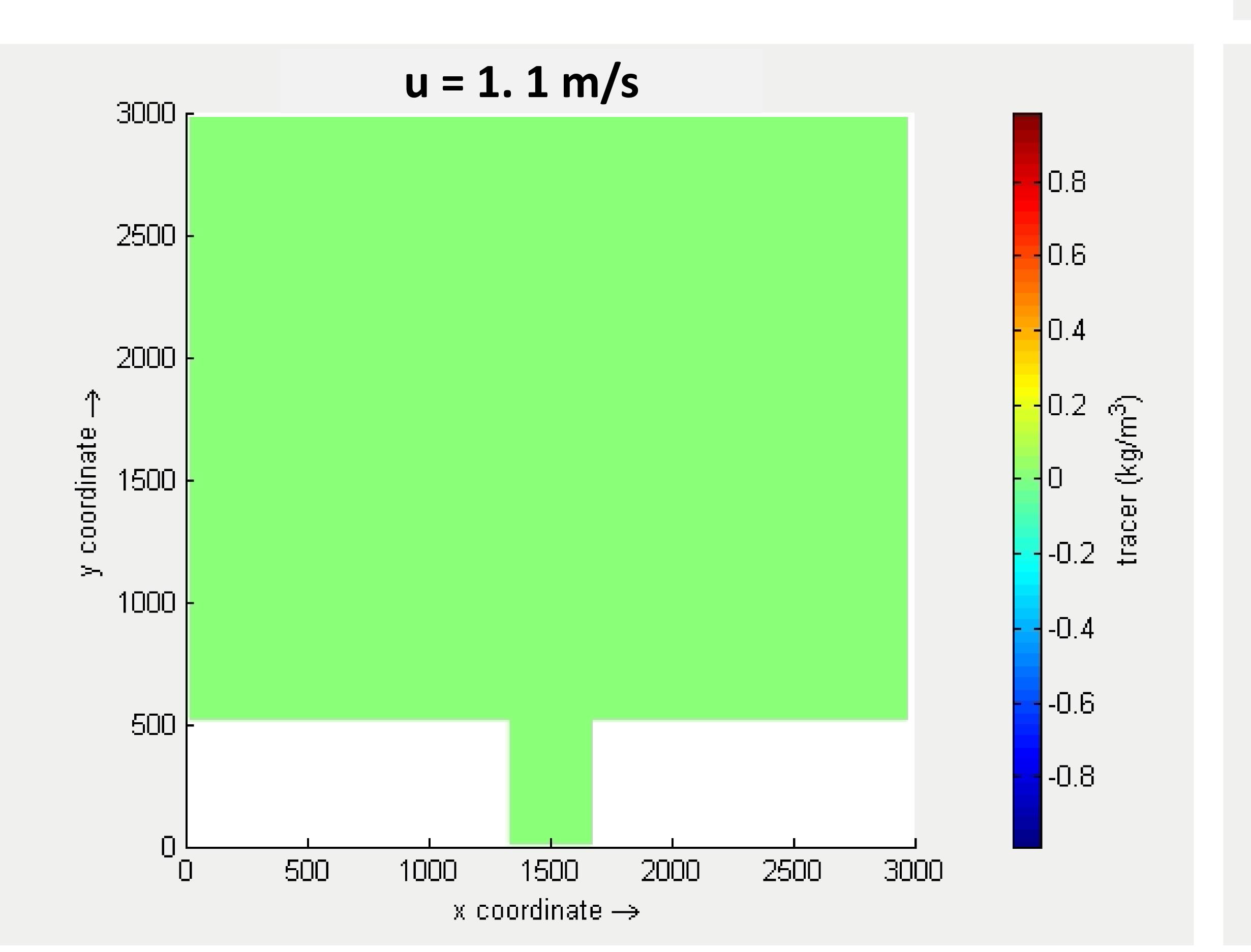
particle characteristics

environmental conditions

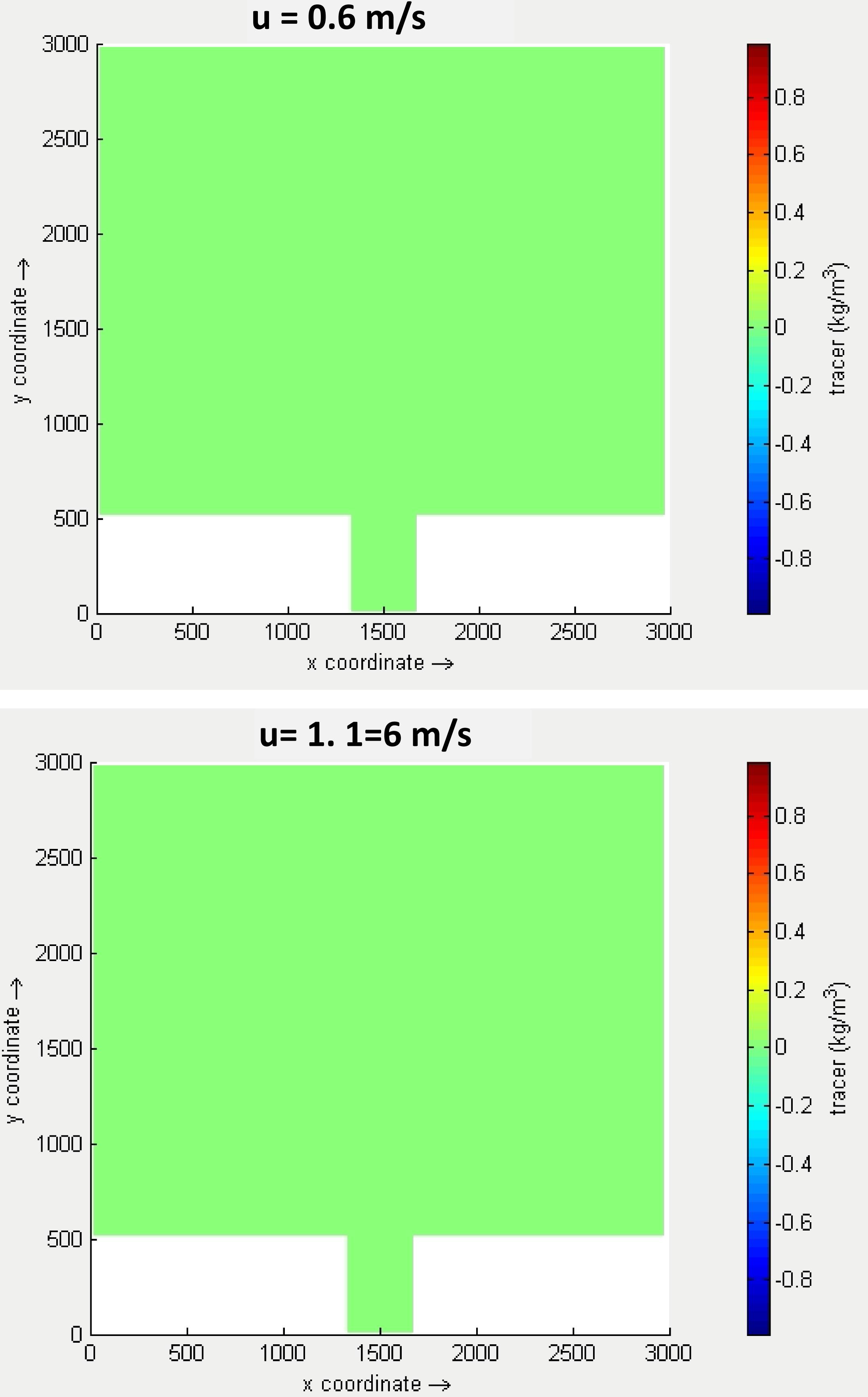


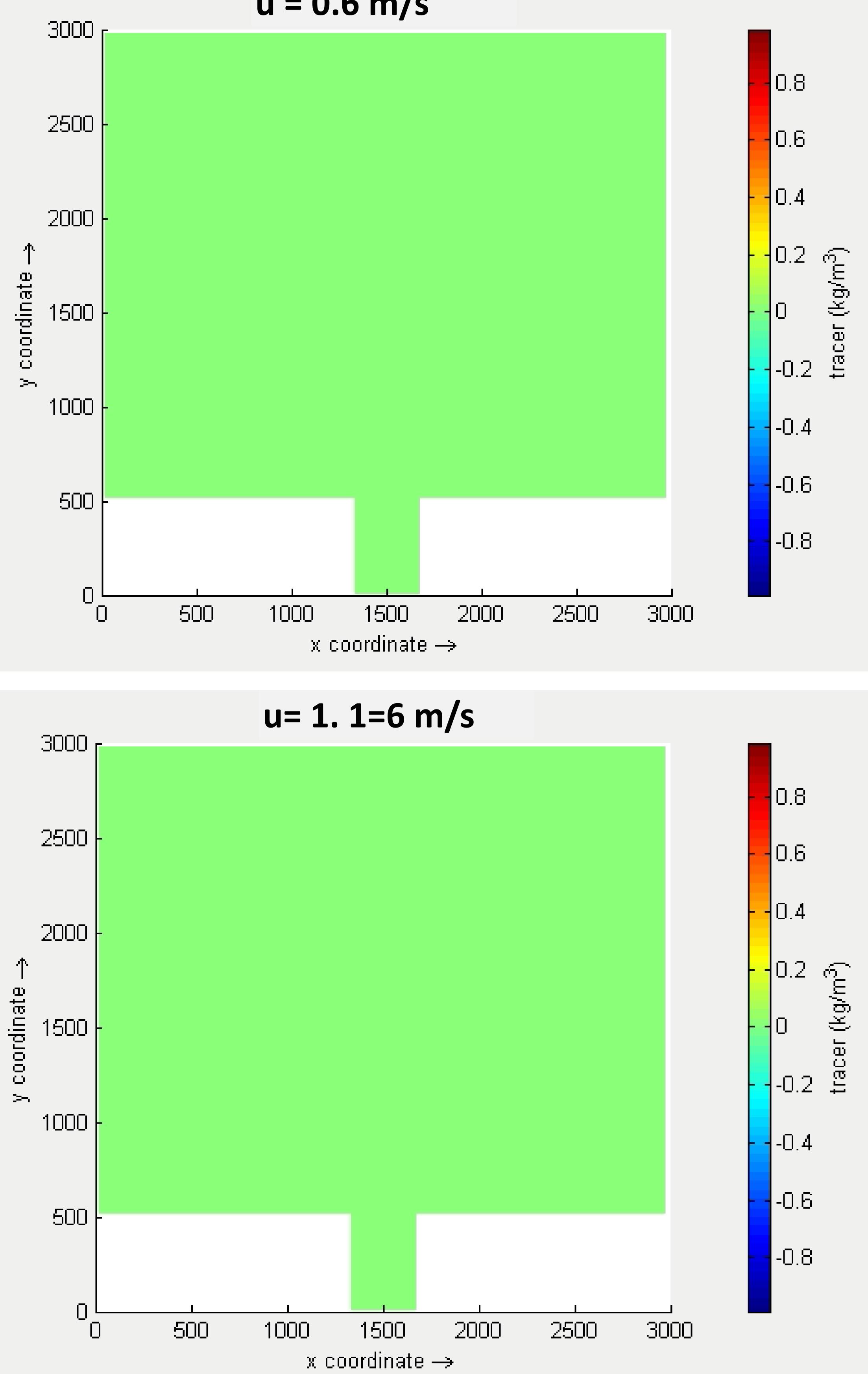
Releasing particles from Estuarine part:

- The advection step due to the shear stresses from currents
- The related to the horizontal random walk step is
 - dispersion.
- Horizontal dispersion due to turbulence., According to
 - turbulence theory this dispersion increases in time.



Numerical part: Delft-3D

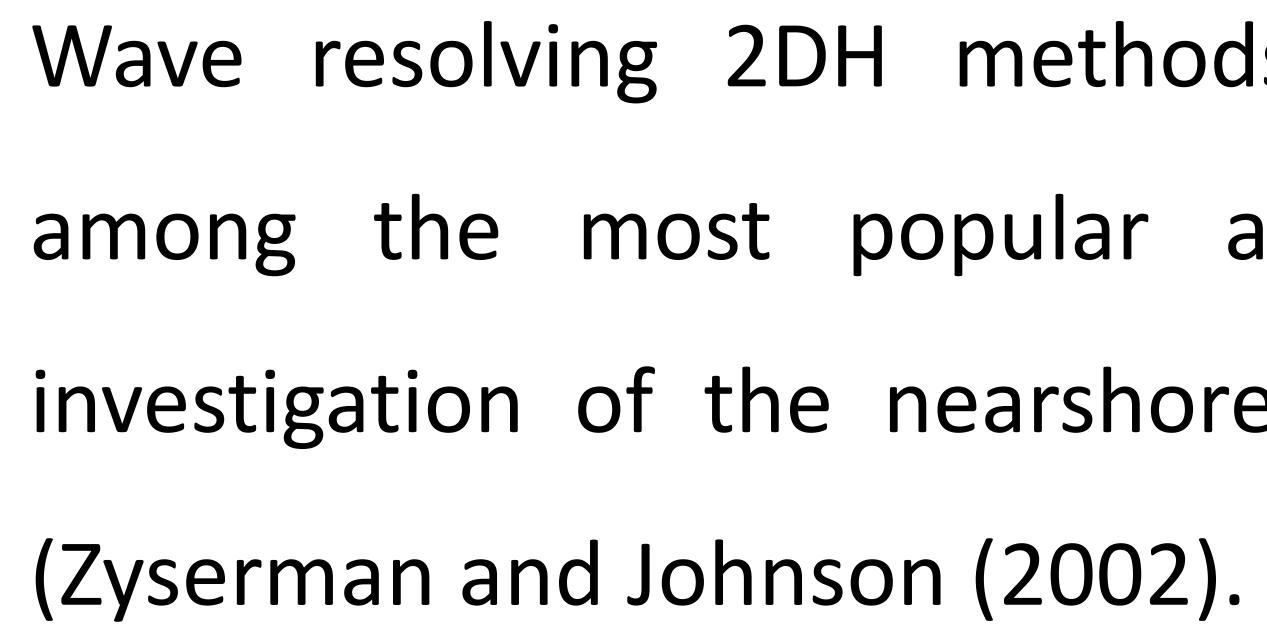


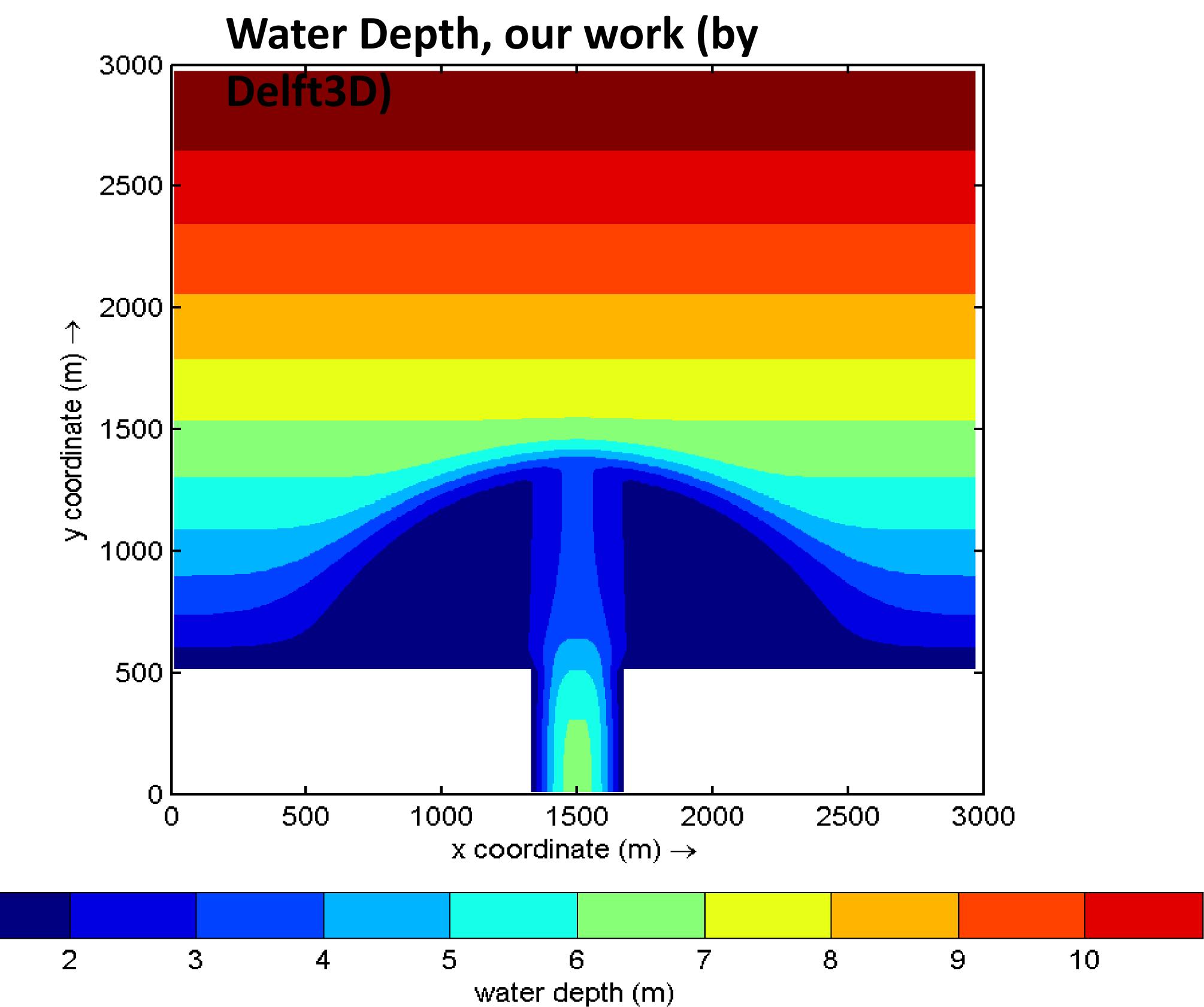




 \checkmark

O To investigate the dynamics of the nearshore zone:



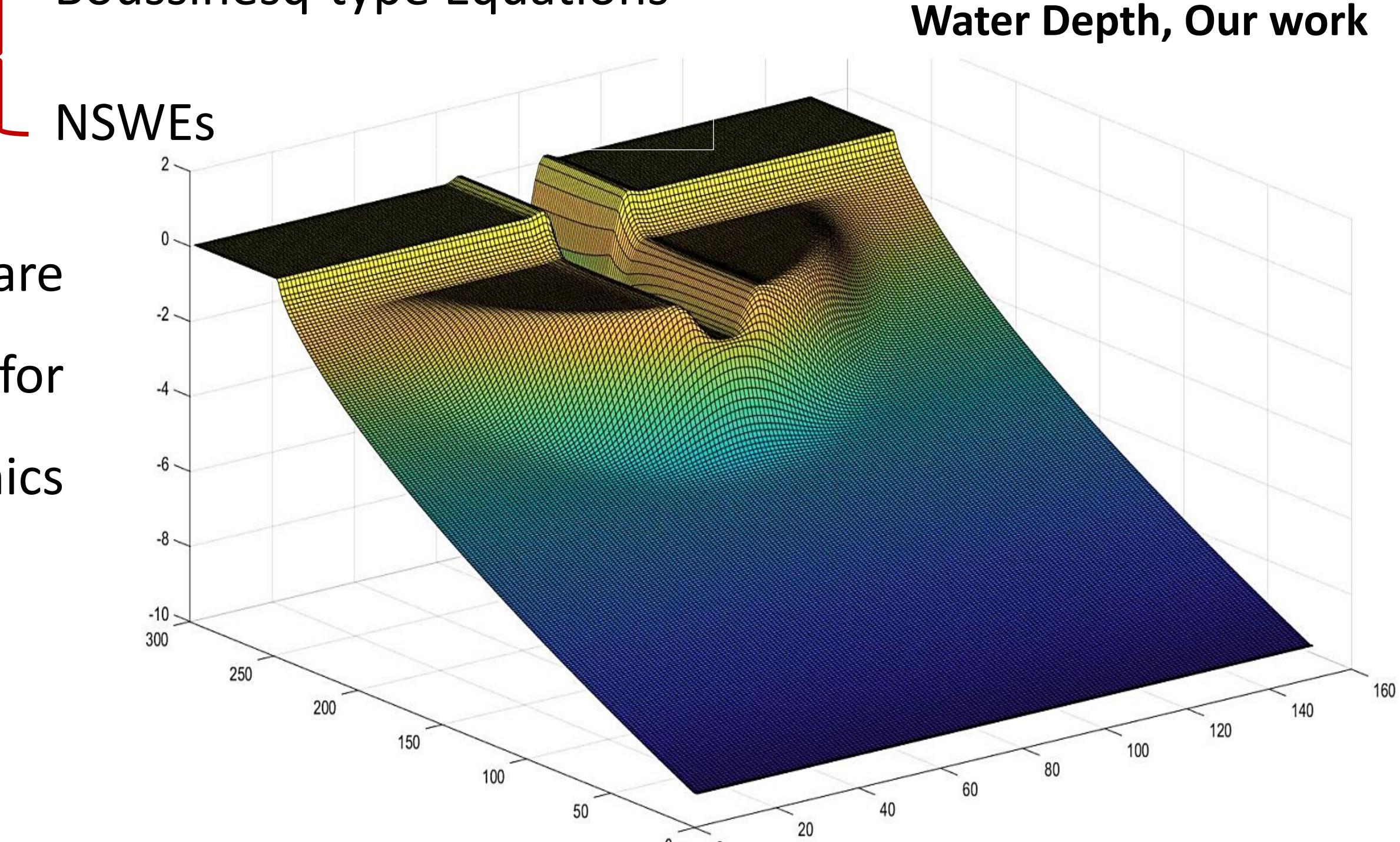




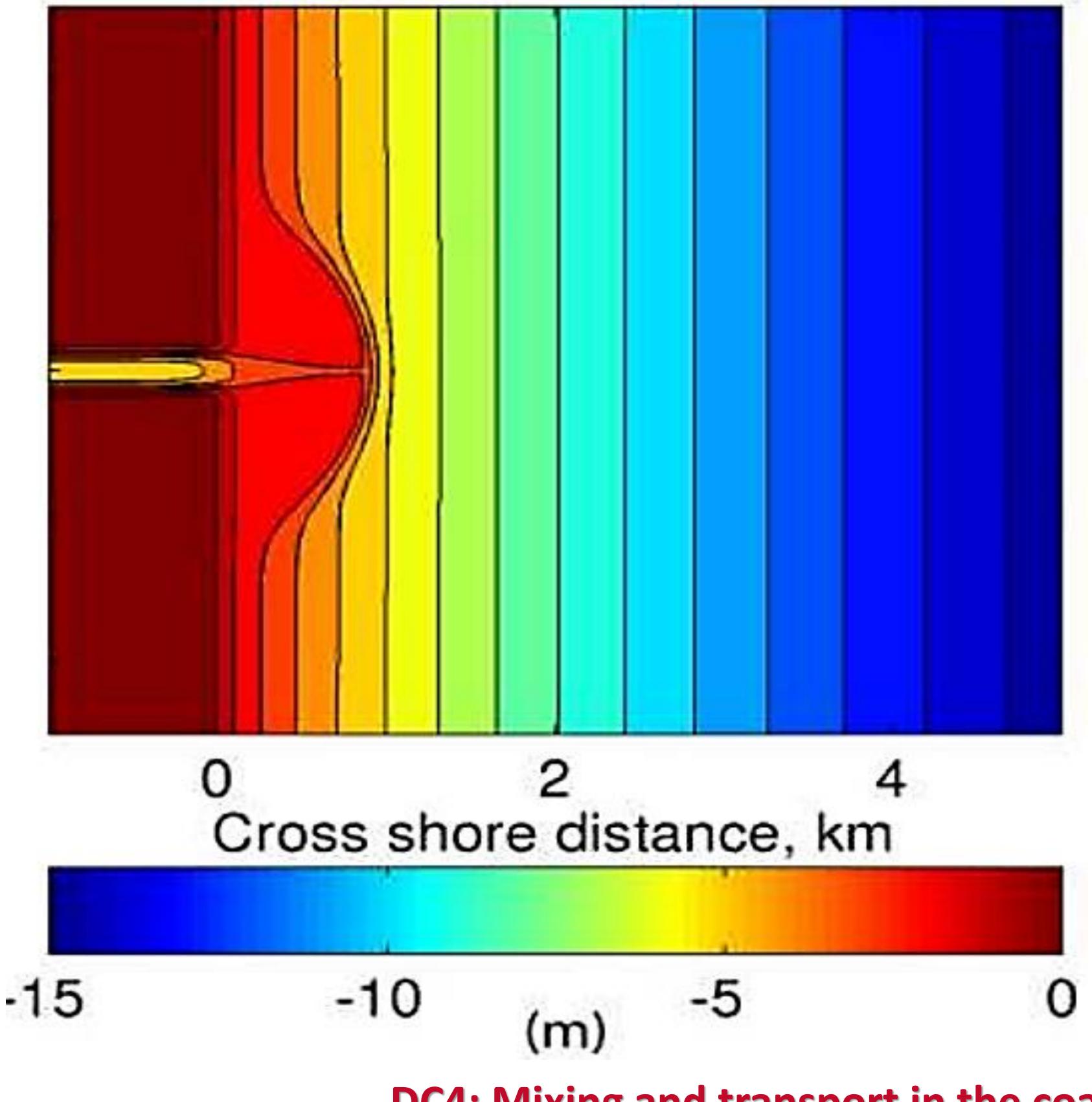
Wave resolving 2DH methods, such as NSWEs, are the most popular approximate models for investigation of the nearshore hydro-morphodynamics

NSWE Solver

Boussinesq-type Equations



Water Depth, Olabarrieta et al. (2014)



DC4: Mixing and transport in the coastal area







- \checkmark How time stepping is applied:
- **Discretization of Equations:**
- respectively.
- 2.
- **Euler Time Stepping:** 3. by the time step dt.

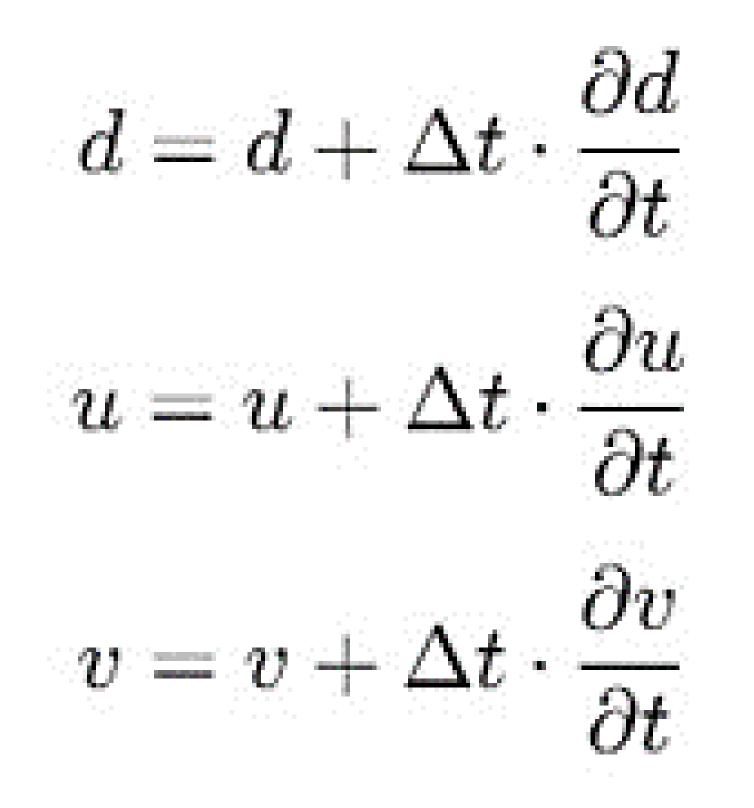
<u>central differencing for spatial derivatives.</u>

The depth d is updated using the continuity equation.

Central Difference for Spatial Derivatives:

The code uses the central difference scheme to compute spatial derivatives:

 \checkmark The Euler method is used to advance the solution in time: The depth d and velocities u, v are updated at each time step by adding their time derivatives scaled



The new values of d, u, and v are then used in the next time step.



U The code **solves** a system of equations for a **2D shallow water model** using finite differences and

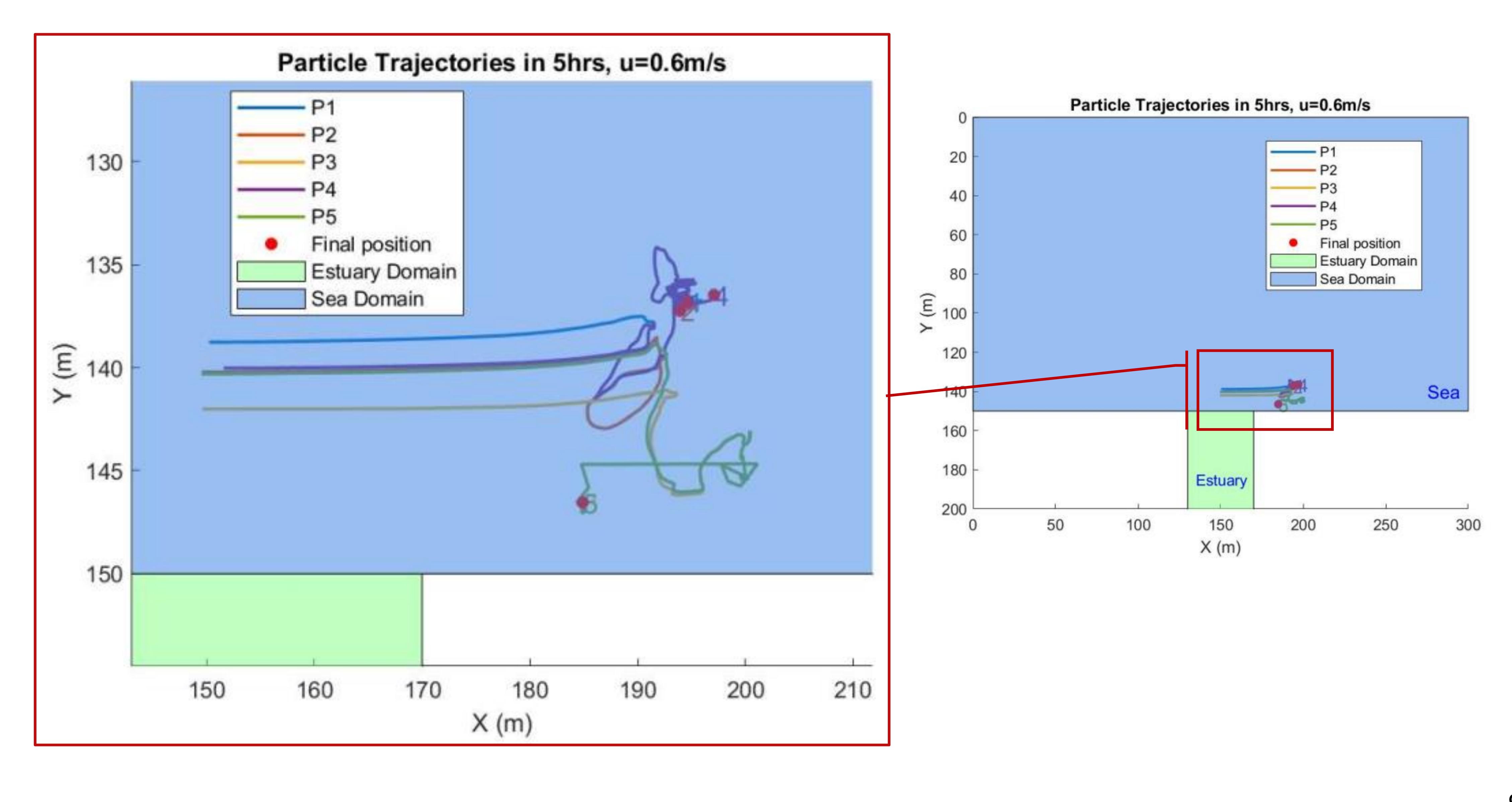
The velocities u and v are updated based on momentum equations in the x and y directions,

DC4: Mixing and transport in the coastal area



D Particle Tracking:

\checkmark Particles are tracked by interpolating the velocities (u_p and v_p) at the particle positions. Y The positions of the particles are updated based on the velocities using a forward Euler step.

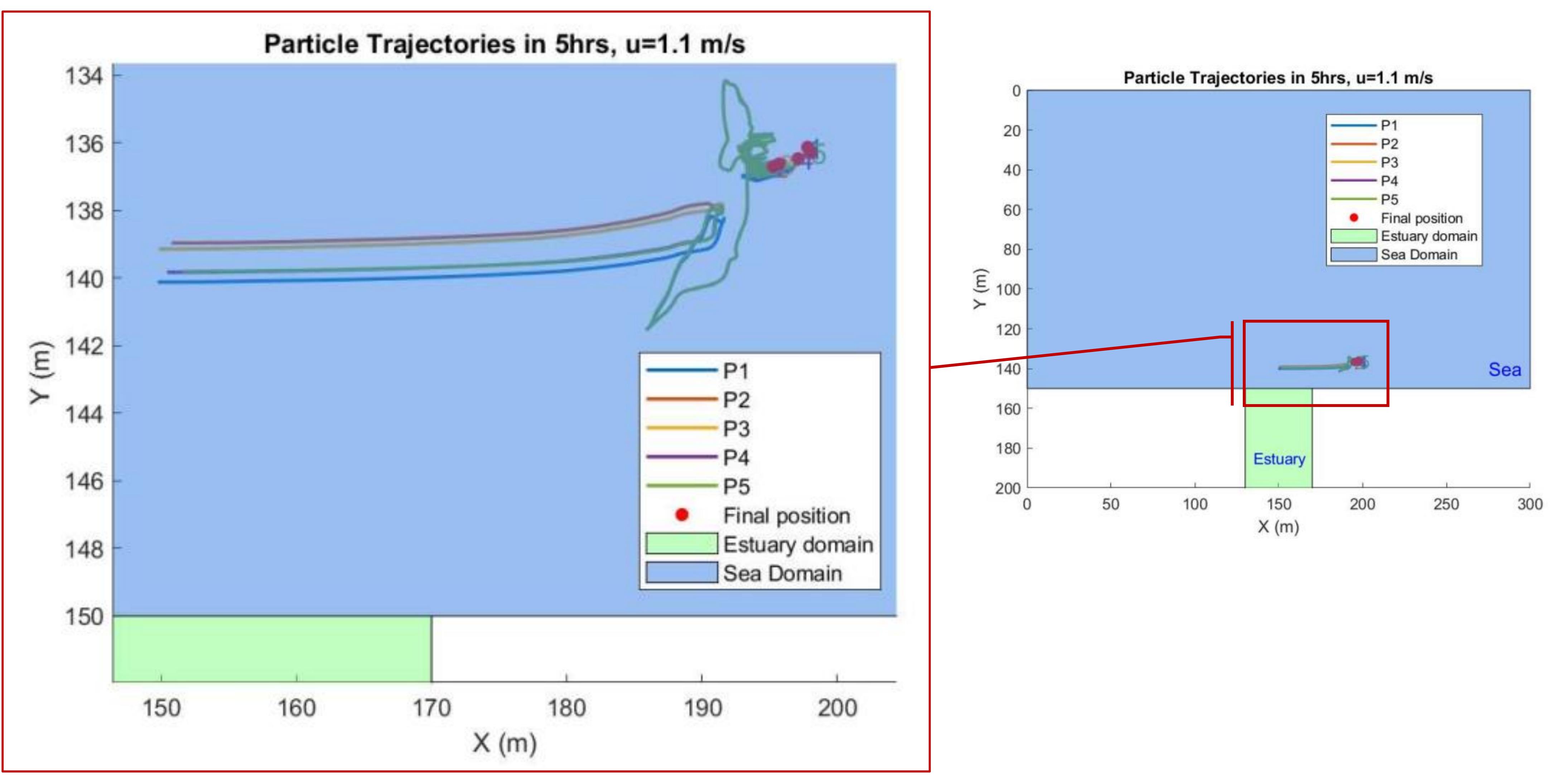


Numerical part: MATLAB Script

DC4: Mixing and transport in the coastal area



Particle Tracking by MATLAB Script:



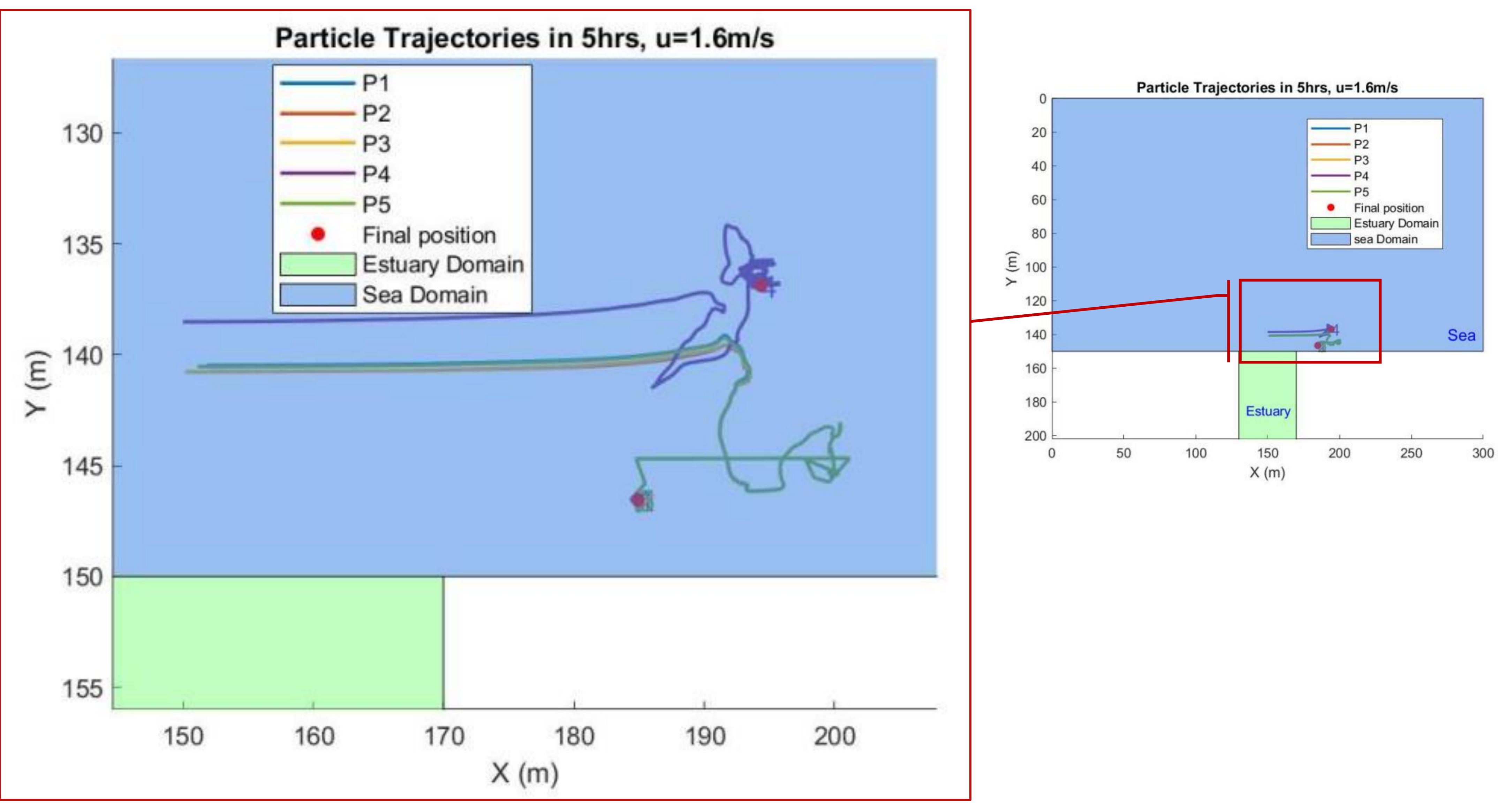


DC4: Mixing and transport in the coastal area



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Particle Tracking by MATLAB Script:





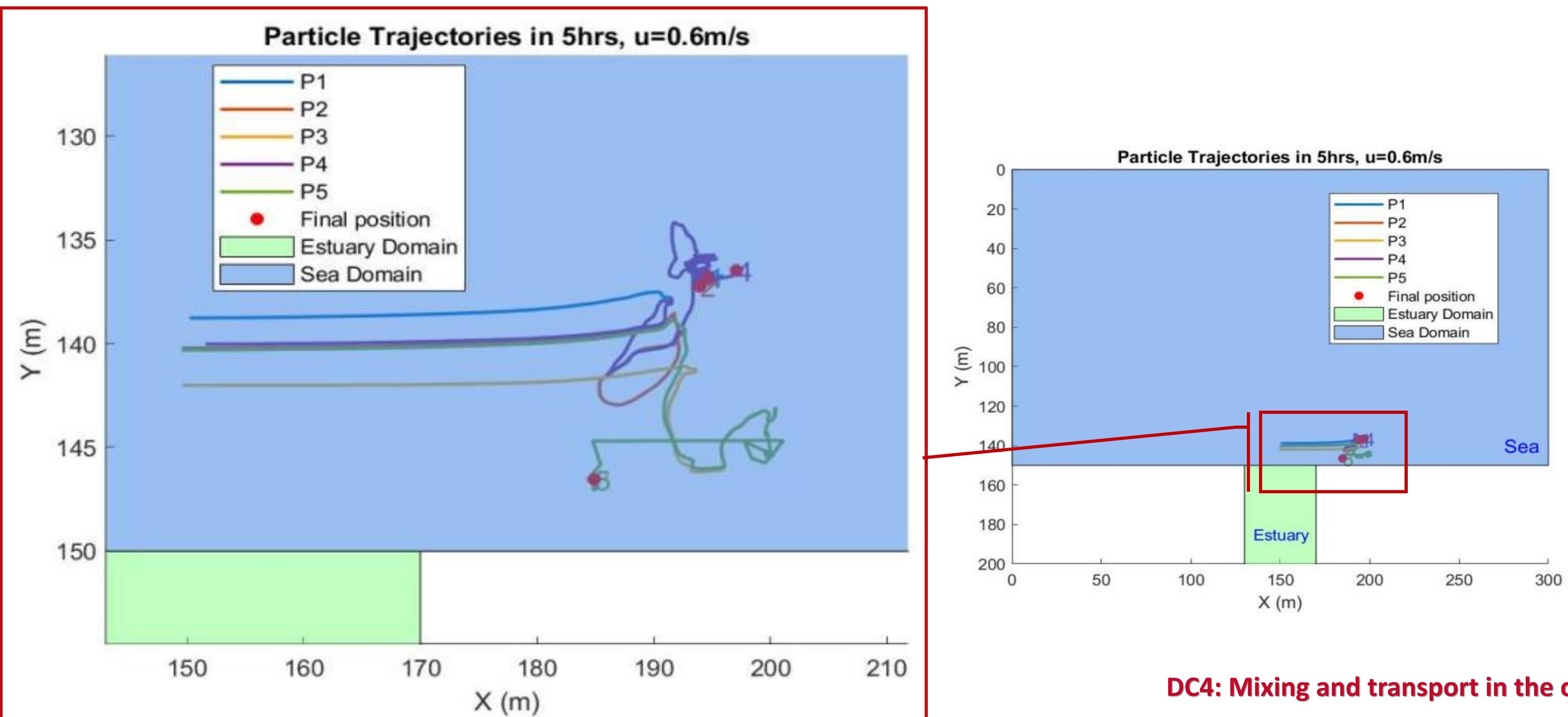
DC4: Mixing and transport in the coastal area

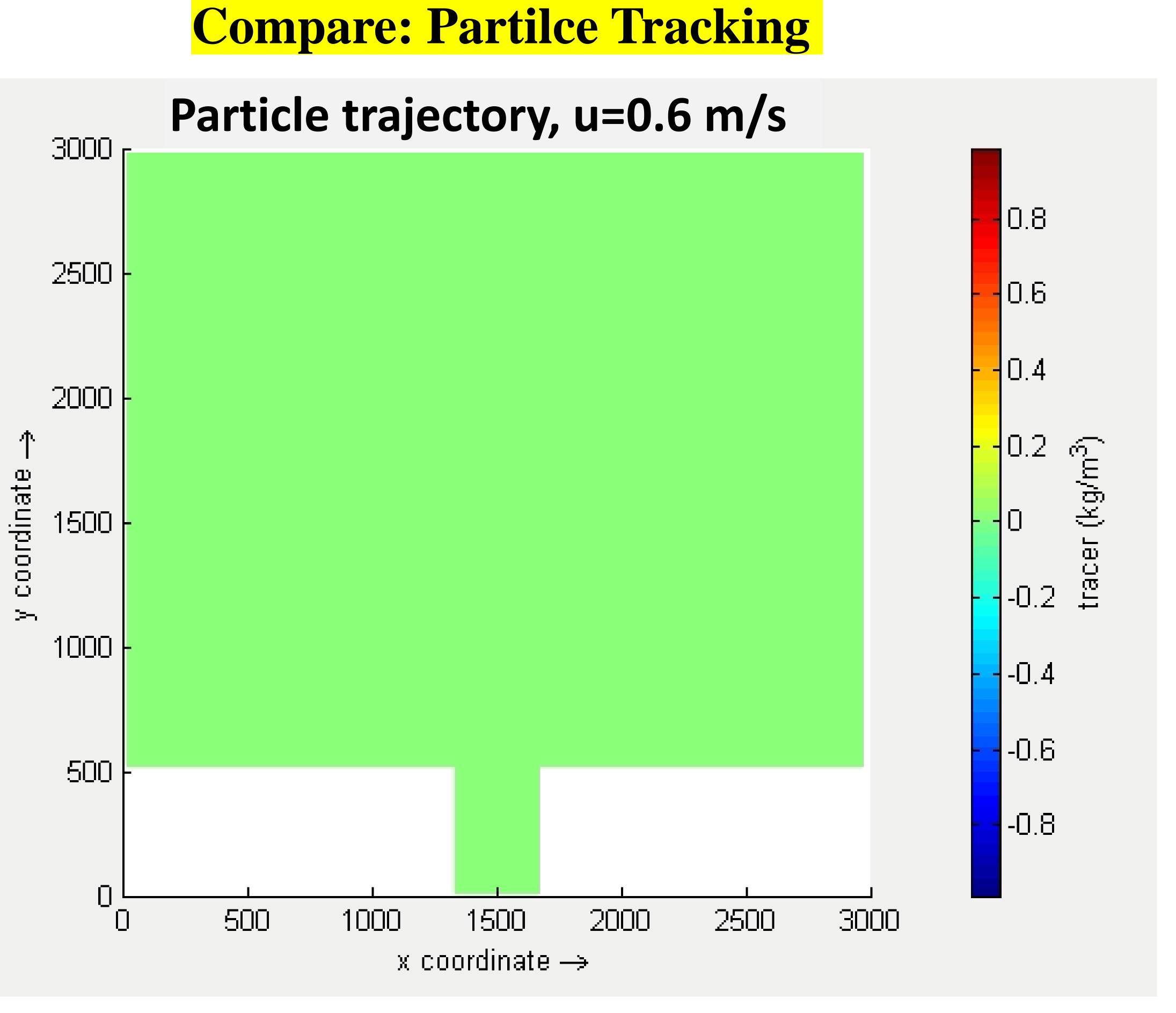




Delft3D Results

MATLAB Script Results





DC4: Mixing and transport in the coastal area

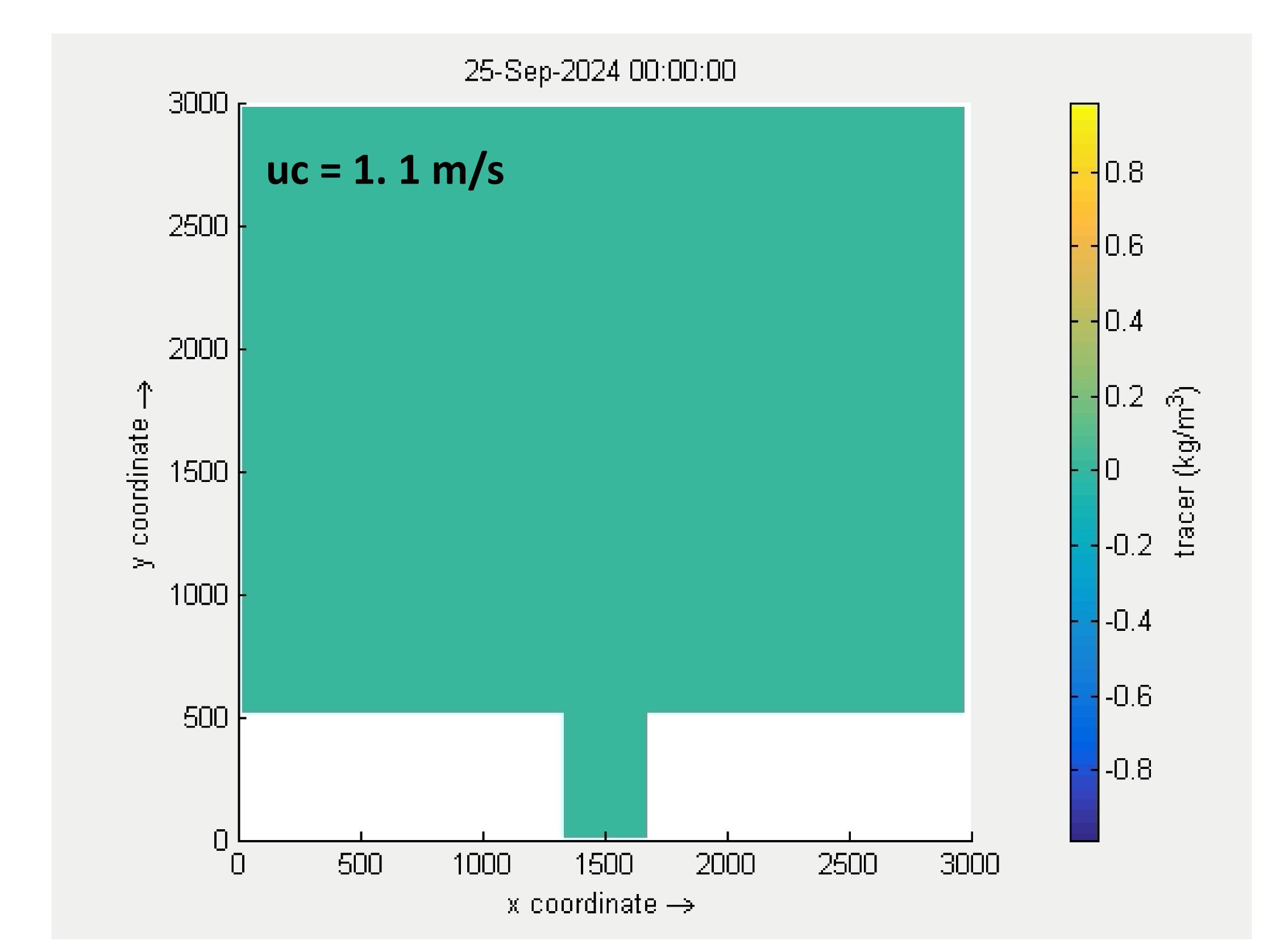


Thank you for your attention!



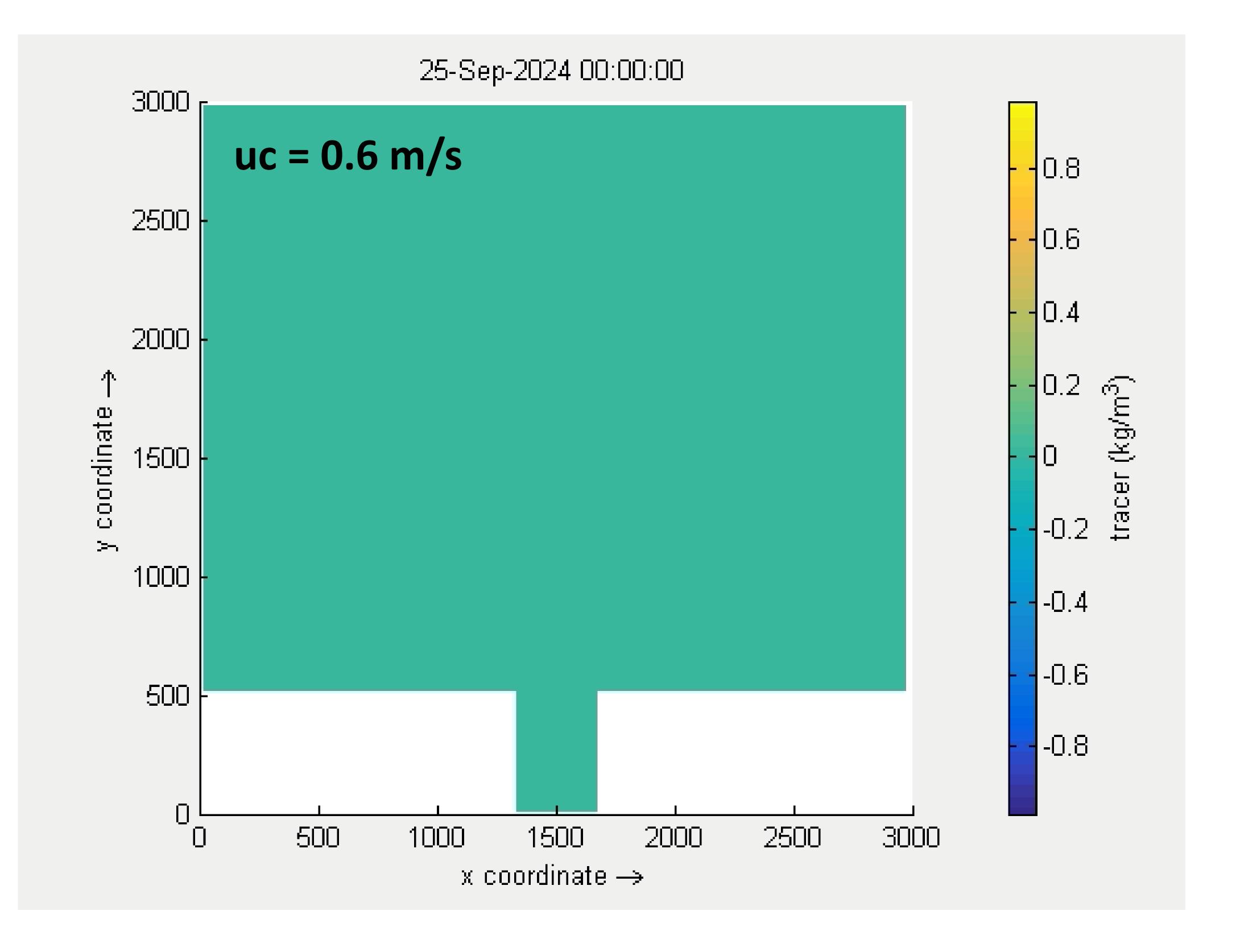
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Comparison of releasing particles from Estuarine part:

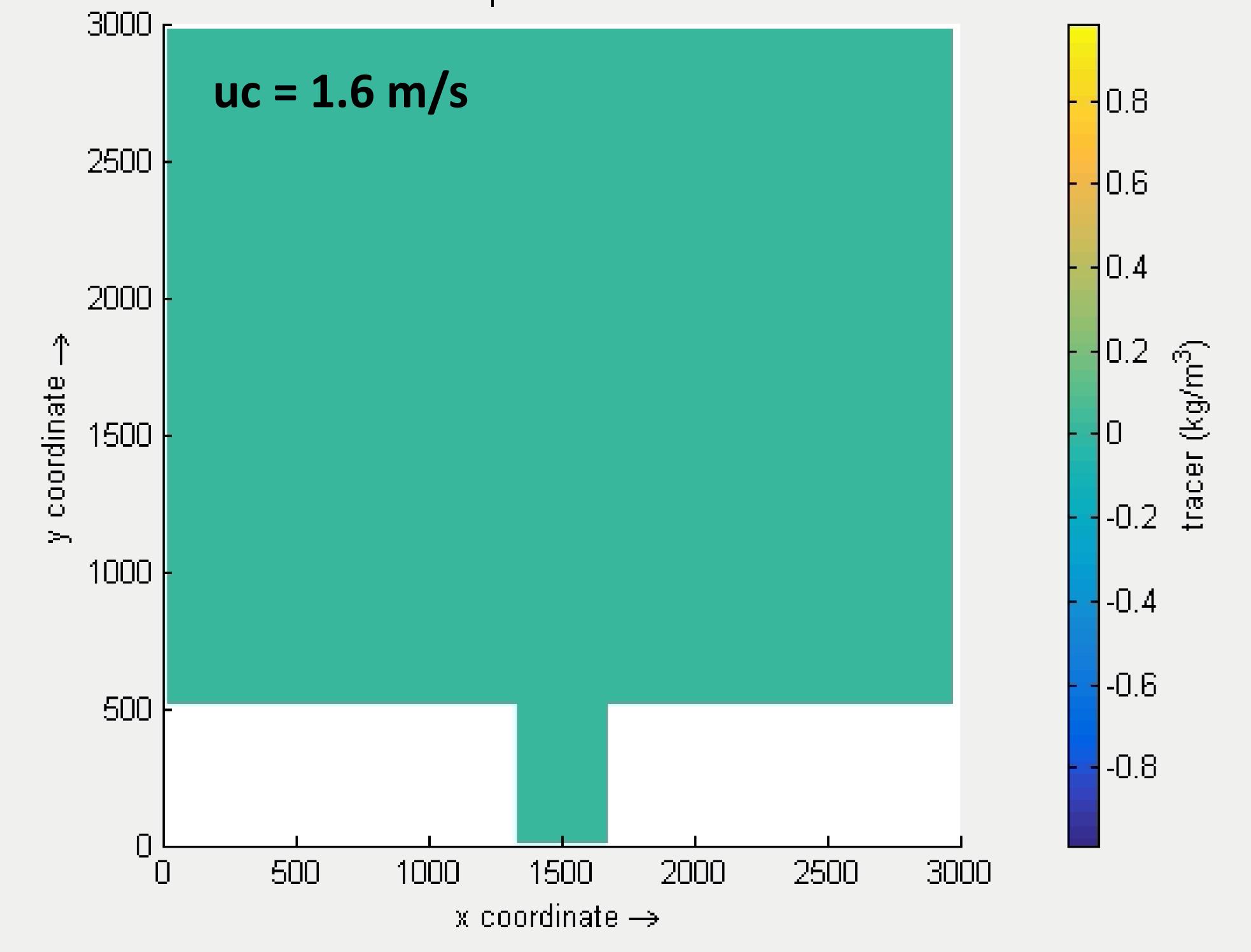




Numerical part: Delft-3D



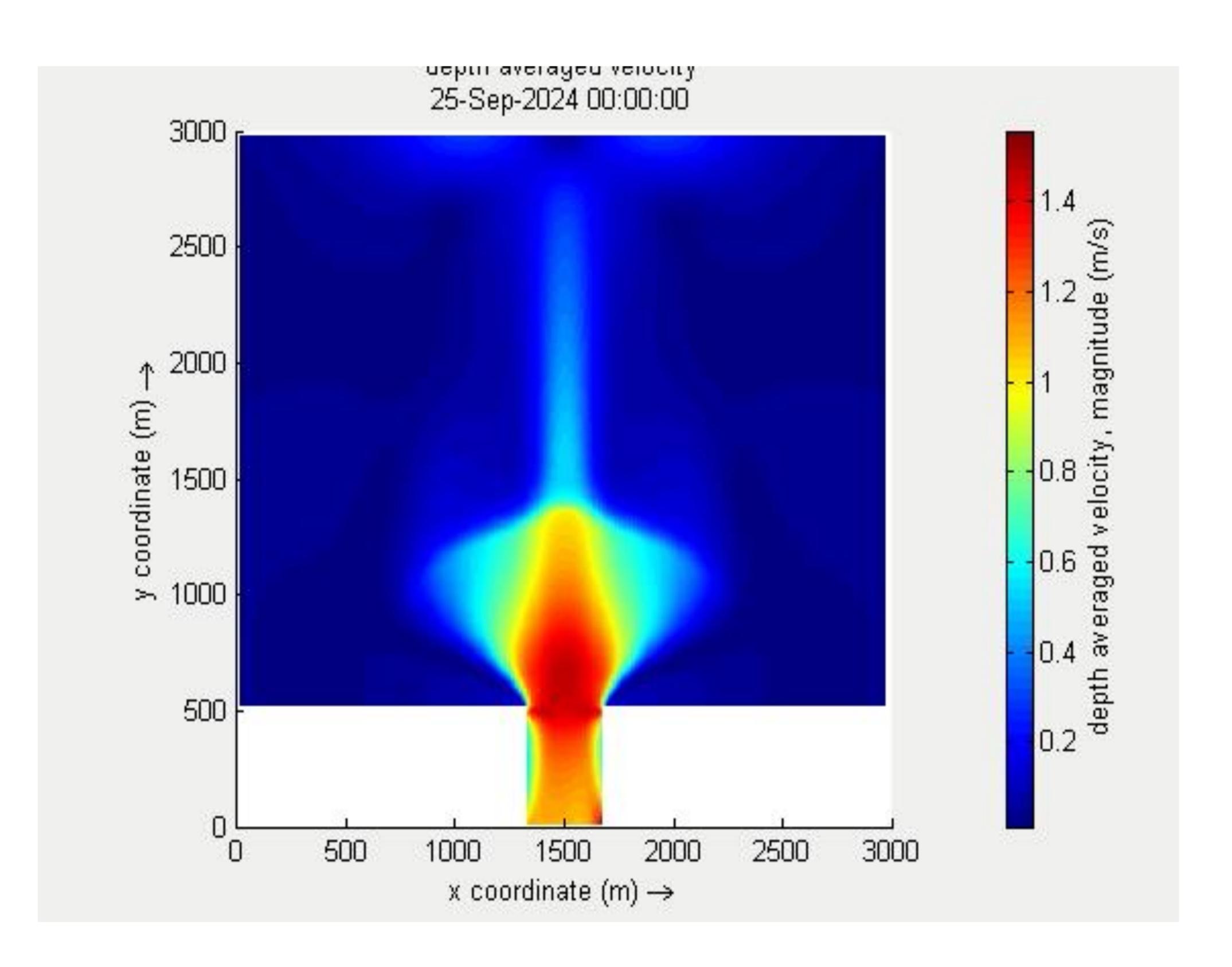


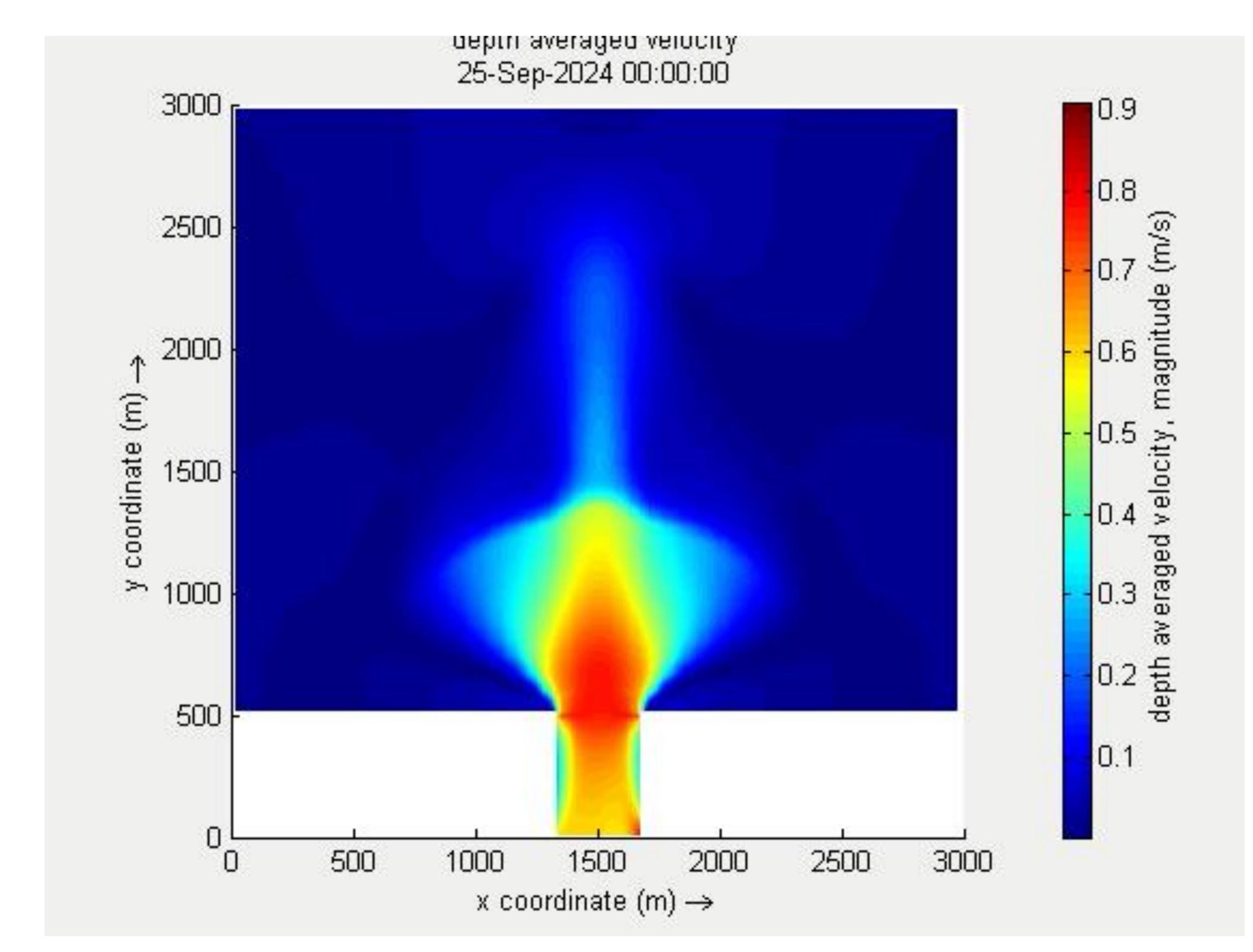




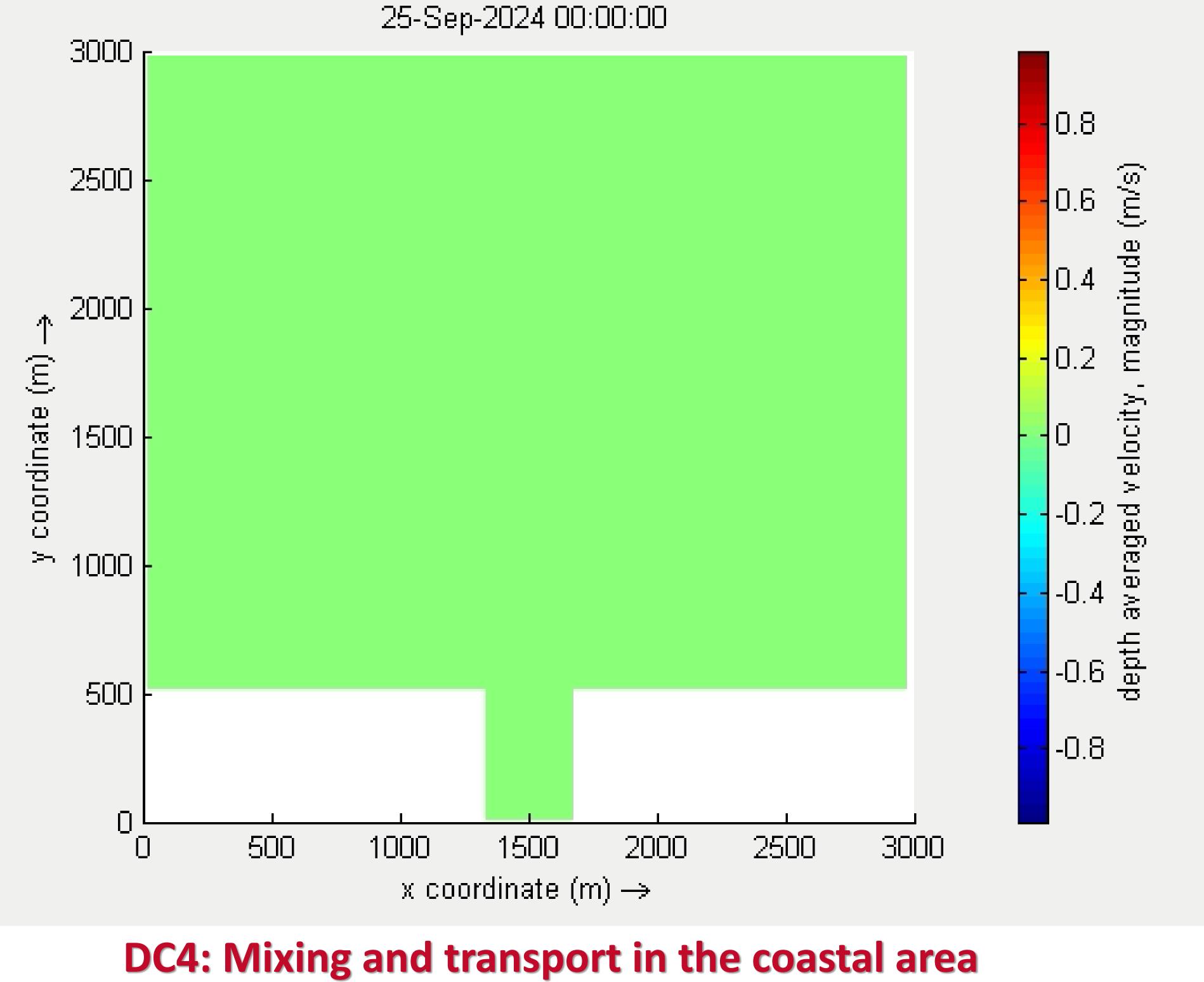


Simulation for all regions:



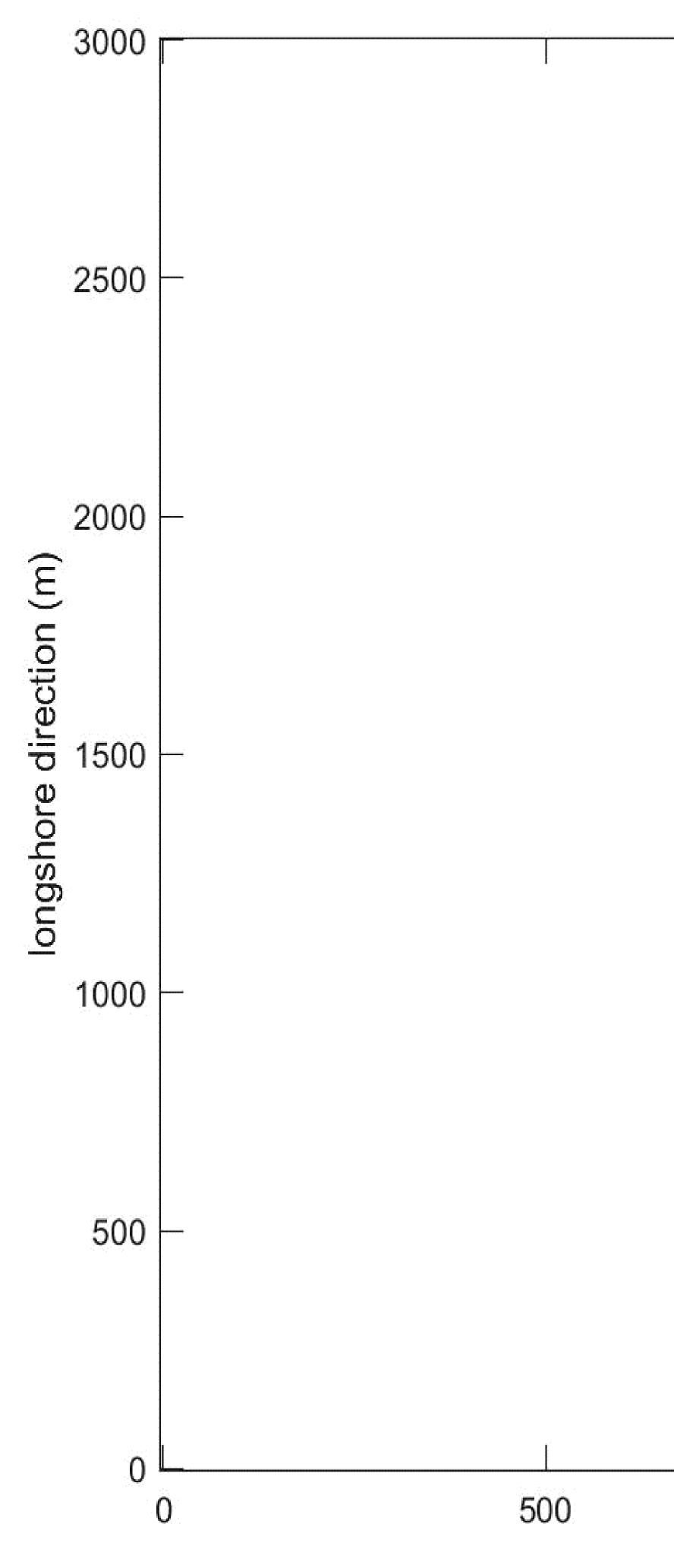


Numerical part: Delft-3D



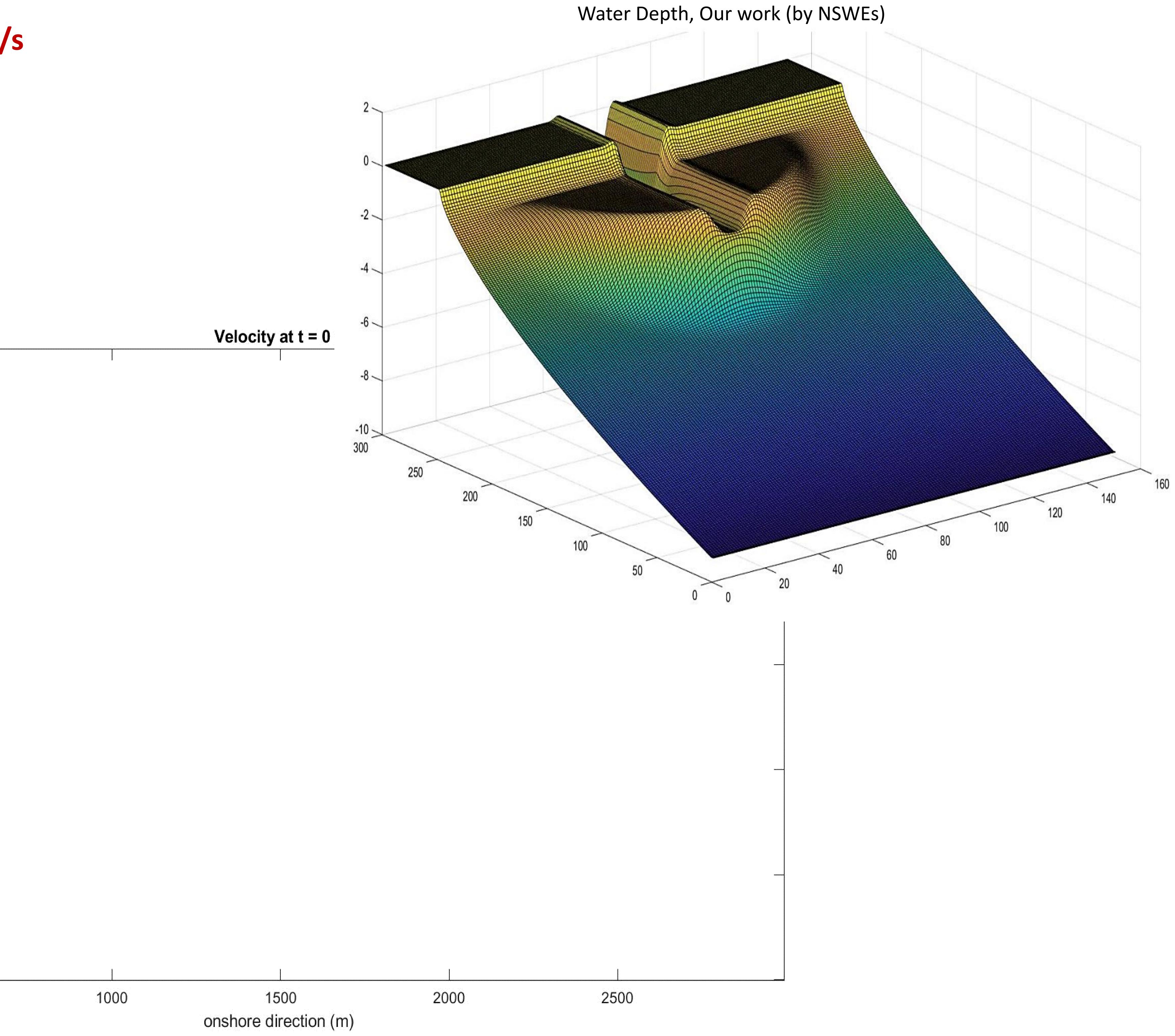


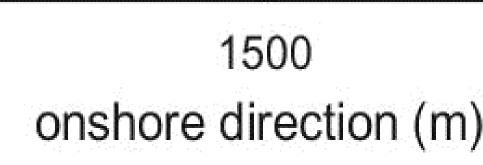
ONSWE Solver: u=0.6 m/s











DC4: Mixing and transport in the coastal area





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Sediment Transport and Morphodynamics in Marine and Coastal Waters with Engineering Solutions

Multi-model approach to scour in dynamic areas

Nishchay Tiwari HR Wallingford, UK

Supervisors and Advisors: Michiel A.F. Knaapen Sina Haeri Richard Whitehouse

Marie Curie Grant Agreement Number: 101072443



Contents

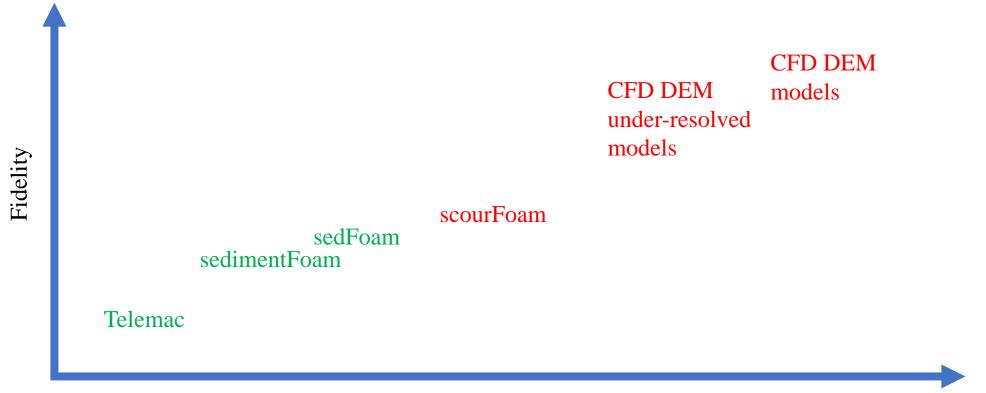
- Objective
 - State of the Art
- OpenFOAM
- sedFoam with OpenFOAM v2012
 - Methodology
 - Experimental Data
 - sedFoam equations
 - Domain and Boundary Conditions
 - Results
- sedimentFoam with FOAM-extend-4.1
 - Methodology
 - sedimentFoam equations
 - FVM to FAM
 - Preliminary Results
- Comparisons between solvers

hrwallingford



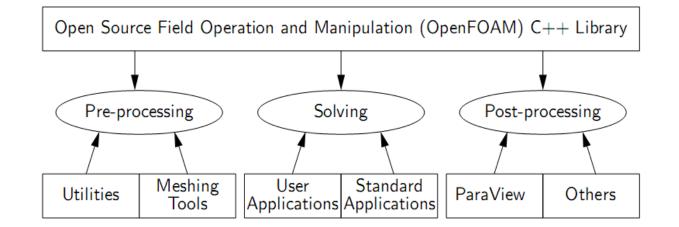
Objective & State of the Art

• To investigate and develop methods for extracting/deriving information from the CFD model to improve the accuracy and resolution of the large-scale simulations.

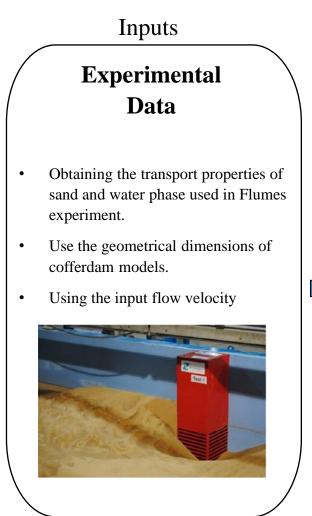


OpenFOAM

- An open-source CFD toolbox that provides ready-to-use **solvers**, **utilities**, and **libraries**.
- Offers a versatile collection of efficient, **object-oriented C++ modules**.
- Utilizes the **Finite-Volume Method (FVM)** to solve systems of partial differential equations on any **3D unstructured** polyhedral cell mesh.
- Supports efficient parallelization.



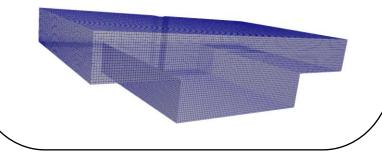
sedFoam: Numerical Modelling



Two-phase Eulerian RANS OpenFOAM

Solver

- Creating 3D domain
- Using sedFoam :
 - Granular Rheology properties (muI)
 - Interfacial properties (drag model)
 - Transport properties
 - Modified Two-phase RAS equations



	Scour Depth
•	Estimation of volume fraction of water and sand to predict the bed formation.
•	Comparison with experiments
I	
	v v √ 1(+12 - 0.03 - 0.02 - 1.01 0 0.01 - 2.3+02

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Experimental Data

- To develop a numerical model for the behaviours of scour using experiments performed on cofferdam models
- General Purpose Flume
 - Dimensions (25m long, 2.4m wide and 0.9m deep)
 - 0.5m deep bed with medium grained sand.
 - Flow discharge, Flow velocity and Water levels are recorded.
 - Duration of experiment: 50-75 hours



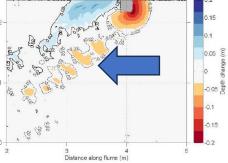


Figure 2: Test 1 result (flow from right to left); a. photograph of the post-test result (black and red graduations at 10 mm spacing) and b. difference in bed elevation between start and end of test.



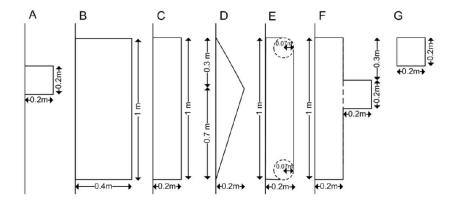


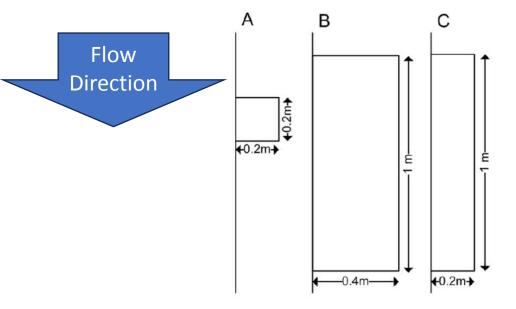
Figure 3: Cofferdam shapes used in the model experiments. The wall is on the left hand side of Structures A to F. The length of the structure is the distance along the wall and the width the distance the structure protrudes from the wall.

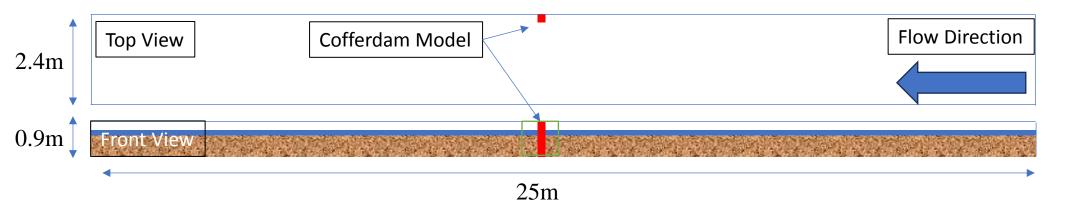
Richard J.S Whitehouse, 2021

hrwallingford

Experimental Data

- Inputs from Experiments
 - Grains (quartz sand):
 - d₁₀= 0.326mm
 - $d_{50} = 0.525$ mm
 - d₉₀= 0.673mm
 - Flow speeds : 0.177m/s, 0.244m/s (currents)
 - Water depths : 0.2m,0.1m
- Geometries of cofferdams
- Dimensions of the Flume





7

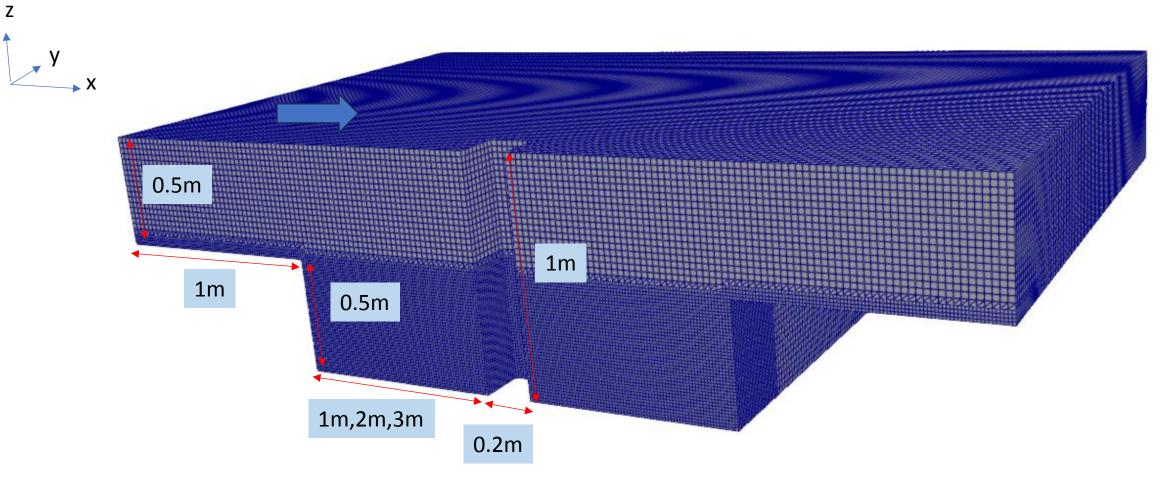
sedFoam

Closure Models for Stress Tensors

- Turbulence Models
 - SedFoam uses different turbulence closures for fluid flow, such as k-ε, k-ω, and a simple mixing length model, to capture the effects of turbulent eddies on sediment transport.
- Granular Stress Models
 - SedFoam implements granular stress models to simulate dense granular flows. The kinetic theory of granular flows and the μ(I)-rheology (derived from the Jop et al., 2006 model) are commonly used.
 - In dense flows, the granular stress is influenced by **particle-particle collisions** and **interparticle friction**, represented by the effective viscosity, which is a function of the shear rate and pressure.
 - Unlike the kinetic theory of granular flows (which works well for dilute conditions), the μ(I) rheology is phenomenological and based on dimensional analysis, focusing on frictional contacts.



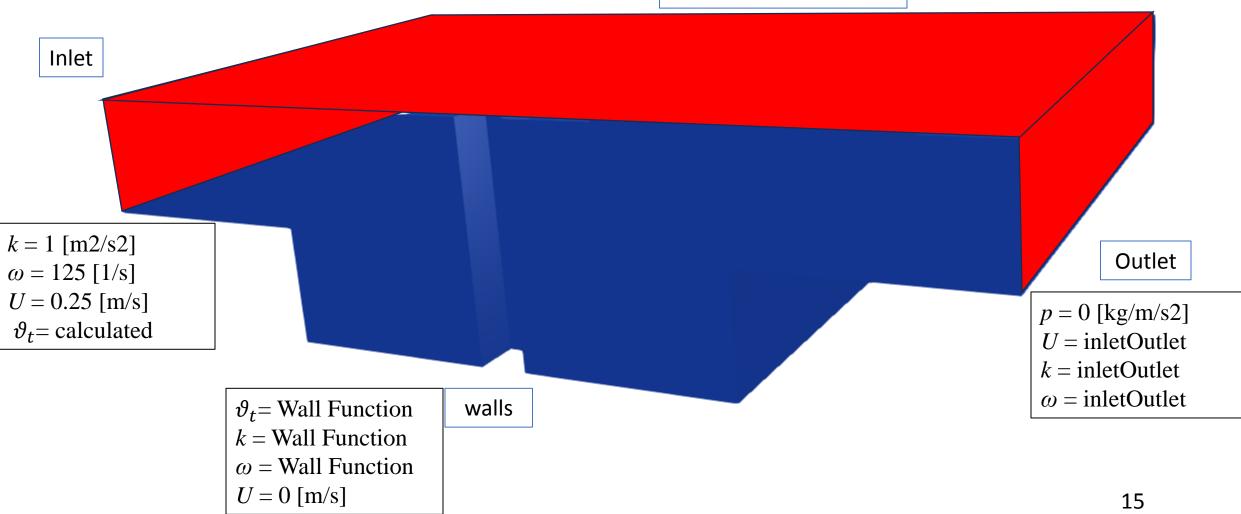
Domain and Boundary Conditions





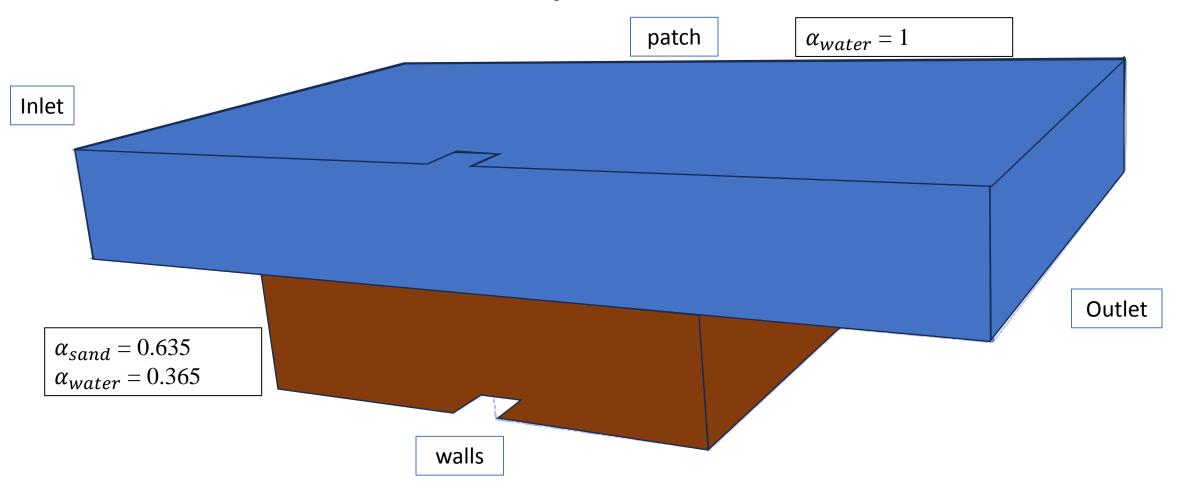
Domain and Boundary Conditions

Patch - zeroGradient



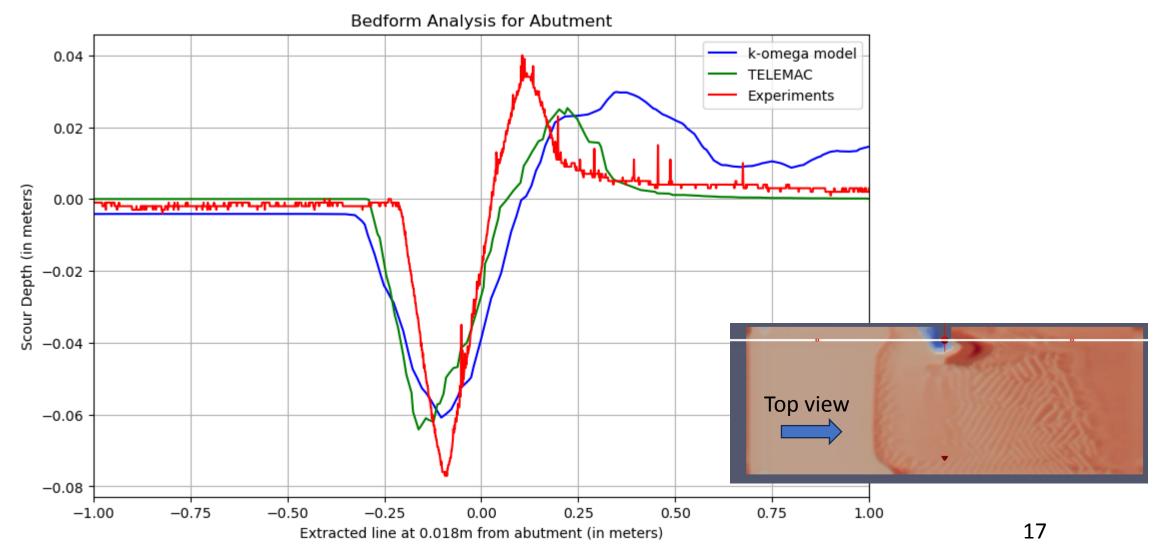


Domain and Boundary Conditions



Results

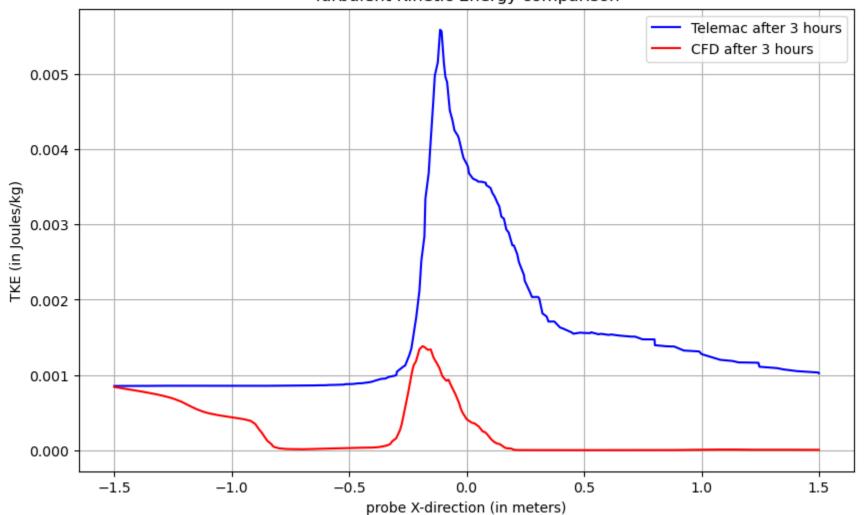
Bedform comparison at 3hr (scour depth: surface elevation)



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Results

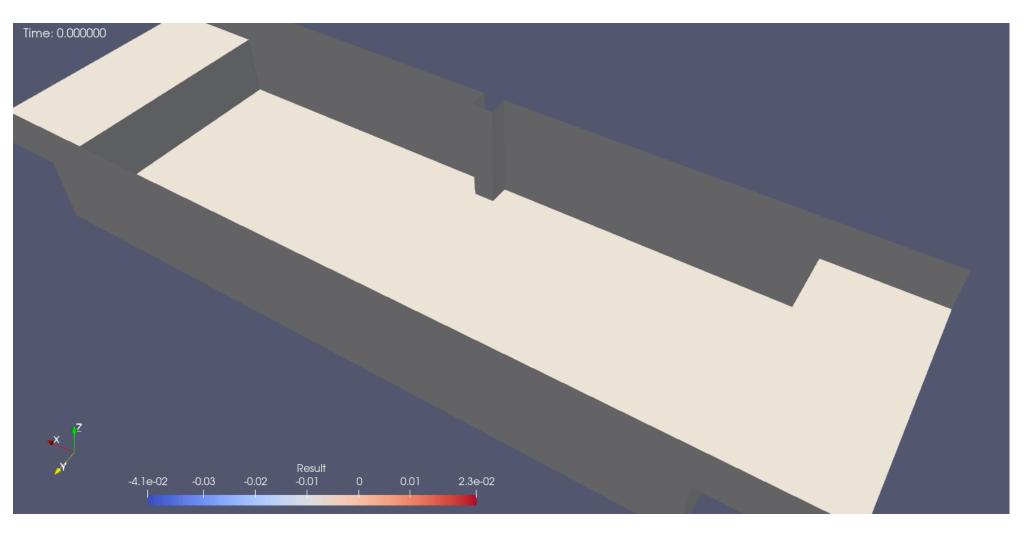
TKE profile at 3hr



Turbulent Kinetic Energy comparison

Results

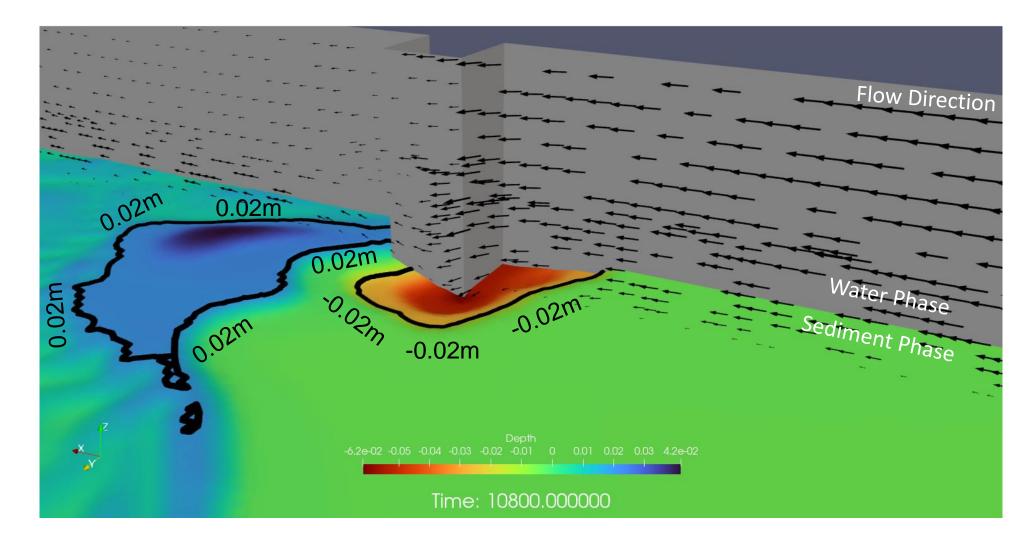
Bedform evolution (Surface Elevation)



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Results

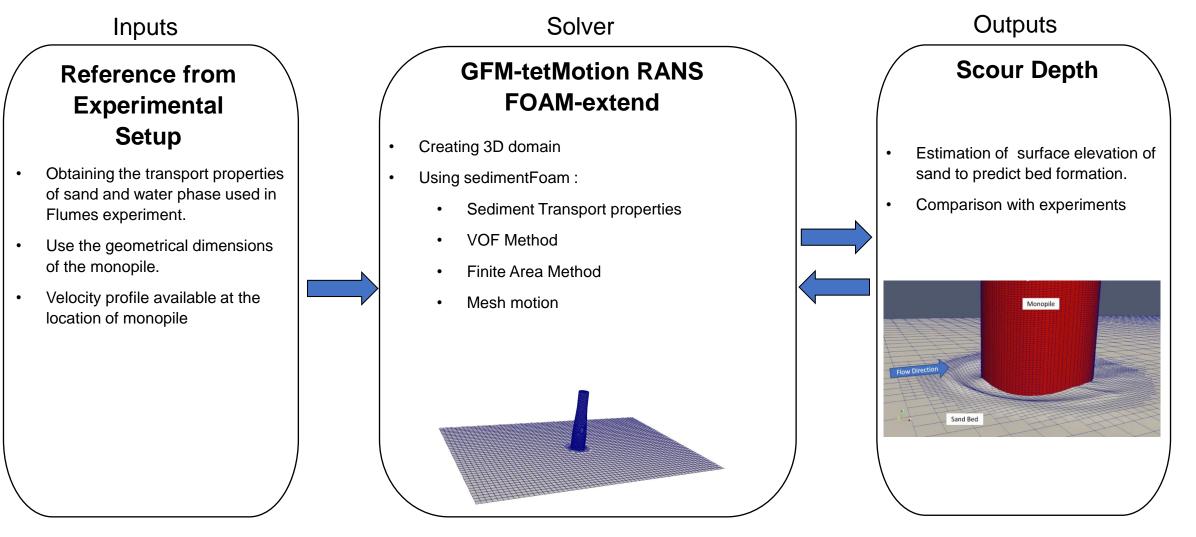
Bedform evolution (Surface Elevation) after 3 hours



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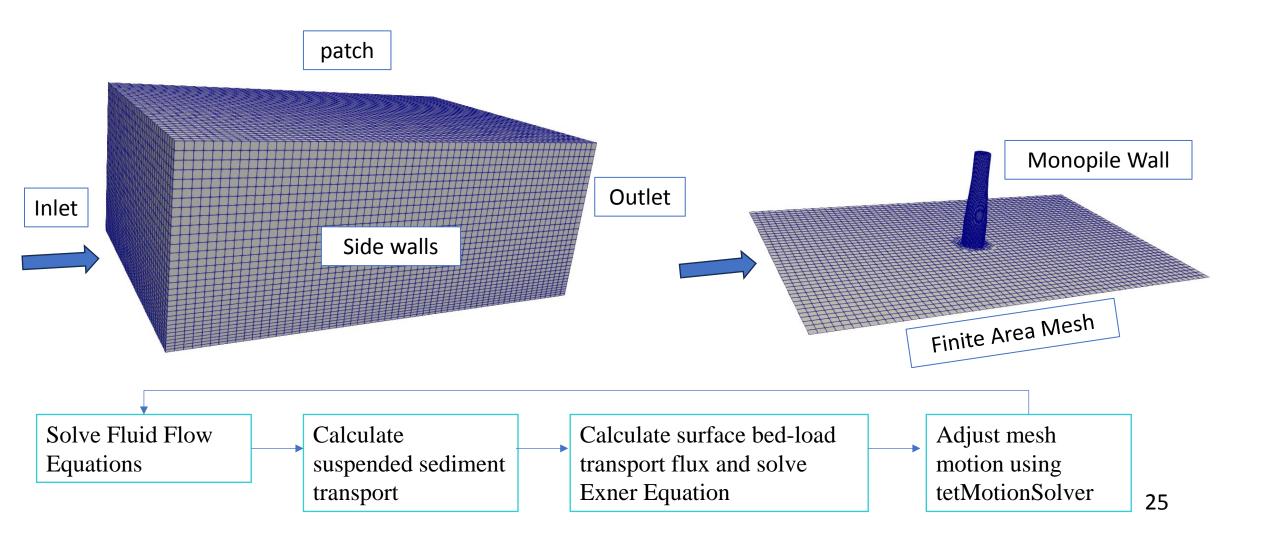
sedimentFoam: Numerical Modelling





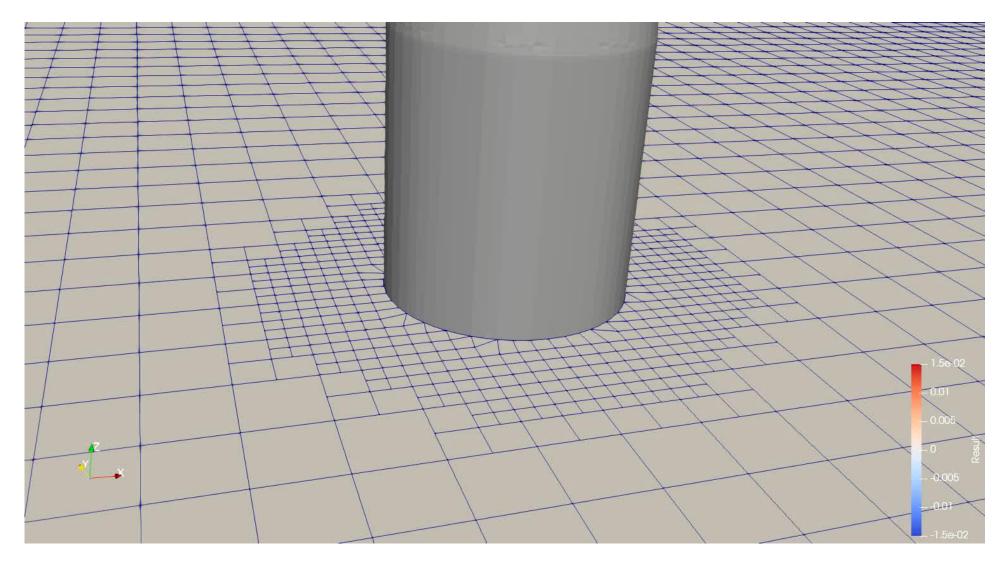
hrwallingford

FVM to FAM



Preliminary Results

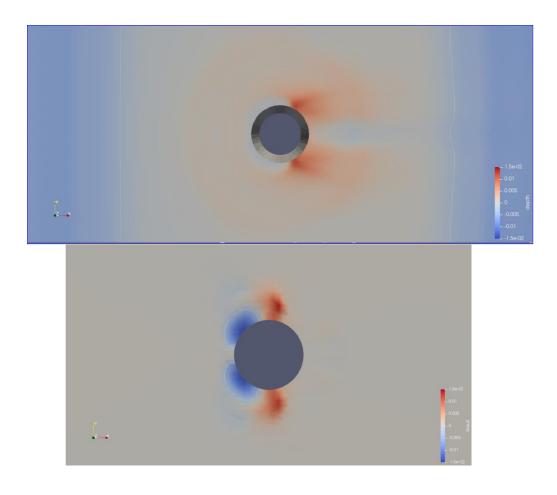
(80 seconds)

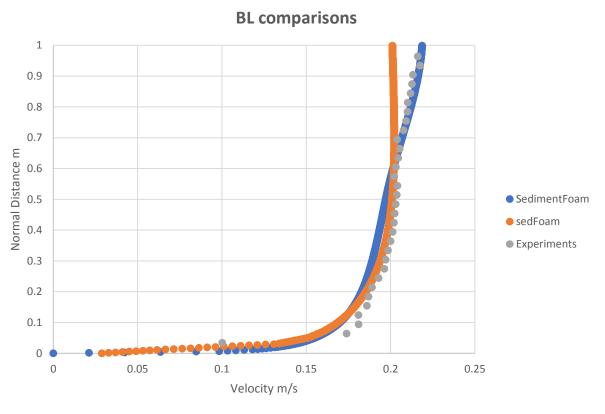


(80 seconds)

sedFoam vs sedimentFoam

hr wallingford





At 600 seconds





Thank You









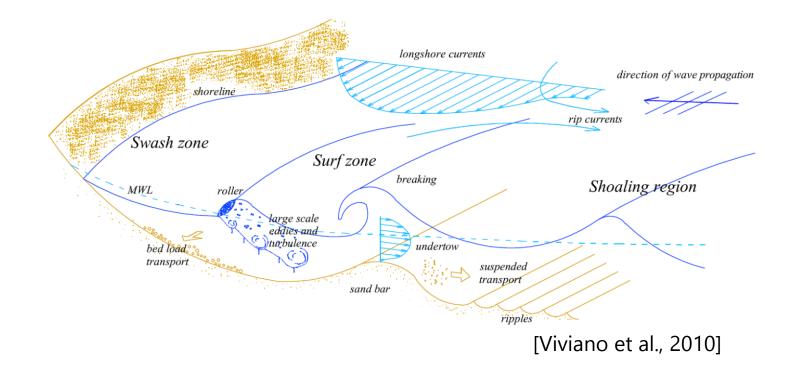
SEDIMARE DC MEETING 07.11.2024-08.11.2024 Delft, the Netherlands

Nearshore Wave Processes by Remote Sensing

Name Muhammed Said Parlak Maurizio Brocchini Supervisors Nicholas Dodd Matteo Postacchini

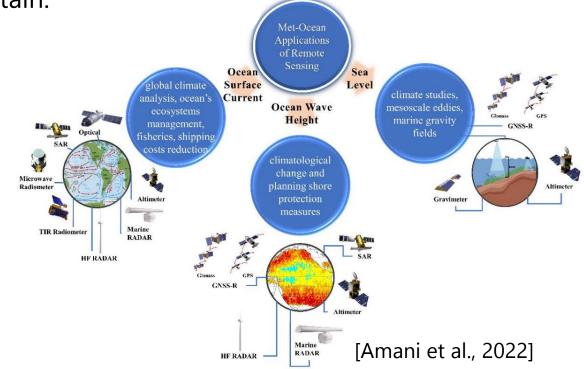
Introduction

- \rightarrow Nearshore dynamics are dominated with nonlinear mechanisms due to multiple interactions.
- \rightarrow Wave propagation is one of the crucial driver of these interactions and nearshore sea state.
- \rightarrow Observation of sea state is essential for human activities and marine environment.



Introduction

- \rightarrow Nearshore measurement with in-situ tools is a tough task due to complex processes.
- \rightarrow They have many components and need regular maintenance and calibration.
- → Remote sensing tools judicious alternatives since they have larger coverage area, and they are easier to maintain.



Study Site

 \rightarrow Senigallia, Italy

3

- Senigallia Harbour, Misa River Estuary
- Strong interactions between sea and river



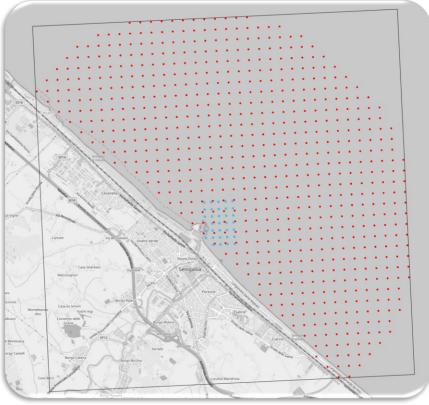
[Brocchini et al., 2017]

Study Site

 \rightarrow Senigallia, Italy

- SGS: 5 cameras, 2 Hz sampling, 10 min-long
- XBR: X-Band RADAR, 0.5 Hz sampling, 63 scans



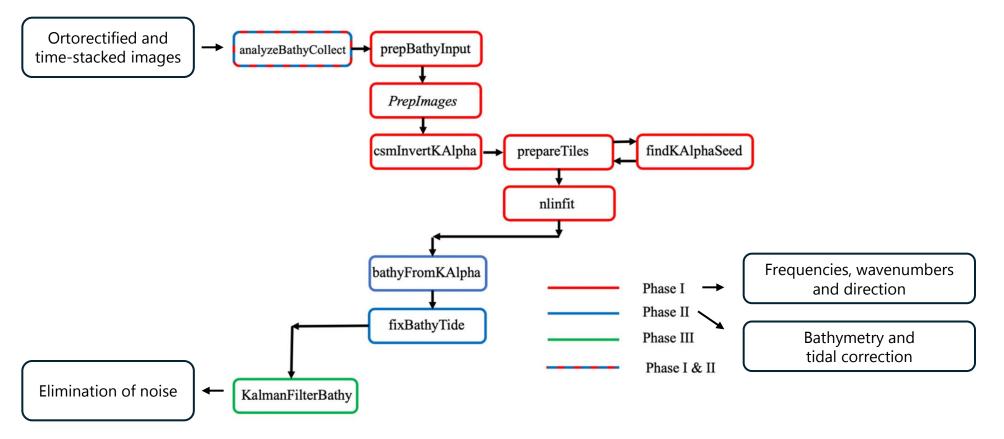


Locations of XBR and SGS

Coverage area

 \rightarrow Processing data obtained from remote sensing tools (cBathy).

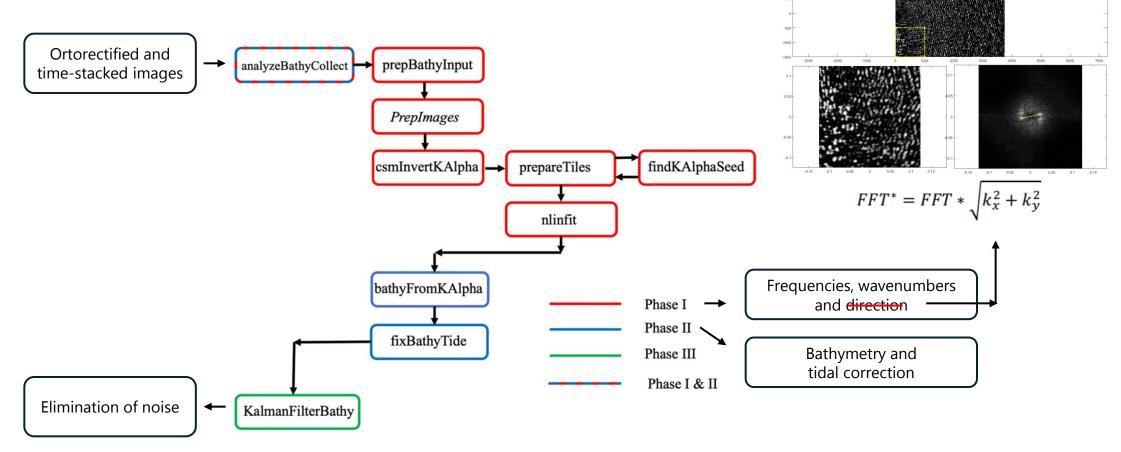
- cBathy algorithm flowchart



[Holman et al., 2013]

 \rightarrow Processing data obtained from remote sensing tools (cBathy).

- cBathy algorithm flowchart



[Holman et al., 2013]

 \rightarrow Processing data obtained from remote sensing tools (cBathy).

- Assumed that the most coherent frequency is close to the peak frequency, f_p .
- Hence:

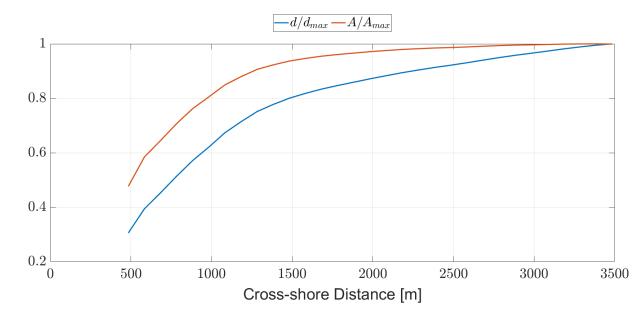
$$C_p = \frac{\omega_p}{k} = \frac{2\pi f_p}{k} \qquad \lambda_p = \frac{2\pi}{k}$$

- C_p is the peak celerity, ω_p is the peak angular frequency, k is the wavenumber, and λ_p is the peak wavelength.

 \rightarrow Processing data obtained from remote sensing tools (cBathy).

- Significant wave height, H_s , is approximated with an emprical approach based on water depth, d, and C_p .

$$H_{s} = \frac{C_{p}^{2}}{g \; A^{*} \; 1.8} \to A^{*} = \begin{cases} d^{\frac{1}{2}}, & d < d_{th} \\ d^{\frac{1}{3}} - d^{\frac{1}{3}}_{th} + d^{\frac{1}{2}}_{th}, & d \ge d_{th} \end{cases}$$

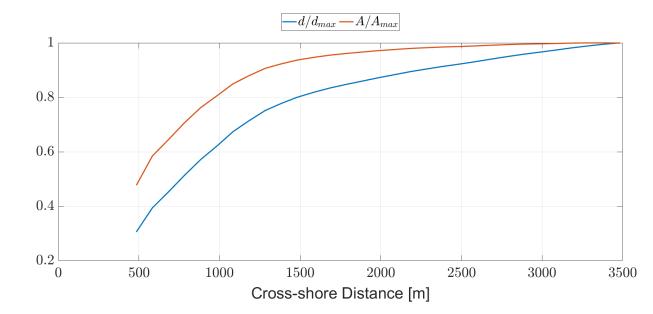


 \rightarrow Processing data obtained from remote sensing tools (cBathy).

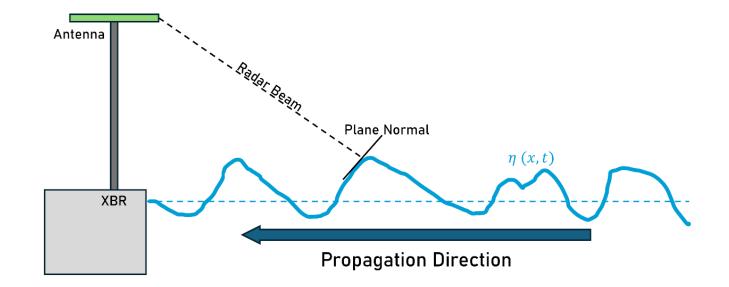
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- d_{th} depends on the bathymetric gradient e.g., 0,2 in this case

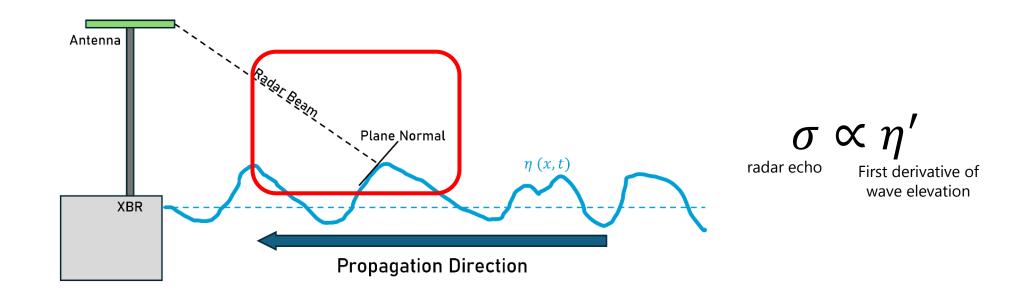


- The wave field is visible to XBR when the waves are relatively energetic.
- Radar return echo is maximized at the wave fronts.



 \rightarrow Reconstruction of wave propagation (preliminary).

- Assumed that radar signals are correlated with wave slopes.



 \rightarrow Reconstruction of wave propagation (preliminary).

- Processing of radar signals:

 $\hat{\sigma} = \sigma - \bar{\sigma}$ Mean extraction $\sigma^* = \hat{\sigma} + |\min(\hat{\sigma})|$ Shifting signal above the zero axis $\boldsymbol{P} = \max(\sigma^*)_{lc}$ Local maxima Δ $6 \overline{}^{\times 10^4}$ 5 σ^* 3 200400 600 800 1000 0 Cross-shore distance [m]

 \rightarrow Reconstruction of wave propagation (preliminary).

- Crest, trough slopes are arranged to resemble a sawtooth wave formation.
- Assumed that trough slopes are half of the average adjacent crest slopes.

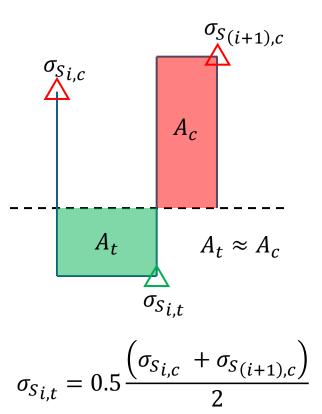
 $\sigma_{S(i+1),c}$

 $\sigma_{S_{i,c}}$

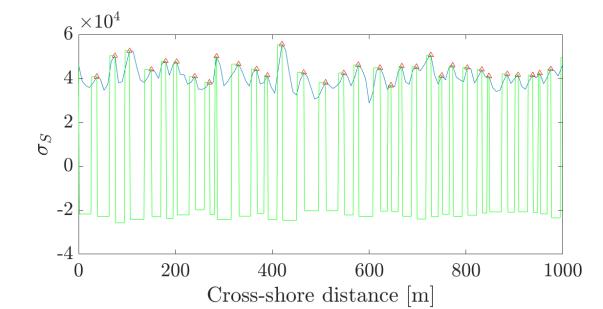
 $\begin{aligned} & \stackrel{\Delta}{\sigma_{S_{i,t}}} \\ & \sigma_{S_{i,t}} = 0.5 \frac{\left(\sigma_{S_{i,c}} + \sigma_{S_{(i+1),c}}\right)}{2} \end{aligned}$

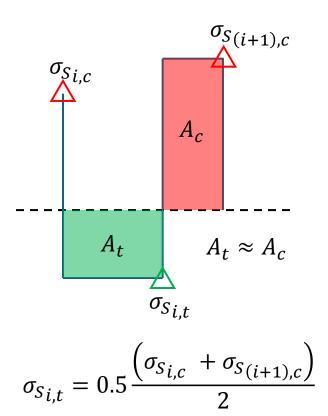
10

- Crest, trough slopes are arranged to resemble a sawtooth wave formation.
- Assumed that trough slopes are half of the average adjacent crest slopes.
- The *x* location is determined to ensure zero-mean signal.

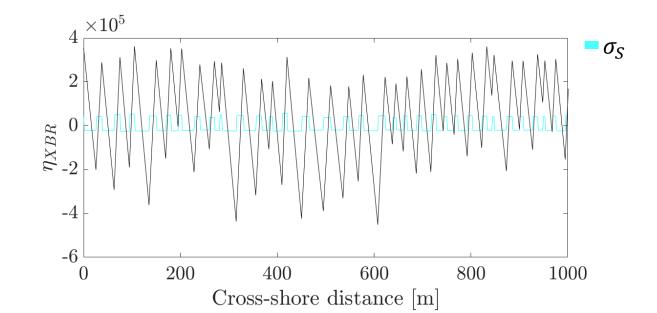


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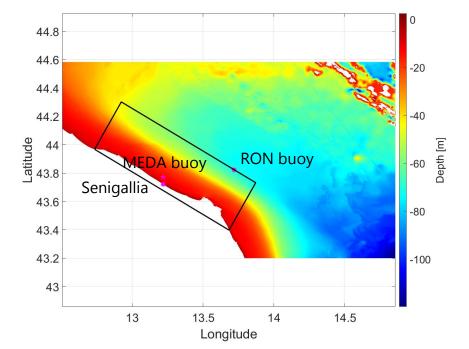


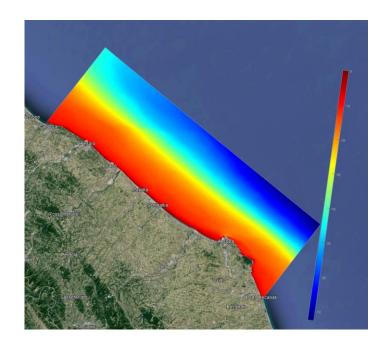


- Crest, trough slopes are arranged to resemble a sawtooth wave formation.
- Assumed that trough slopes are half of the average adjacent crest slopes.
- The *x* location is determined to ensure zero-mean signal.
- Signal is integrated in spatial domain.



- \rightarrow Model Chain (SWAN).
 - Spatial domain is determined by considering the available buoy locations.

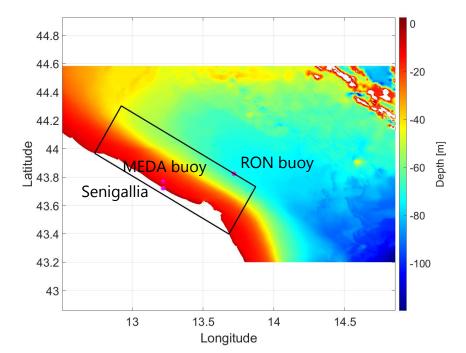


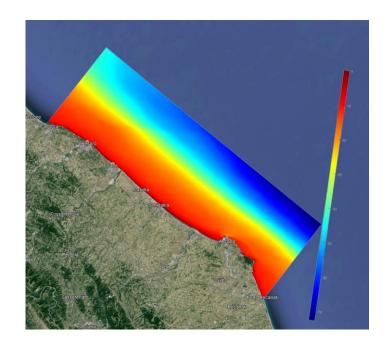


\rightarrow Model Chain (SWAN).

- Spatial domain is determined by considering the available buoy locations.
- 3 open boundaries forced with wave characteristics obtained from RON.
- The grid is forced with winds from COPERNICUS database



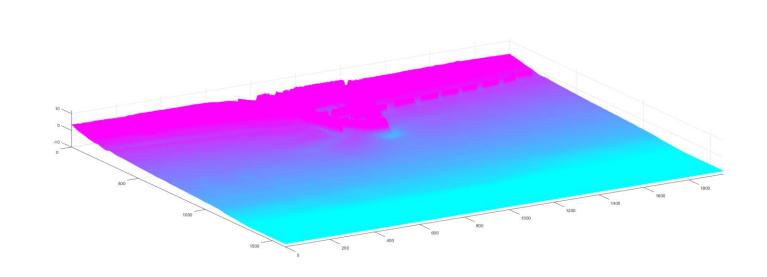




\rightarrow Model Chain (FUNWAVE).

- Spatial domain is constructed around the harbour.





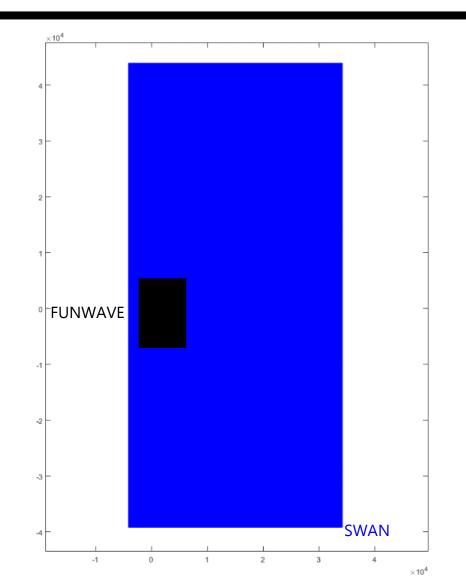
\rightarrow Model Chain (FUNWAVE).

- Spatial domain is constructed around the harbour.
- WAVEMAKER input is obtained from SWAN results

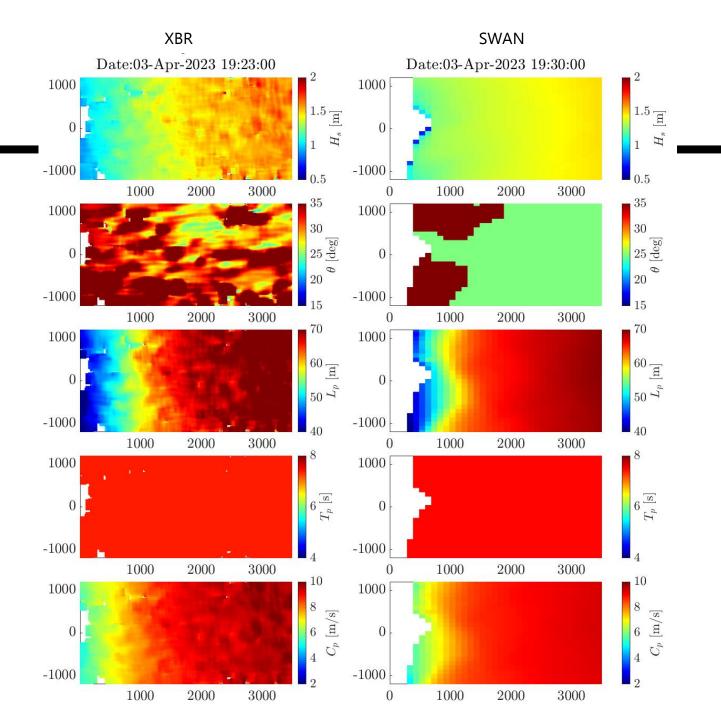
Properties of WM_IRR WAVEMAKER

Parameter	depth [m]	$f_P [Hz]$	$H_{m0}[m]$	$\theta_P \left[deg \right]$	<i>Υτμα</i>	f _{min} - f _{max}
Value	8.17	0.1337	1.27	-12.57	1.5	0.04 – 1.0

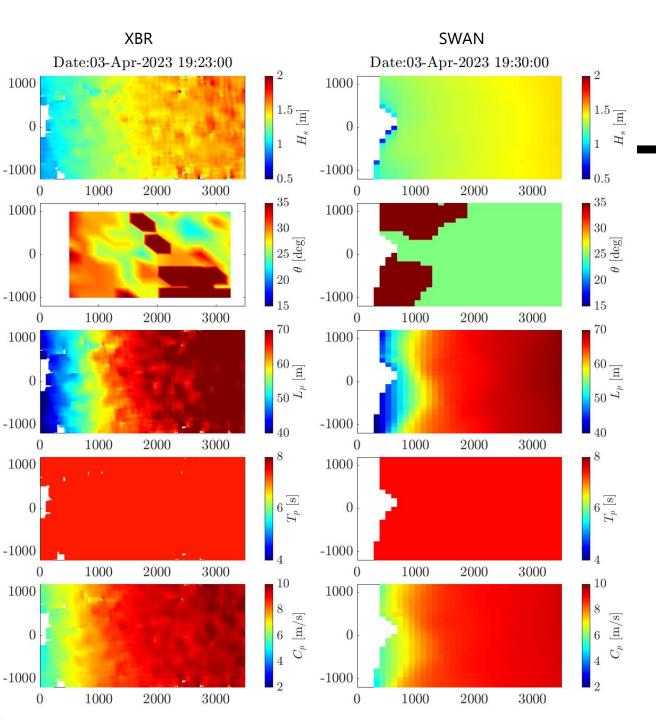
dx=dy=1 m Grid Size= 1599x1957



- \rightarrow XBR sea state estimation performance.
 - L_P , T_P and C_P are captured well



- \rightarrow XBR sea state estimation performance.
 - L_P , T_P and C_P are captured well
 - θ_P estimation is better with new approach



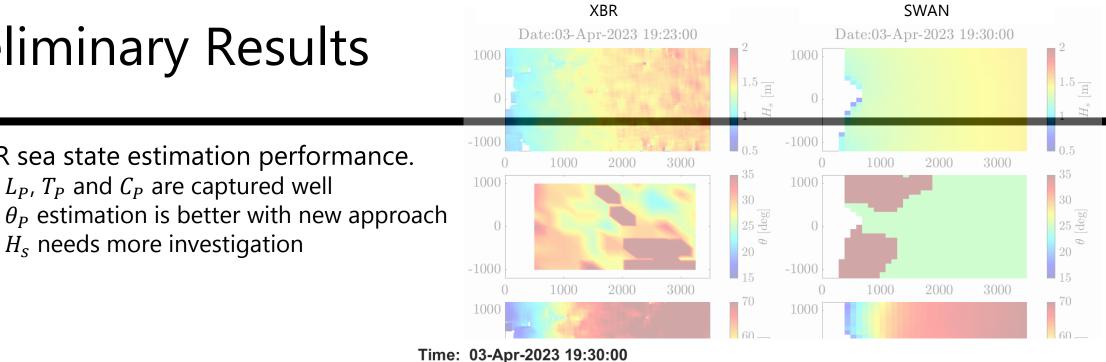
 \rightarrow XBR sea state estimation performance.

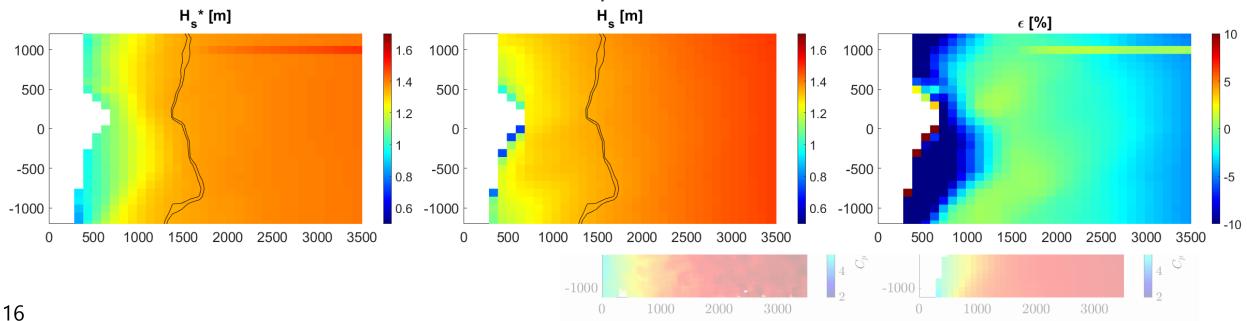
- L_P , T_P and C_P are captured well

 H_s needs more investigation

-

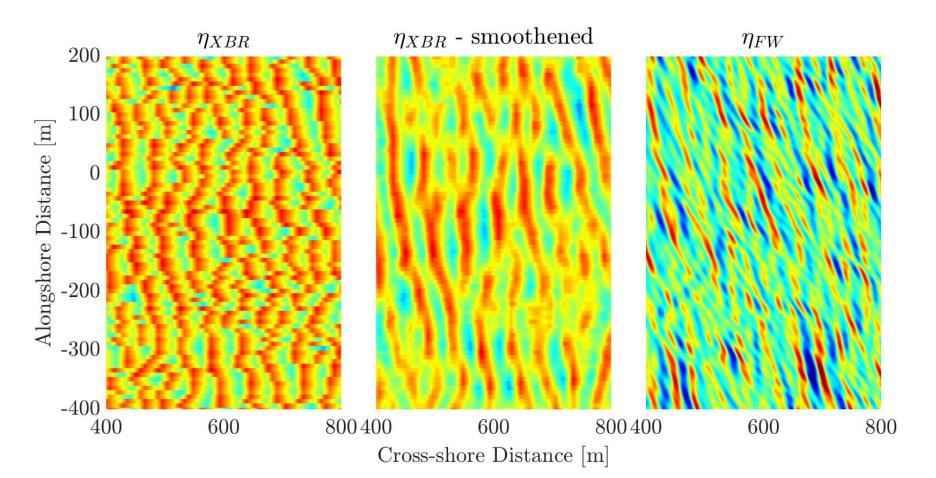
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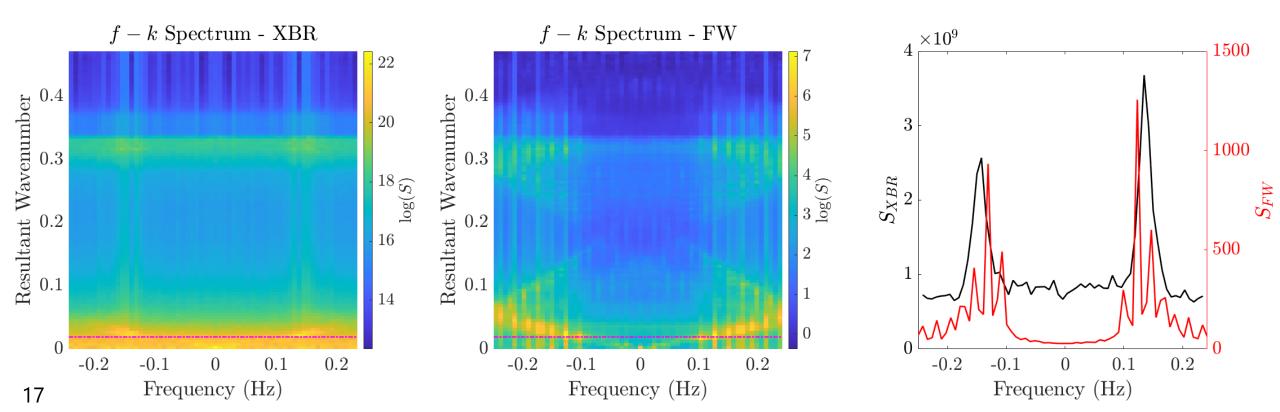
 \rightarrow XBR wave propagation.

- Proposed approach is compared to FUNWAVE results

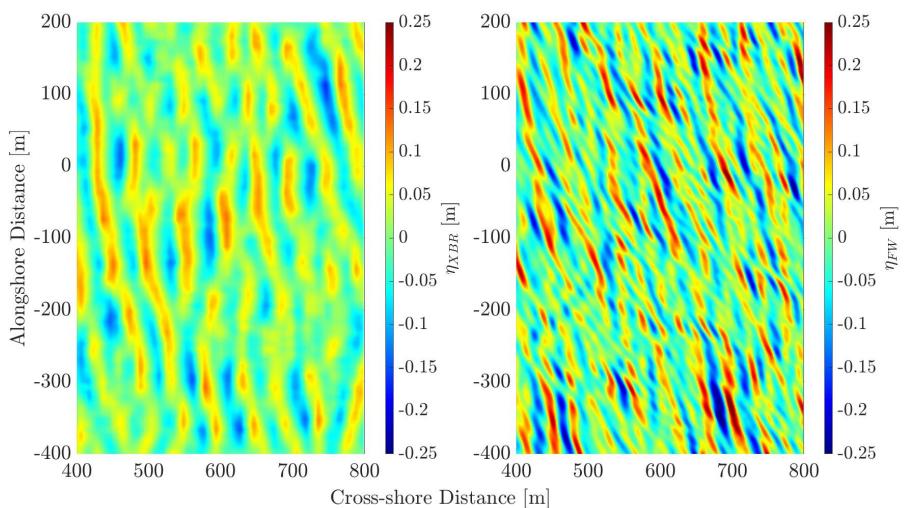


 \rightarrow XBR wave propagation.

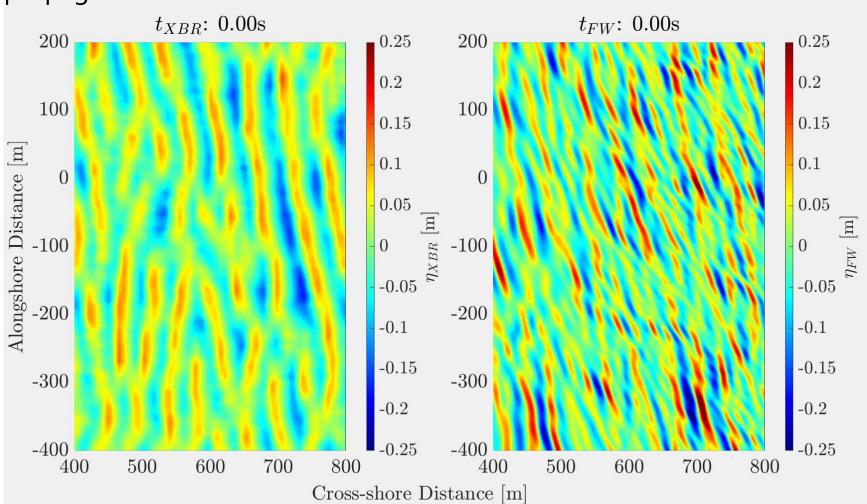
- Scaler determined by f - k spectra



 \rightarrow XBR wave propagation.

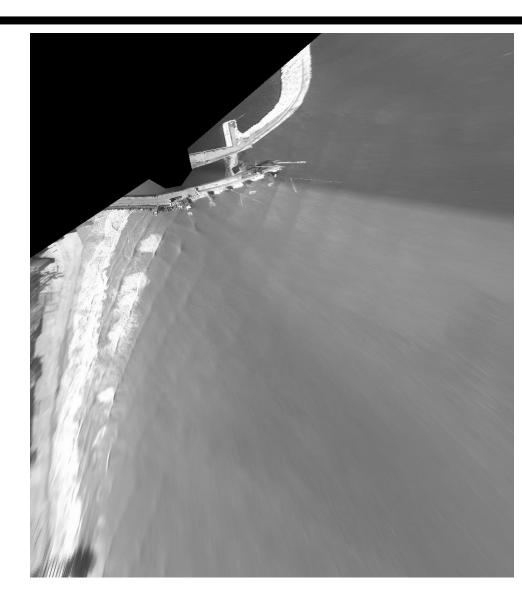


 \rightarrow XBR wave propagation.

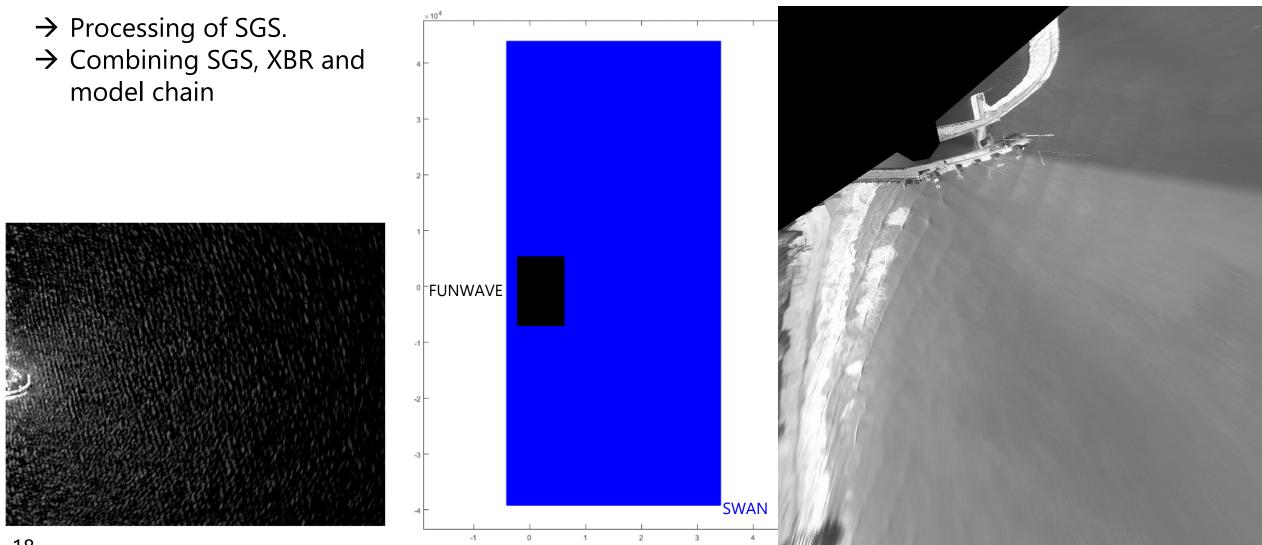


Ongoing Work

 \rightarrow Processing of SGS.



Ongoing Work











SEDIMARE DC MEETING

07.11.2024-08.11.2024 Delft, the Netherlands

Nearshore Wave Processes by Remote Sensing

Muhammed Said Parlak <u>m.s.parlak@univpm.it</u>

https://sedimare.eu

Acknowledgement: This project has received funding from the European Union's (EU) Horizon Europe Framework Programme (HORIZON) under Grant Agreement No 101072443 as a MSCA Doctoral Network (HORIZON-MSCA-2021-DN-01) of

UNIVERSITY OF TWENTE.

SEDIMARE 2023 - 2027

Sediment Transport and Morphodynamics in Marine and Coastal Waters with Engineering Solutions

SEDIMARE PROJECT_DC #3

Netherlands Meeting

EROSION AND TRANSPORT OF SAND-SILT MIXTURES

PhD Candidate:Nguyen, Thi To Van (Van)Promotor:P.C. Roos (Pieter)Co-promotor:J.J. van der Werf (Jebbe)

Date: 2024 Nov 07th

Contents

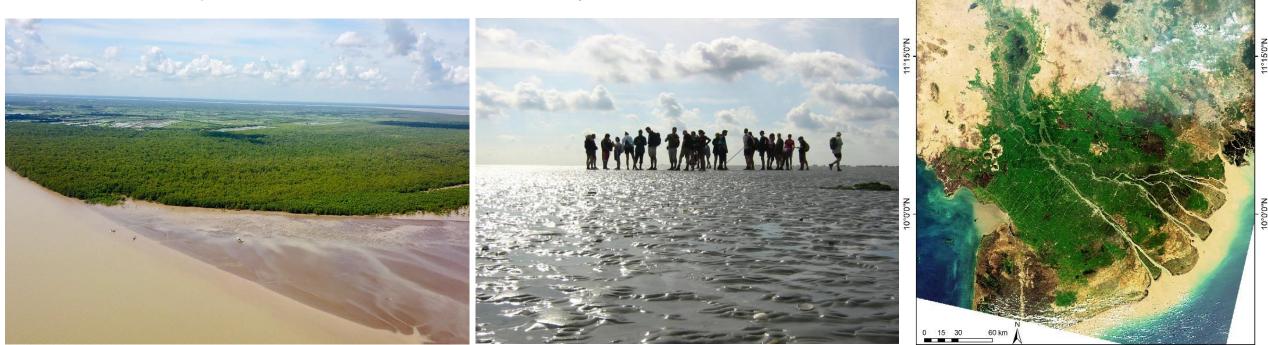
1. Introduction

- 2. Research progress
- 3. Upcoming plans
- 4. Conclusion

1. Introduction

In nature, especially in coastal and fluvial systems, most of sediments are mixes of sand and fines (clay and silt).

Recently, there have been more studies focusing on the transport of sand-mud mixtures. However, most of these studies treated clay and silt collectively as mud (*Mitchener and Torfs, 1996; Van Ledden, 2003; Jacobs, 2011; Winterwerp et al., 2012; Colina Alonso et al., 2023*).



Silty sediments in Mekong Delta, Vietnam [Photo by <u>MangLub project</u>] Mudflat, Wadden Sea, Netherlands.

Satellite images of the Mekong Delta taken by Envisat

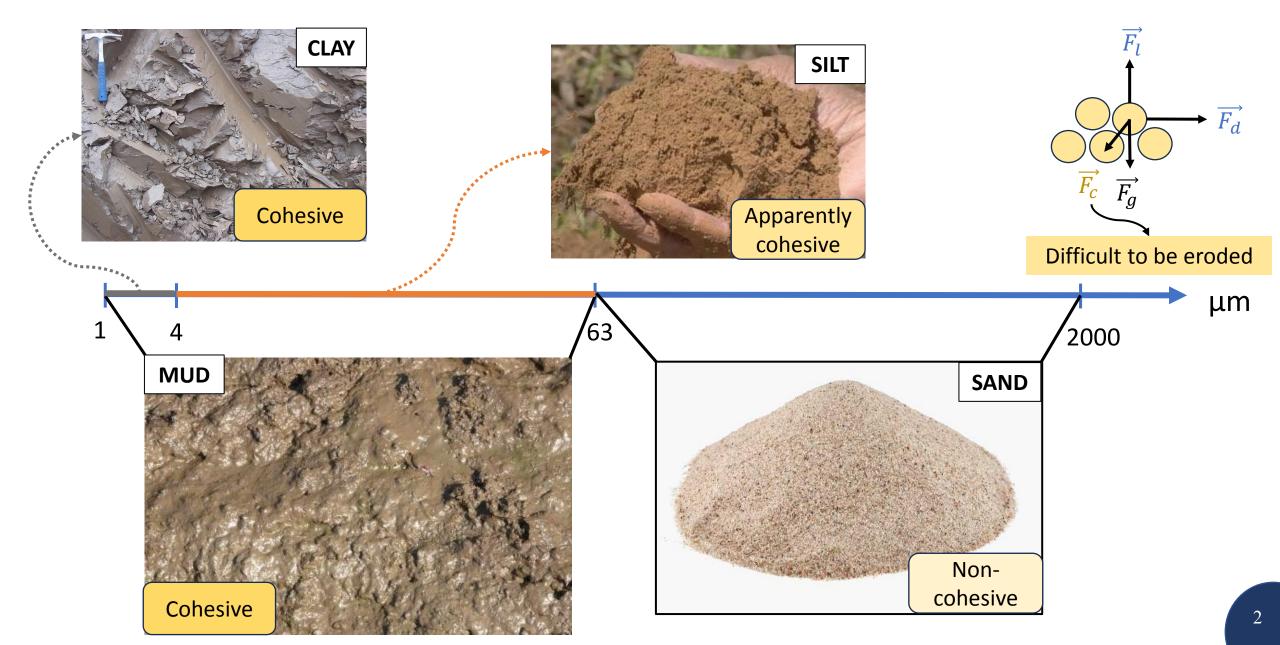
106°15'0"E

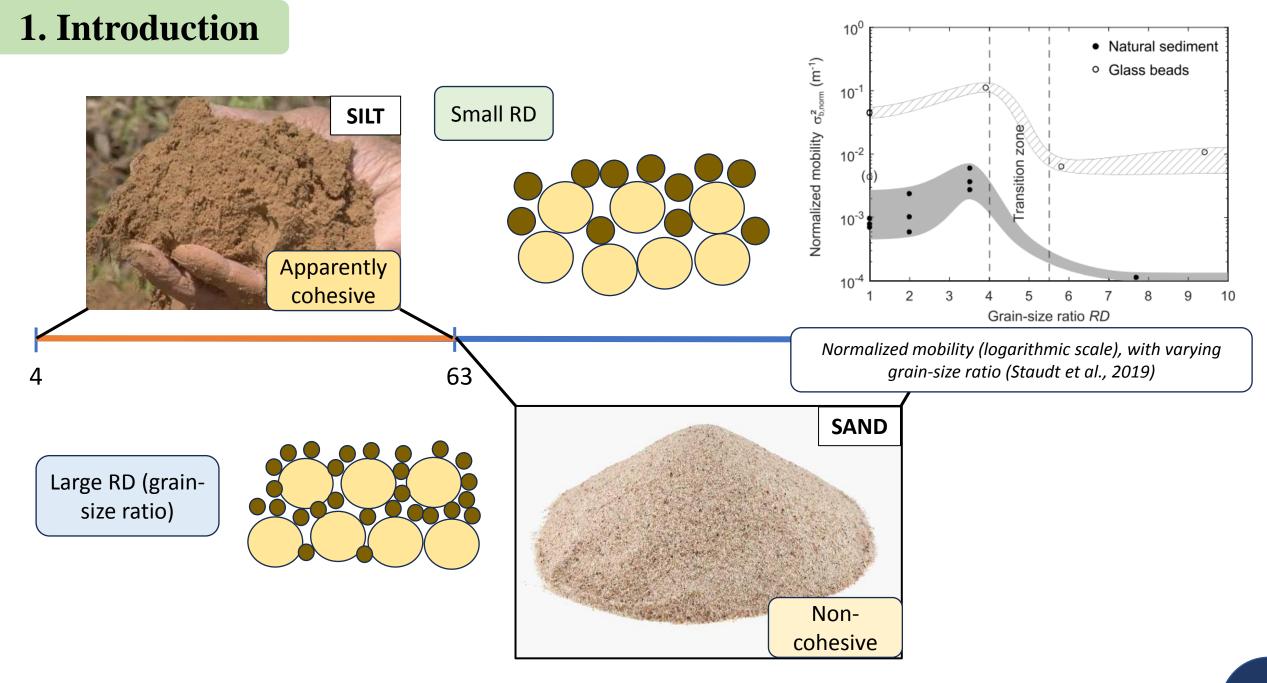
106°15'0"E

05°0'0"

105°0'0"E

1. Introduction





The second preparatory experiment in the small-scale oscillatory tunnel, also known as the **Aberdeen Mini Tunnel** (AMT). For our experiment, we will use solid glass beads as the material. These beads are round and well-sorted.

PISTON

Observations, results, and experiences obtained from the preparatory experiments will inform configurations conditions, bed and the procedures of the main AOFT experiments. MWL 🔻 **Sediment fractions** Silt content Velocity Mixtures (%) conditions (m/s)Sand ($D_{50} \approx 150 \, \mu m$) Sa 0 0.2 0.3 253 cm 1 3.5 cm 0.2 Sand ($D_{50} \approx 150 \, \mu m$) SaCs20 20 235 cm and Coarse silt ($D_{50} \approx$ 0.3 50 µm) SEDIMENT BED SaCs40 40 0.2 WINDOWS 0.3 Sand ($D_{50} \approx 150 \, \mu m$) SaMs20 20 0.2 15 cm 7 cm END RAMPS and Medium silt ($D_{50} \approx$ 3.5 cm **1**3.5 cm 0.3 25 µm) 15 cm 25 cm SaMs40 40 0.2 The Aberdeen Mini Tunnel (AMT) at the University of Aberdeen. 0.3

Water

pump

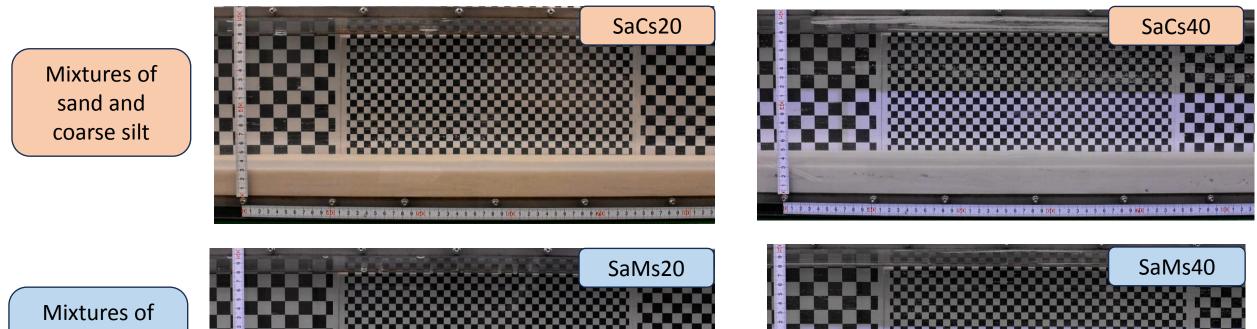
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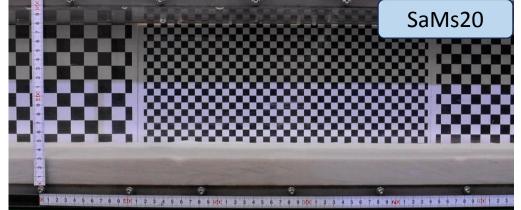
Flow condition: T = 3s, $U_{max} = 0.2$ m/s

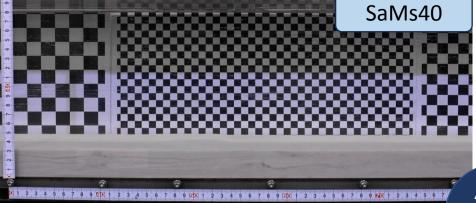
- Bedform development over approximately 450 flow cycles (~23 mins) for different beds was observed.
- For the mixed bed containing **40% silt**: **Higher** suspended sediment • concentration (SSC) and ripples developed more slowly compared to others.





sand and medium silt

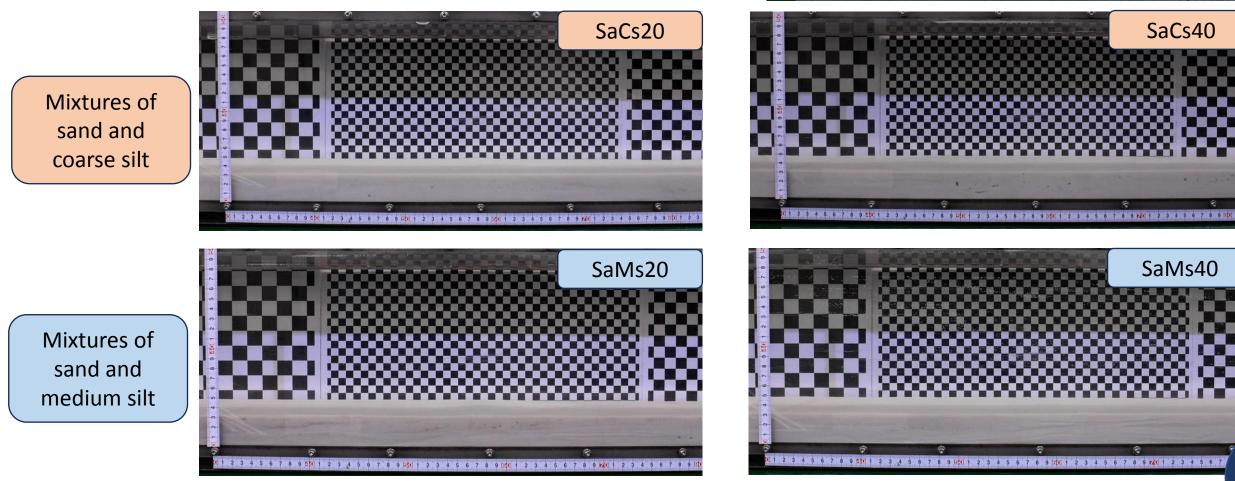




5

Flow condition: T = 3s, $U_{max} = 0.3$ m/s

- Bedform development over approximately 450 flow cycles (~23 minutes) was observed for different beds.
- Under higher-velocity conditions, a high suspended concentration was easily seen during the initial cycles.
- **Ripples were forming faster** compared to the 0.2 m/s condition.

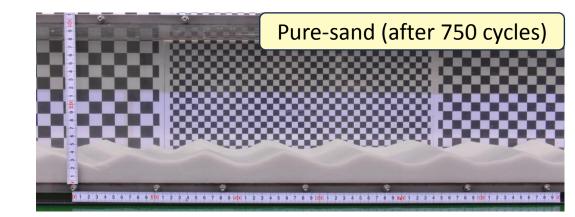


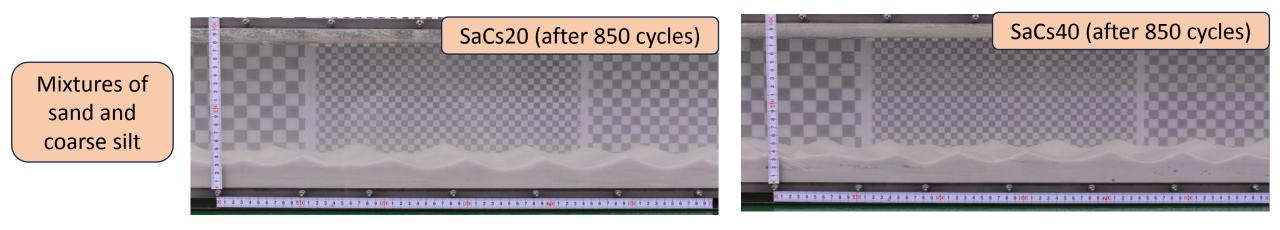
Pure-sand

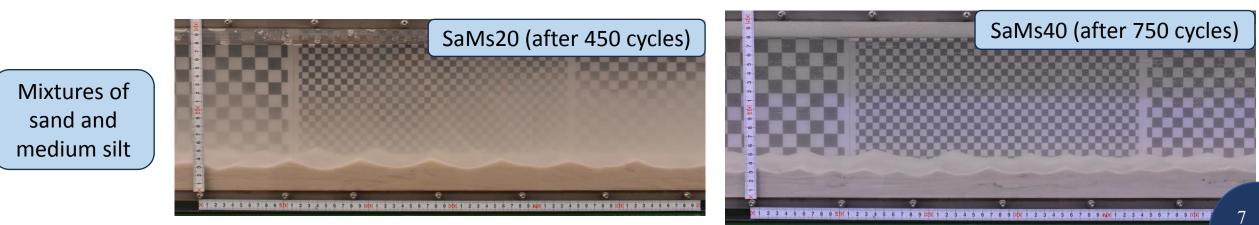
6

Flow condition: T = 3 s, $U_{max} = 0.2$ m/s

- Bedform of different beds at **the end** of each experiment.
- 3D ripples were observed in mixed beds.
- Ripples of 40%-silt-content mixtures had smaller scales compared to 20% mixtures and pure-sand bed.



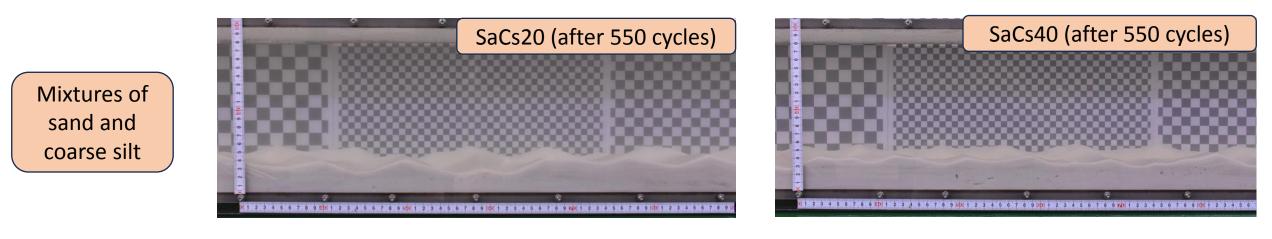




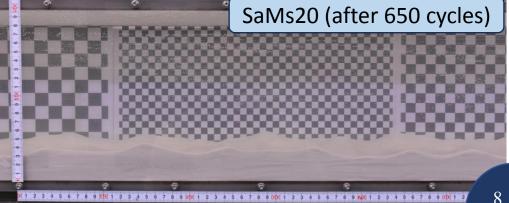
Flow condition: T = 3 s, $U_{max} = 0.3$ m/s

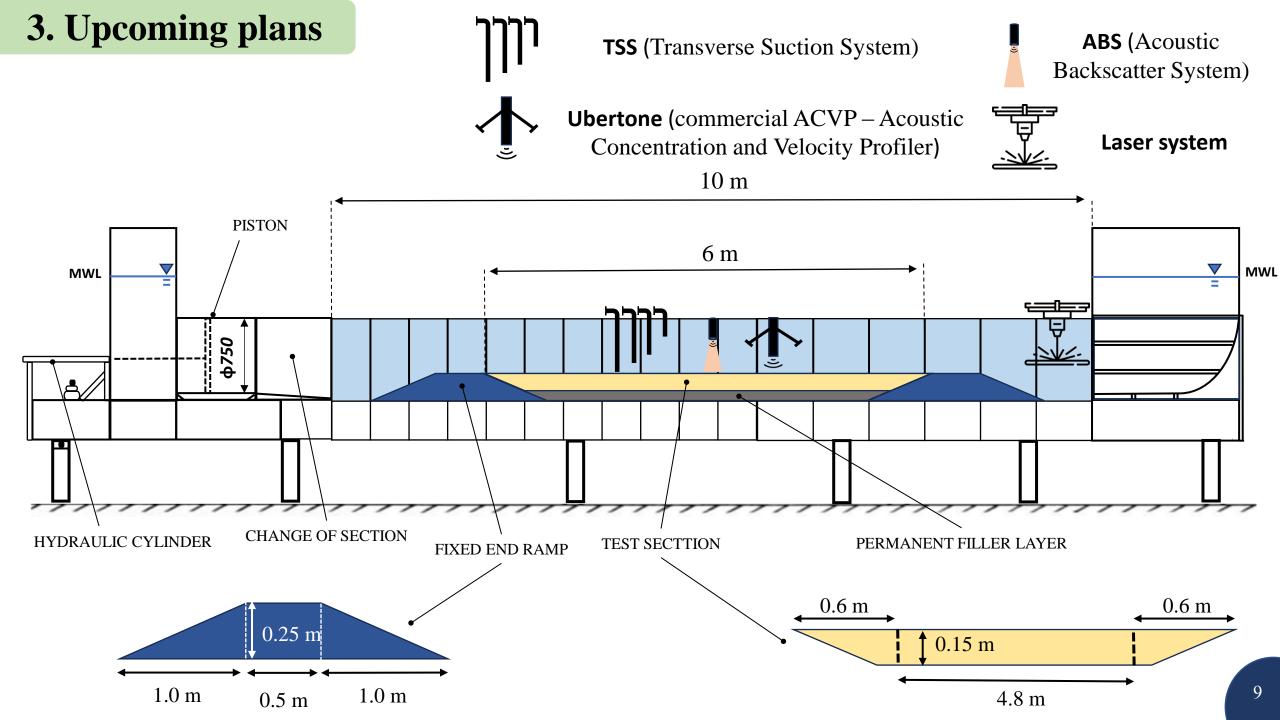
Bedform of different beds at the end of each experiment. Compared to 0.2-m/s condition, 3D ripples developed faster with higher velocity. Suspended sediment concentrations were also higher (visual observations).











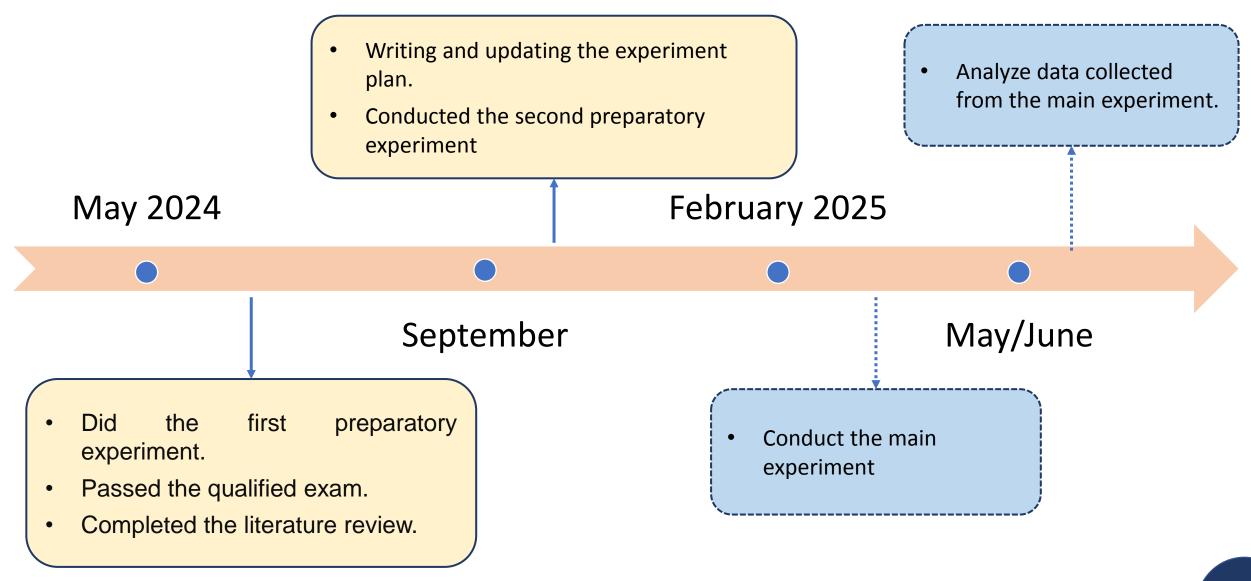
4. Conclusion

We have a nice preparatory experiment.

We are updating the experiment plan and preparing to have the main experiment conducted in Feb 2025.

THANK YOU FOR YOUR ATTENTION

Source: <u>freepik</u>

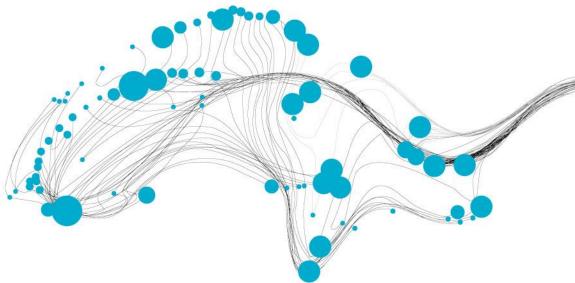


Ideas for SEDIMARE presentation:

Give introduction about sand-silt mixtures, effect of RD => fillingpores effect and hiding-exposure effect (coarser silt can enhance mobility of sand, and oppositely, fine enough silt can decrease mobility of the bed). Mention finding of Lange 2024 using the transport stage. Mention the difference between study of Lange and Staudt and Barzke

SEDIMARE DC8

Morphodynamics of Breach Growth and Bank Erosion using Laboratory Experiments



By Siyuan Wang

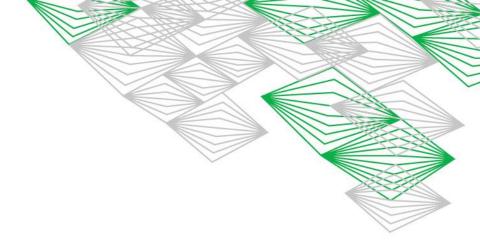
2024-11-06







CONTENT



- Introduction
- Bank Erosion and Dam Breaching
- Dam Breaching Experiment Design
- Previous Dam Breaching Experiments in UCLouvian

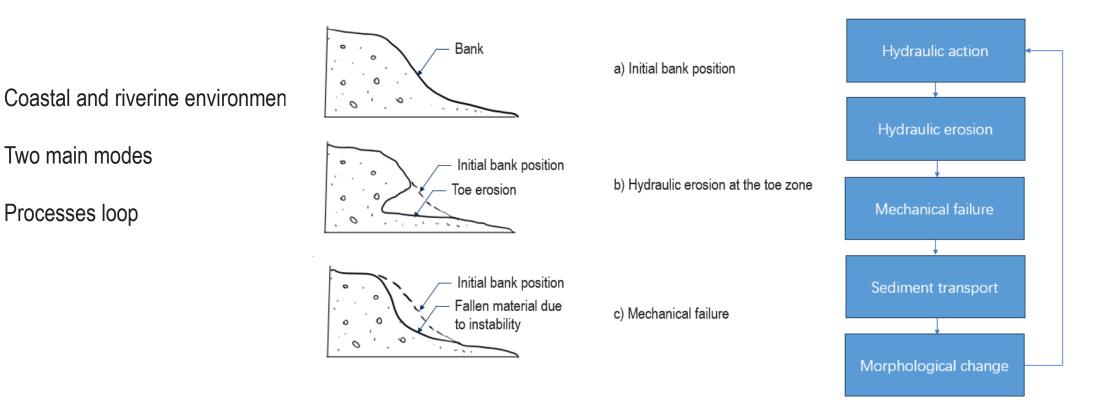
Introduction

-

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-

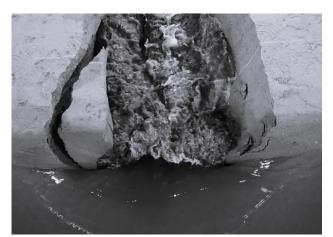
Bank erosion refers to the process of **removing material** from a bank which is the land alongside **a body of water** in geography.



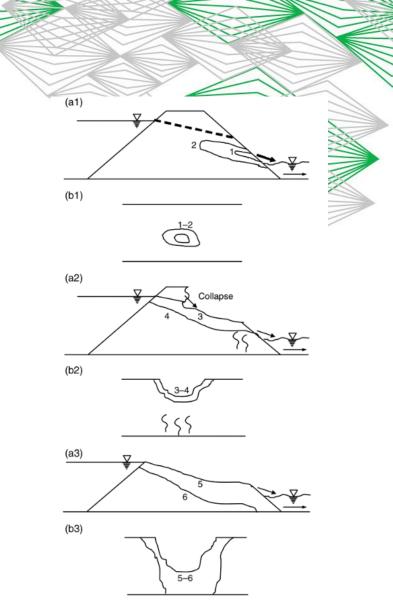
Bank Erosion and Dam Breaching

Similarities

- Environment conditions
- Hydraulic forces (in coastal setting)
- Processes



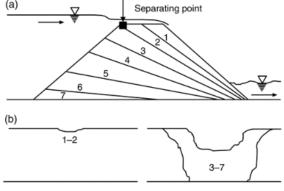
Ebrahimi, 2024

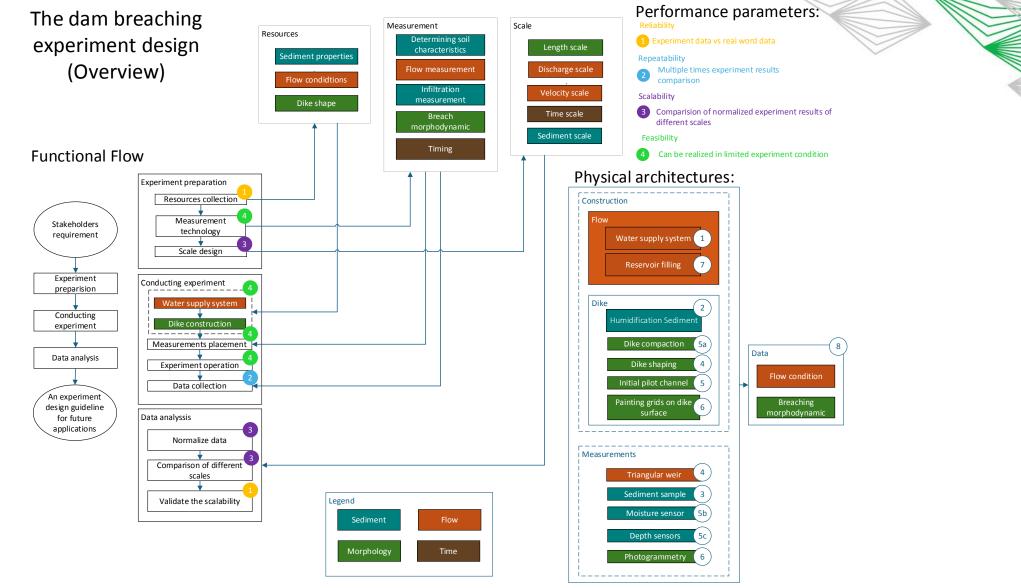


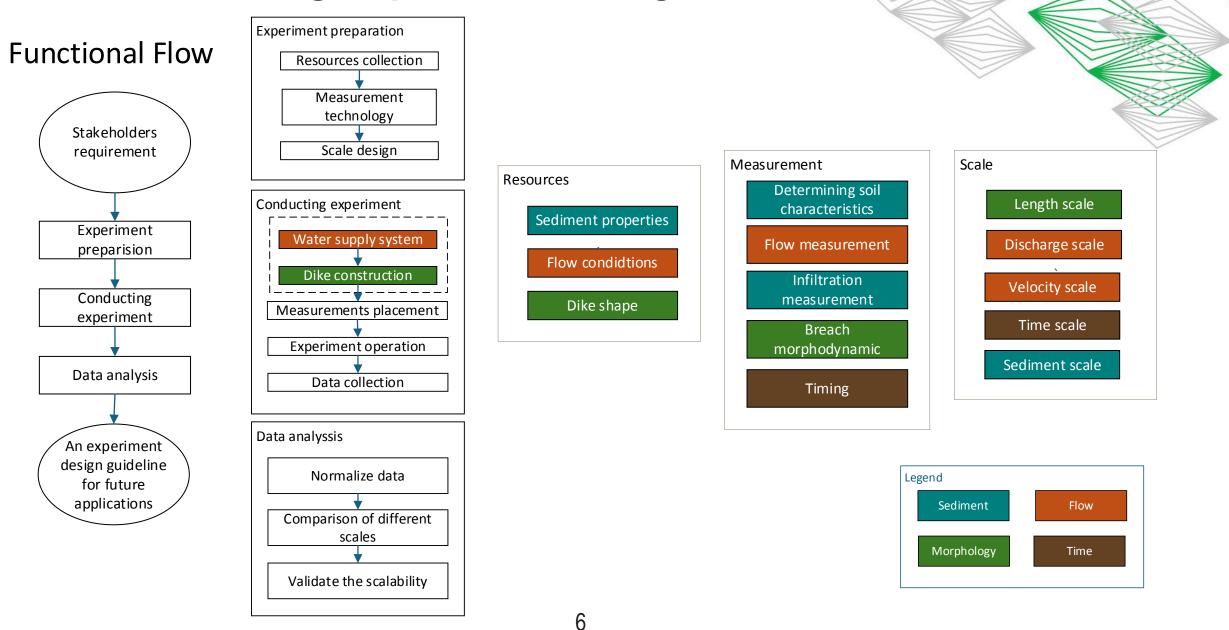
Zhang et al., 2016

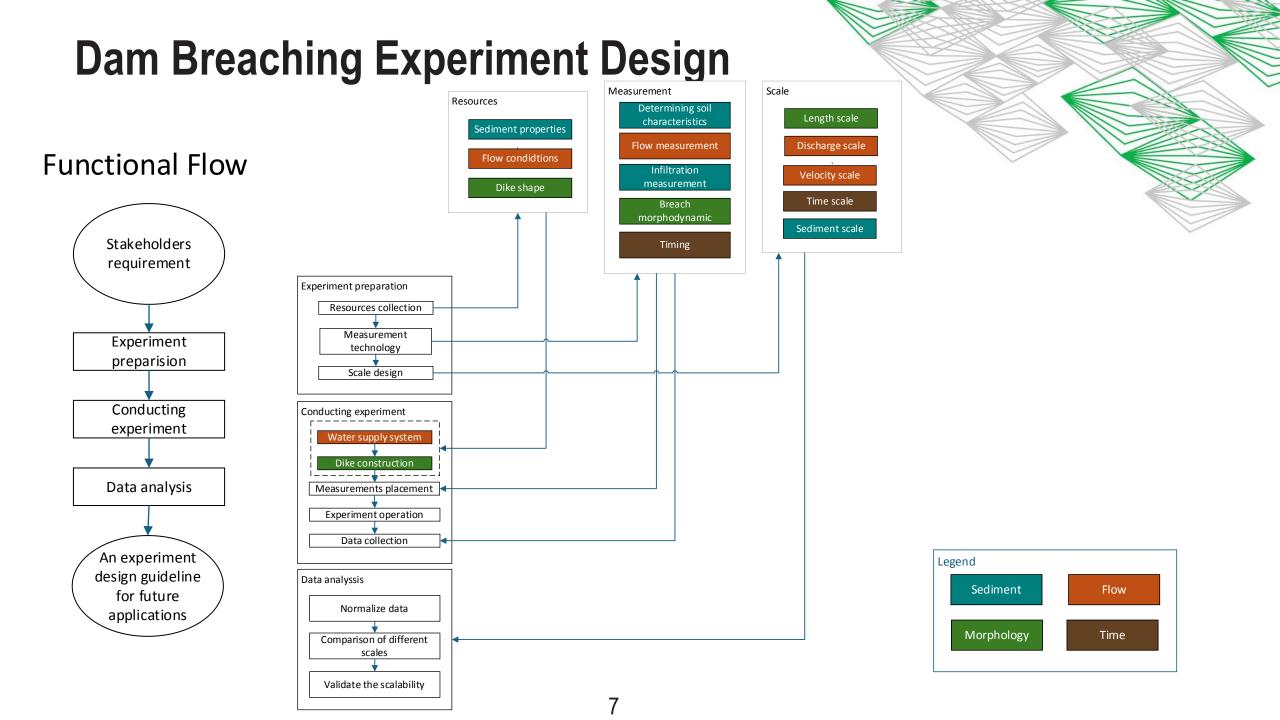
Differences

- Initiations
- Hydraulic forces (in riverine setting)
- Processes period

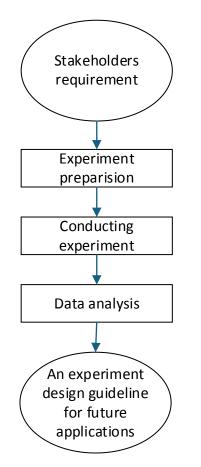


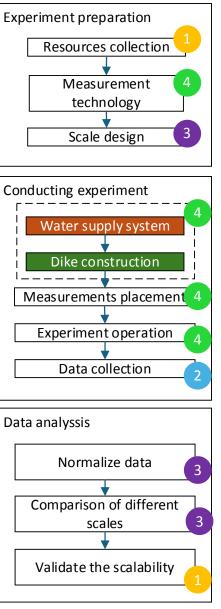


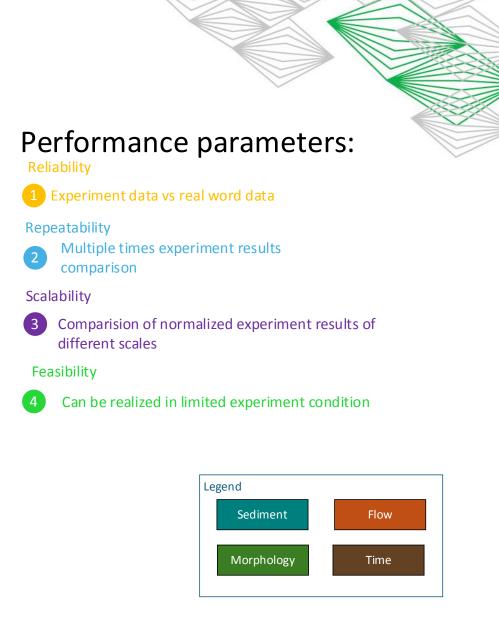




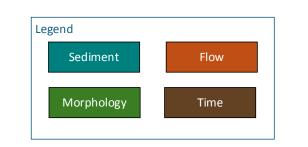
Functional Flow







Physical architectures: Construction Flow Water supply system 1 Reservoir filling 7 Dike 2 Humidification Sediment Dike compaction 5a Data Dike shaping Flow condition Initial pilot channel 5 Painting grids on dike Breaching 6 surface morphodynamic *i* Measurements 4 Triangular weir Sediment sample 3 Moisture sensor 5b Depth sensors 5c Photogrammetry 6 9



PREVIOUS DAM BREACHING EXPERIMENTS IN UCLOUVIAN

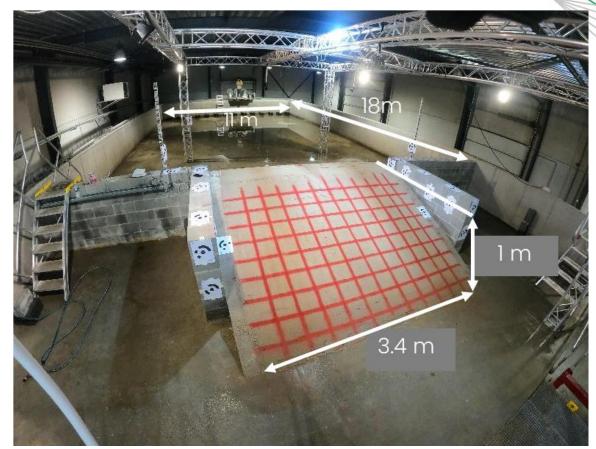
Experiment Scales

Delpierre, 2014

H_s=0.2 m Coarse sand d₅₀=1.7 mm H_s=0.2 m Fine sand d₅₀=0.7 mm

Small scale

Ebrahimi, 2024



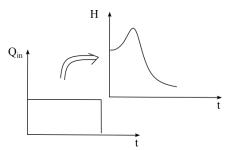
Medium scale

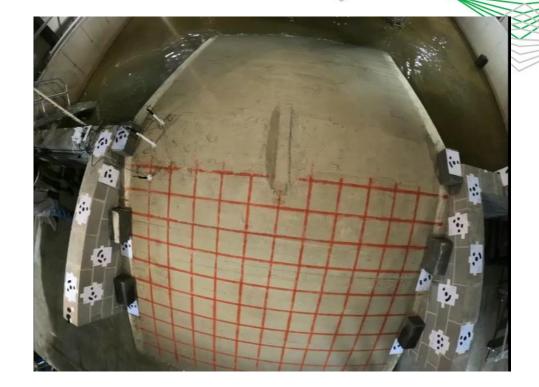
PREVIOUS DAM BREACHING EXPERIMENTS IN UCLOUVIAN

Water supply system

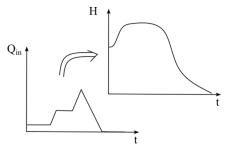


Constant inflow



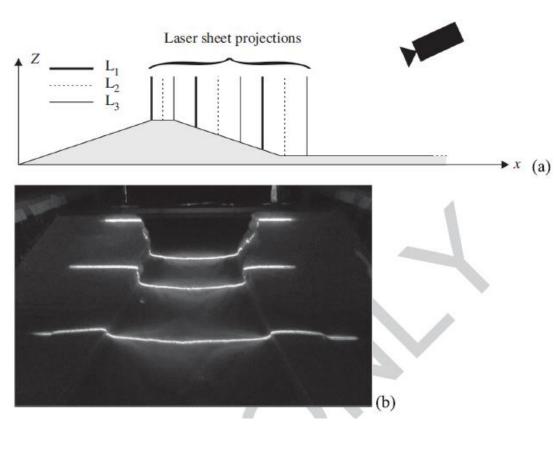


Constant water level Or constant inflow



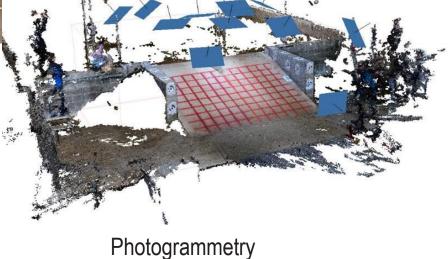
PREVIOUS DAM BREACHING EXPERIMENTS IN UCLOUVIAN

Measurement techniques



Laser-sheet and camera measurements









SEDIMARE

ent Transport and Morphodynamics in Marine and Coastal Waters with

Characterization of stratification and near-bed dense layers in high-density sediment-laden flow

D.C 9- Eloah Rosas

Promoter

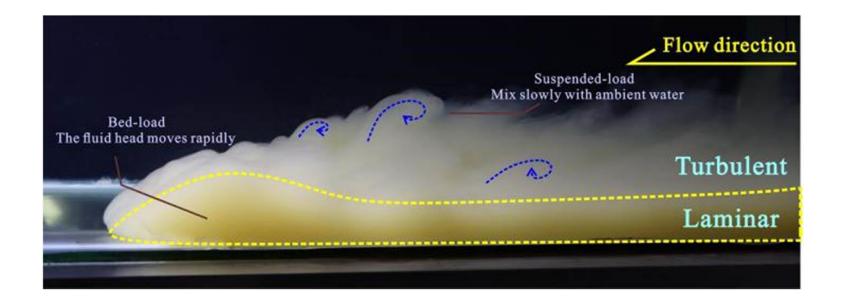
Benoit Spinewine

Sandra Soares

Sediment-laden flow is characterized by unsteady, highly dense sediment concentrations and vertical stratification.

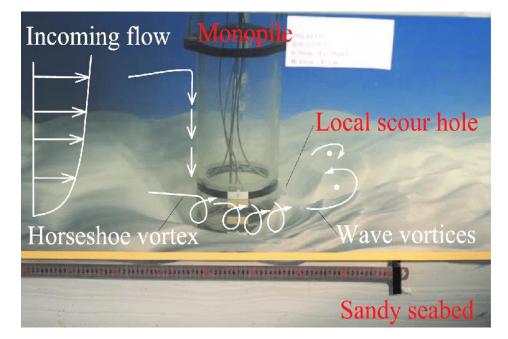
High-density sediment-laden flows is common in natural systems like rivers, deltas and coastal environments, and it can be initiated by natural or anthropogenic factors in the marine environment, such as:

Tsunamis, Earthquakes, Storms waves, Submarine Landslides, Dredging.



Studying the stratification and near-bed dense layers in high-density sediment-laden flows is crucial for several engineering applications, particularly in coastal and offshore infrastructure:

- 1. Wave-induced scour hole development around infrastructure such as wind turbines;
- 2. Storm-induced sheet flows on beaches causing dune instabilities and coastline retreat;
- 3. Turbidity currents, debris flows or submarine landslides posing a significant threat to pipelines or large submarine power cables



Flow pattern involving scour around monopile. Gao, F., & Qi, W. (2022).

Studying the stratification and near-bed dense layers in high-density sediment-laden flows is crucial for several engineering applications, particularly in coastal and offshore infrastructure:

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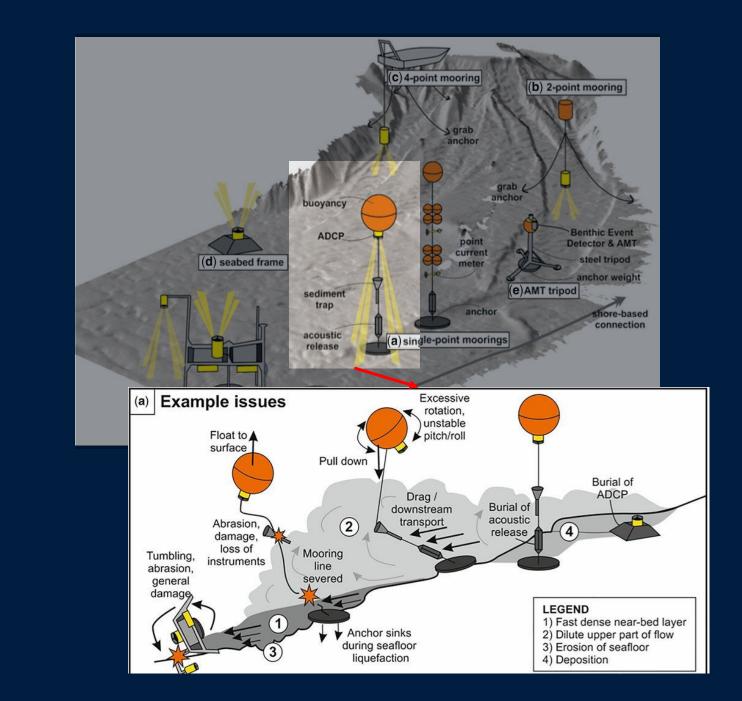




Broken cables due to the turbidity currents and submarine landslides in Kaoping Canyon subsquently two earthquakes. Hsu et al. (2008).

Studying these flows involves a combination of **field observations,** laboratory experiments,

and numerical simulations



Studying these flows involves a combination of field observations, **laboratory experiments**, and **numerical simulations**





Bagnold	Capart	Armanini	
Sumer	Spinewine	Francarollo	
Close Duct Flume	Recirculating tilting Flume (6 m x 0.25 m x 0.70 m) Vertical gate at the middle section		
Pressure sensors CCM Prandtl tube	High Speed CCD Camera Laser light sheet		
Coefficient Friction	 Comprehensive characterization of flow regimes 	 Stratified flow with sub-layers Transitions between sub-layer- based Stokes number 	
Bedload power law velocity in the sheet flow	 Detailed profiles of velocity and granular concentration at different depths Enhanced shallow water equations with vertical details 	 Richardson number constant in collisional. 	

Matousek

Recirculating Tilting Flume (8 m x 0.2 m x 0.27m)

High Speed camera Laser UVP (4 MHz); Prandl tube Ultrasonic water level

- New linear equation
- Measured local velocity and granular temperature in bedload.
- layered structure in sediment with three sublayers based on the dominant transport method
- Transport model with position of the interfaces of sublayers.

-fugro

Bagnold	Capart	Armanini	Matousek
Sumer	Spinewine	Francarollo	
Close Duct Flume	Recirculating tilting Flume (6 m x 0.25 m x 0.70 m) Vertical gate at the middle section		
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fugro

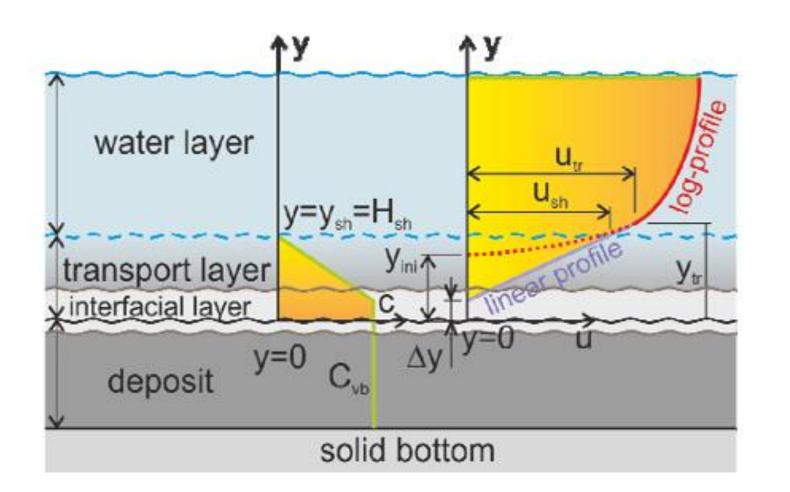
Bagnold	Capart	Armanini	Matousek
Sumer	Spinewine	Francarollo	
Close Duct Flume	Recirculating tilting Flume (6 m x 0.25 m x 0.70 m) Vertical gate at the middle section	Recirculating tilting Flume Conveyor belt recirculating sediment	
Pressure sensors CCM Prandtl tube	High Speed CCD Camera Laser light sheet	High Speed Camera	UVP (4 MHz); Prandl tube Ultrasonic water level
Coefficient Friction Bedload power law velocity in the sheet flow	 Comprehensive characterization of flow regimes Detailed profiles of velocity and granular concentration at different depths Enhanced shallow water 	 Stratified flow with sub-layers Transitions between sub-layer- based Stokes number Richardson number constant in collisional. 	 New linear equation Measured local velocity and granular temperature in bedlo layered structure in sediment of three sublayers based on the dominant transport method
	equations with vertical details	 	Transport model with position the interfaces of sublayers

the interfaces of sublayers.

-**FUGRO**

Bagnold	Capart	Armanini	Matousek
Sumer	Spinewine	Francarollo	Recirculating Tilting Flume
Close Duct Flume	Recirculating tilting Flume (6 m x 0.25 m x 0.70 m) Vertical gate at the middle section	Recirculating tilting Flume Conveyor belt recirculating sediment Dams	(8 m x 0.2 m x 0.27m) High Speed camera Laser
Pressure sensors CCM Prandtl tube	High Speed CCD Camera Laser light sheet	High Speed CCD Camera	UVP (4 MHz); Prandl tube Ultrasonic water level
Coefficient Friction Bedload power law velocity in the sheet flow	 Comprehensive characterization of flow regimes Detailed profiles of velocity and granular concentration at different depths Enhanced shallow water equations with vertical details 	 Stratified flow with sub-layers Transitions between sub-layer- based Stokes number Richardson number constant in collisional. 	 New linear equation Measured local velocity and granular temperature in bedload. layered structure in sediment with three sublayers based on the dominant transport method Transport model with position of the interfaces of sublayers.

-fugro



Matousek

Recirculating Tilting Flume (8 m x 0.2 m x 0.27m)

High Speed camera Laser UVP (4 MHz); Prandl tube Ultrasonic water level

- New linear equation
- Measured local velocity and granular temperature in bedload.
- layered structure in sediment with three sublayers based on the dominant transport method
- Transport model with position of the interfaces of sublayers.

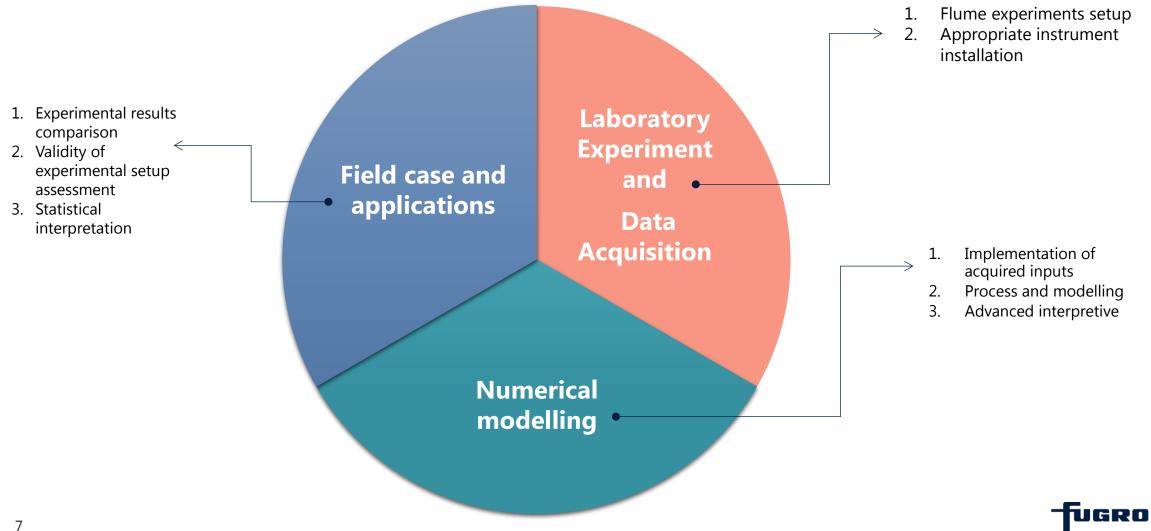
fugro



- 1. Propose a **multi-layer shallow-water modelling** framework, and propose **practical parametrizations of interlayer exchange processes** that may capture the behaviour of dense near-bed layers;
- 2. Calibrate these relations with well-controlled small-scale laboratory experiments
- 3. Confront the results with documented case studies or direct field-scale observations

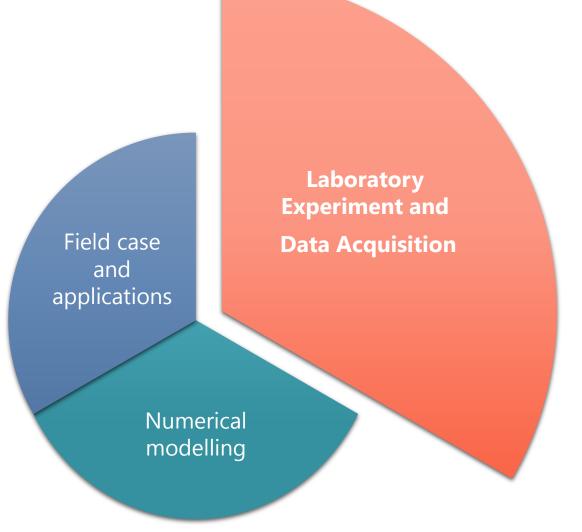


Three-tier approach



Three-tier approach

- **Setup 1**: "Submerged Granular Column Collapse" experiments at UCL
- Experimental setup readily available
- Detailed measurements of sediment velocities in transport layer
- **Setup 2:** "modular tilting / adverse slope flume " at Fugro or UCLouvain
- Configuration 1: rapid tilting of flume with a sediment layer initially at rest:
- Configuration 2: steady liquid-granular flow on adverse slope (water goes up, sediment goes down)



fugro

Thorpe Flume

Channel

- Length 1830 mm
- Width 100 mm
- Height 30 mm
- Glass walls

Flume

- Tilting using Eletric Jack
- Tilting angle 4.13 in 0.45 s

Duponcheel et al. 2017





Modifications Thorpe flume

New Channel

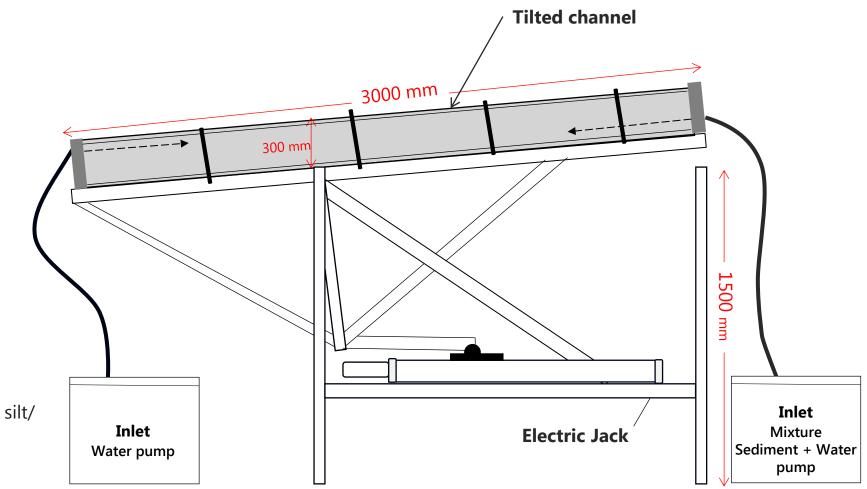
- Length 3000 mm
- Width 100 mm
- Height 300 mm
- Transparent Plexiglas sidewalls

Tilting

- Tilting using Electric Jack
- + 20 °

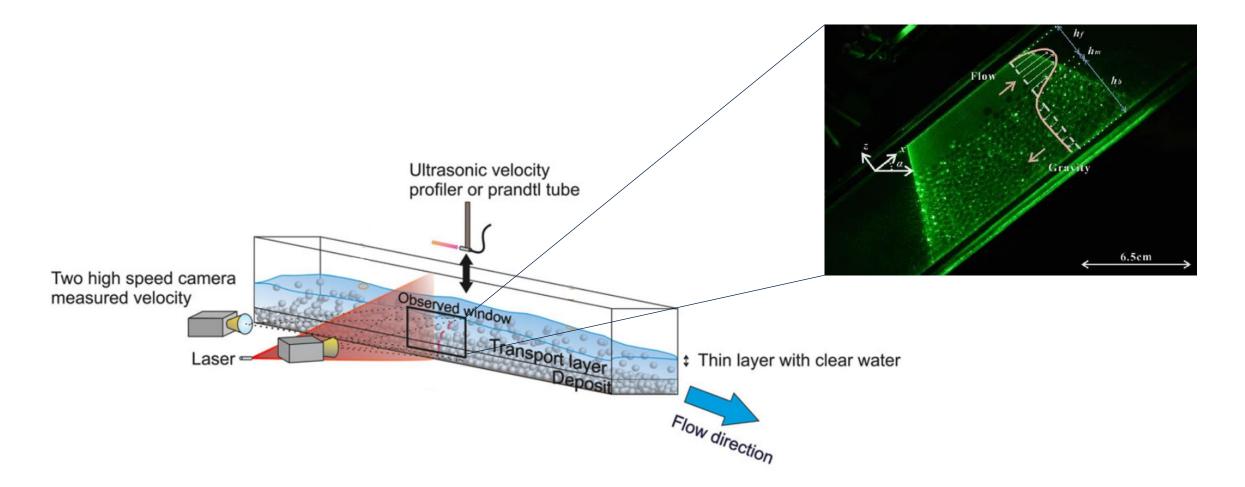
Sediments Type

- Natural sediment (fine sediment silt/ clay)
- Plastic grains (3 4 mm)



-fugro

Measuring Procedures



TUGRO



SEDIMARE

liment Transport and Morphodynamics in Marine and C

Characterization of stratification and near-bed dense layers in high-density sediment-laden flow

Eloah Rosas e.rosas@fugro.com





Overtopping Breakwater For Energy Conversion (OBREC)

Saeed Osouli

Supervisors:

Prof. Matteo Postacchini - UNIVPM

DR. Ivan Sabbioni - MAC

Prof. Maurizio Brocchini - UNIVPM

SEDIMARE

2nd Network Training School: Experimental and Practical Modelling of Sediment Transport and Coastal Morphology

University of Twente/Deltares

5-8 November 2024





Outline

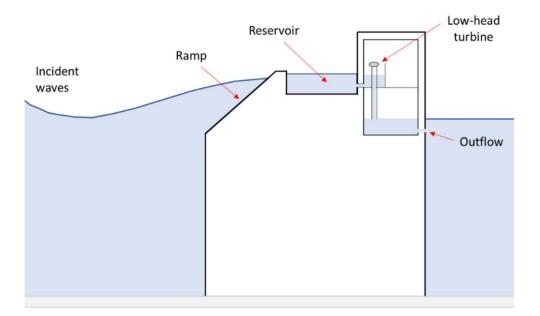


 Introductio
 Comparing and Selecting Data
 Shallow water Solver
 CFD Module
 Conclusion



Motivations

- Advantages that characterize Wave energy:
- □ The high energy density.
- □ The easy prediction of the wave characteristics.
- $\hfill\square$ The reduced energy loss
- Drawbacks:
- □ The high variability of the wave characteristic.
- □ WECs are exposed to large environmental forces.
- □ High production costs compared.



Scheme of the OTD integrated into a vertical caisson breakwater

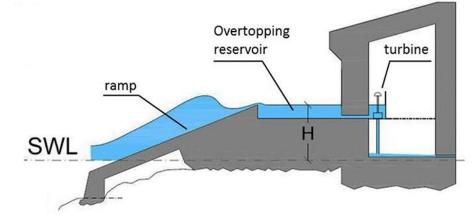
MENTUCCIAldo

Literature review

E

CONCRETE





ARMOUR LAYER (D= 2,695 m ARTIFICIAL ROCKS - OUTSIDE.

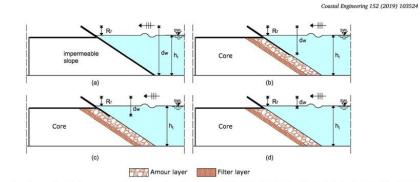


Fig. 3. OBREC configurations analyzed in the present experimental campaign: a) $d_w = 0.274$ m ($d_w^* = 1.0$); b) $d_w 0.166$ m ($d_w^* = 0.6$); c) $d_w = 0.113$ m ($d_w^* = 0.4$); d) $d_w = 0.064$ m ($d_w^* = 0.2$). R_r indicates the crest free-board and h_i indicates the water depth at the model toe.

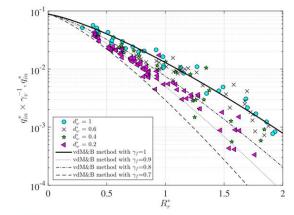
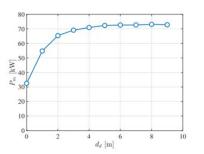
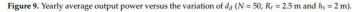
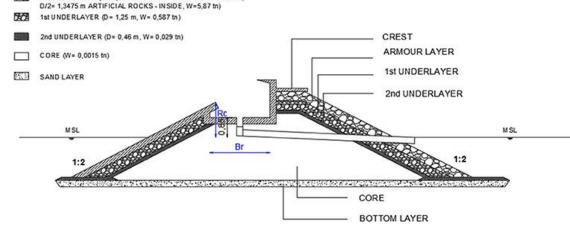


Fig. 8. Comparison of the average wave overtopping rates measured for all configurations tested in the present tests and those estimated by the prediction method of van der Meer and Bruce (2013) adopting four different values of γ_f . The experimental data were corrected using the coefficient γ_V .



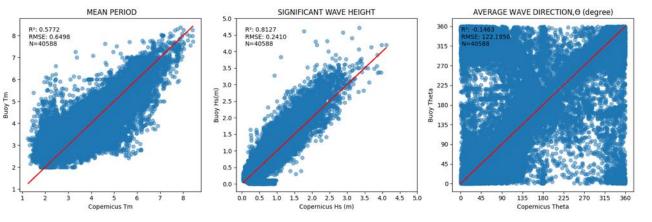
• OBREC





C. Iuppa, et al.

Observed and Model data



Comparison of wave parameters from buoy and Copernicus.



Buoy N Copernicus N 330 3.00% 3.00% 4.50% 4.50% 6.00% 6.00% 7.50% 7.50% 9.00% 9.00% Wave height (m) Wave height (m) [0.1:1.0) [1.0:2.0] [1.0:2.0] [2.0:3.0] [2.0:3.0) [3.0 : 4.0) >4.0 [3.0 : 4.0) >4.0 ς

Wave Rose

MENTUCCIAldo



Wave selection



Copernicus and Buoy

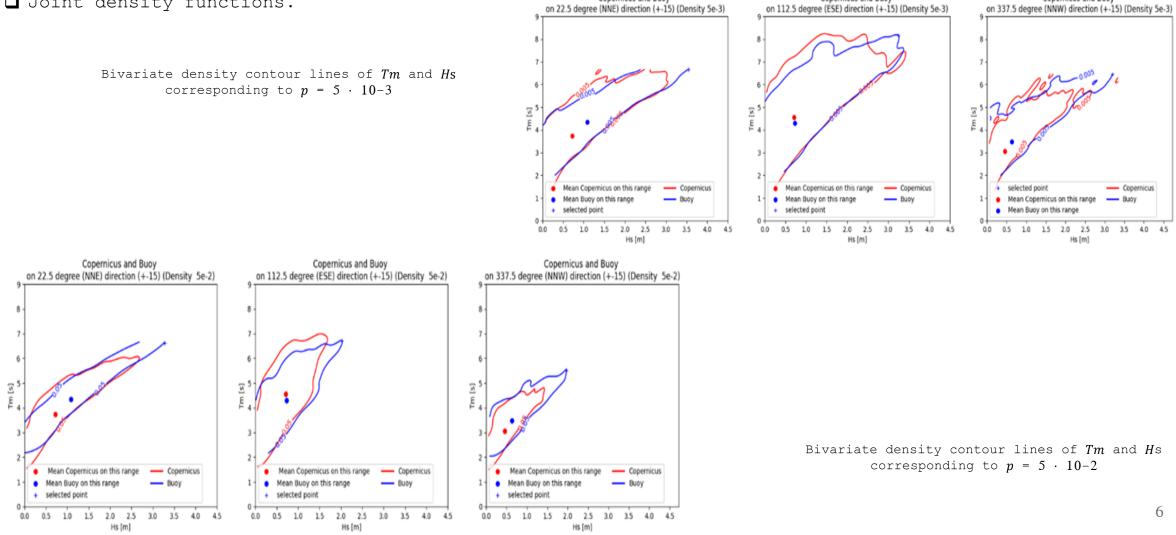
 \Box min, max and average wave Heights based on 10-year average.

Extreme Wave Conditions (EWC), with a 100-year return period, and Average Wave Conditions (AWC).

Copernicus and Buoy

Copernicus and Buoy

□ Joint density functions.





Wave characteristics at the Buoy Location and numerical MENTUCCIALdo domain

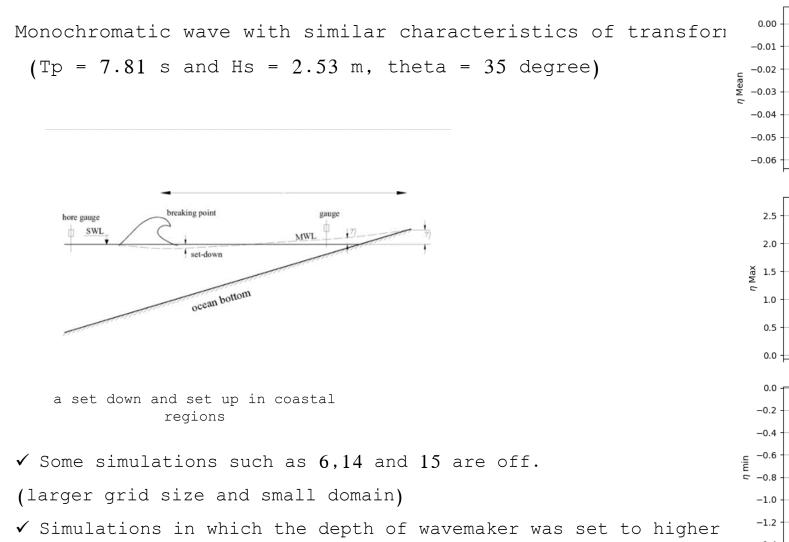


90.0° 121.0°			NNE	ESE	NNW			NNE	ESE	NNW
81.5° 171.5°		Hs (m)	1.09		0.63		Hs	0.90		0.59
36.5°NW NHE		Tm (s)	4.35		3.49		(m)	0.90		0.57
W E Adriatic Sea		α (Respect to	22.5		337.5	Mean	Tp	5.13		4.12
Sdriatic Sea		the North) [°]					(s)			
Hariate Sea 31.0°	Mean	αp 0 (Respect to					ap,h1 [°]	56.00		35.00
90.0°		orthogonal line	01 50		26 50		l J Hs			
		to boundary of the numerical	81.50		36.50		(m)	2.53		1.67
		domain) [°]					("") Tp			
		Hs (m)	3.28		1.97	Density 0.05 (s)		7.81		6.51
		Tm (s)	6.62		5.52		ap,h1			
AT HE DOLLAR STREET, SAN DE		α (Respect to					[°]	35.00		27.00
		the North) [°]	22.5		337.5		Hs			
	Density 0.05	$\alpha p0$ (Respect to					NNE			NNW
Inclination due to the boundary.		orthogonal line					12			1A
boundary.		to boundary of	81.50		36.50		1			17
		the numerical					11			/ / h
		domain) [°]					¥ /		У	
Phase-averaged (ROMS-SWAN)		Hs (m)	3.54		3.2		14/	36.5°		/
\square Analytical model (Goda 201		Tm (s)	6.67		6.46	81.5° /	Y /		$\setminus \vee /$	0
Analytical model (Goda 201		α (Respect to	22.5		337.5	01.5			\mathbf{V}	
		the North) [°]				1×		<	35.0°	
	Density 0.005	αp 0(Respect to				-56.0°		\vee	7 /	
		orthogonal line	01 50		36.50			, Y		NI.
	Wave charact	teristics at the Buo the numerical	y Local io	on.		¥/.,	driatic Sea	V /		-11
		domain) [°]				1/		\vee /		
								\sim		/W

Transformation and Deformation of Sea Waves.

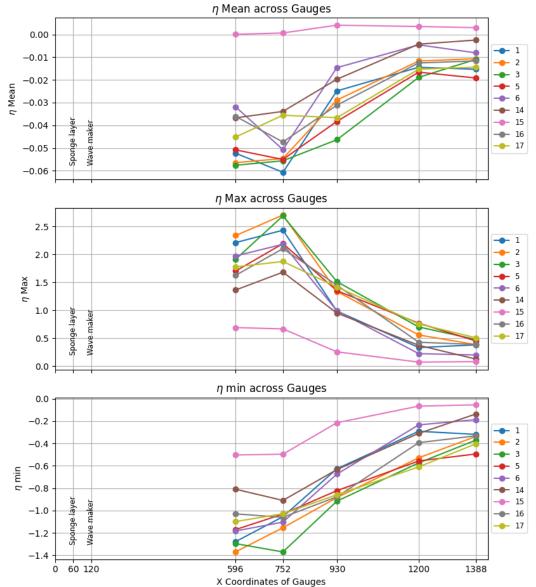






value provides much more energy. (3 and 17)

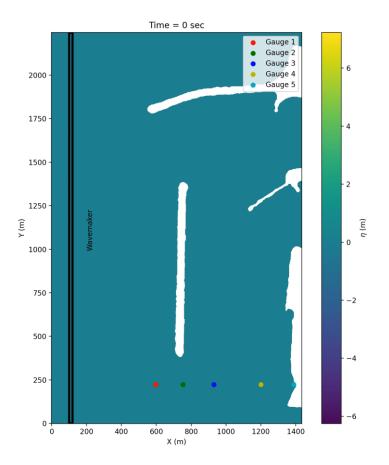
Configuration 3 was selected for next simulations.

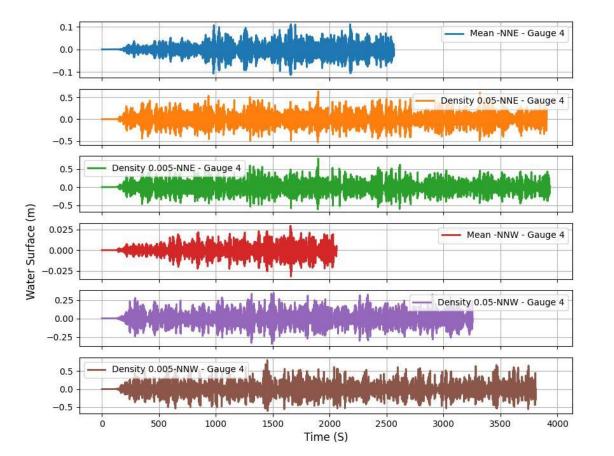






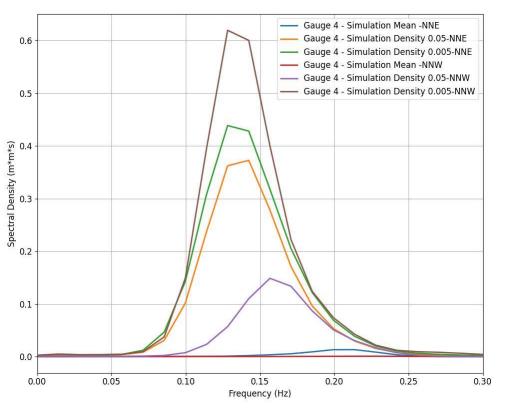
Simulation time : 500 waves for each scenarios.

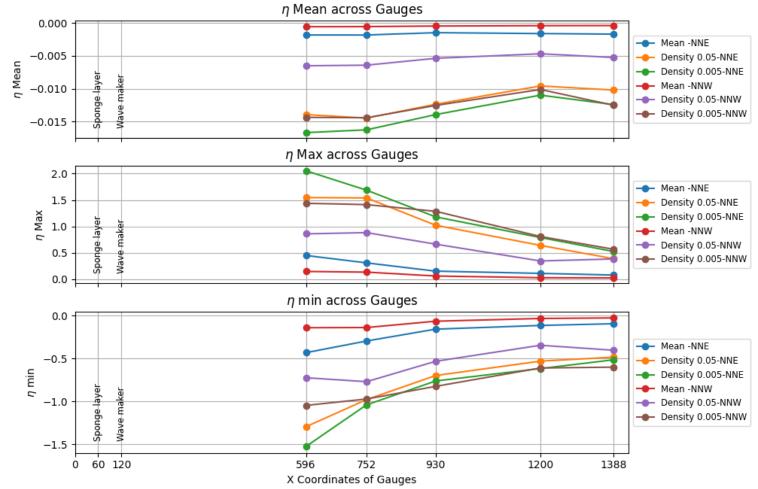












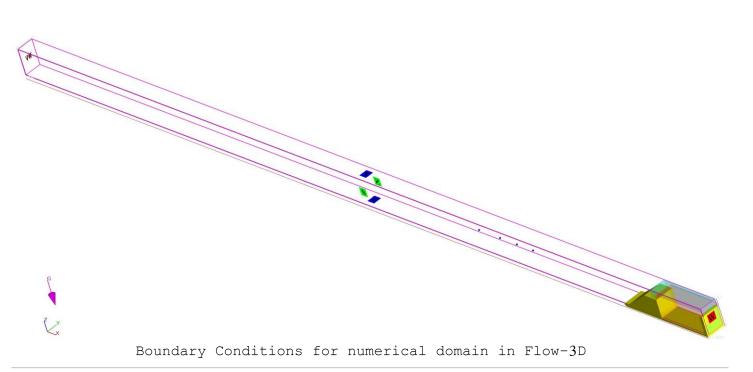


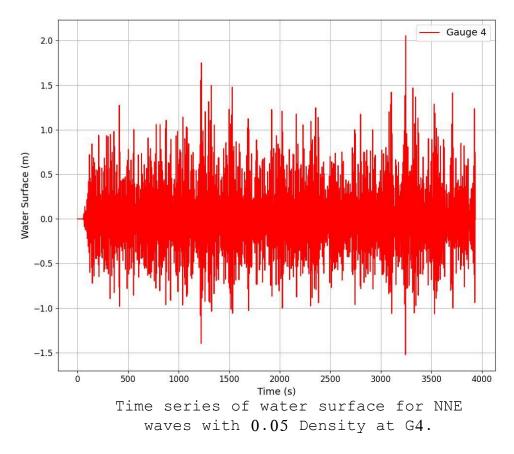
FLOW-3D



11

- Free surfaces are modeled with the Volume of Fluid (VOF) technique.
- Utilizing Navier-Stokes-based numerical simulation.
- Structured mesh can only be used.
- NNE waves with 0.05 density was used as Wave maker.
- Numerical domain was 227 m length, 4.4 m width and 7 m height.
- Water depth was set to 3 m. Simulation Time were 200 seconds.
- AutoCAD was used to create STL files as structures.



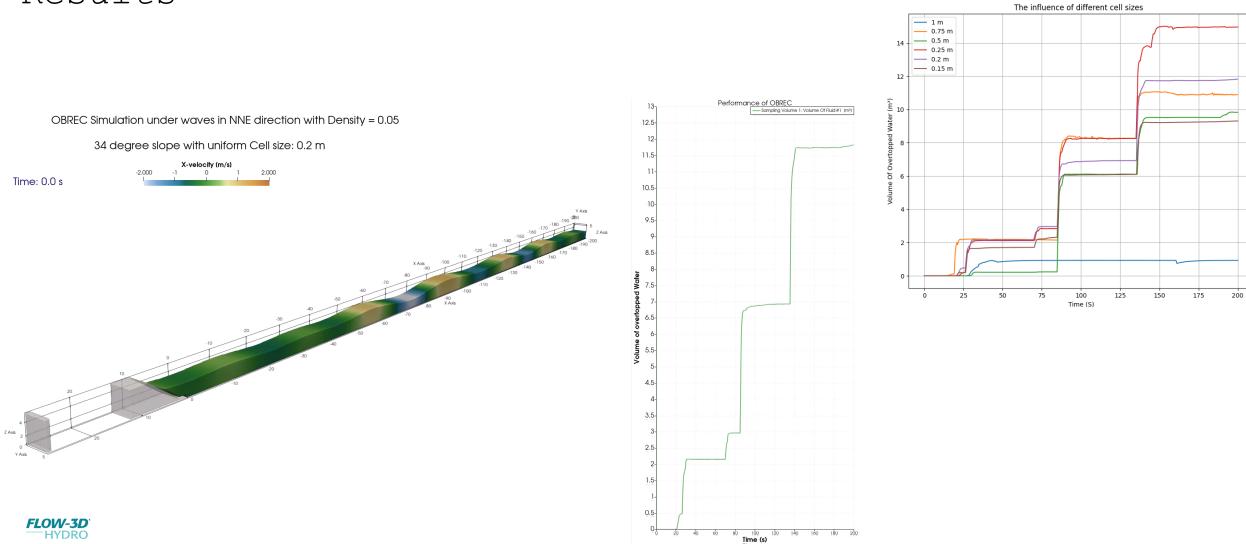




FLOW-3D



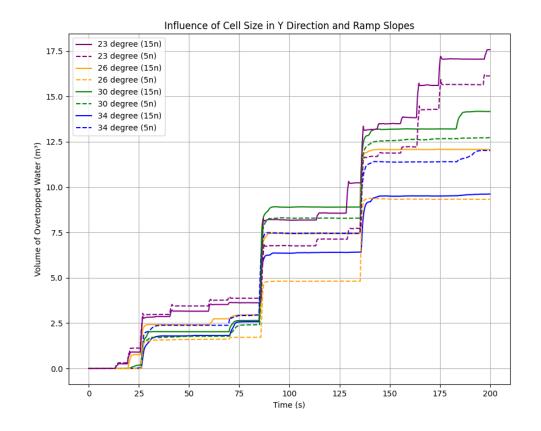
Results

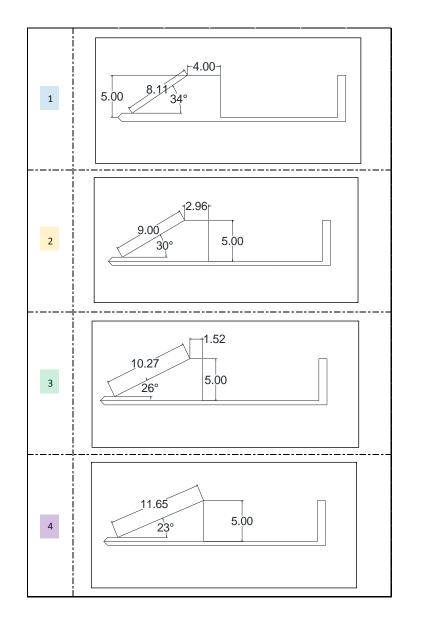


FLOW-3D- Effect of 2D approach and Ramp MENTUCCIALdo slopes

Results

All simulations have cell sizes in X and Z direction with 0.2 m and different cell number in Y direction.





A CHILDREN C

Conclusion and Discussion





□ PDF diagrams was used to wave selection.

Analytical model (Goda 2010) was used to transfer waves.

- □ FUNWAVE-TVD could represent set-down which is a non liner behavior of waves.
- □ FLOW-3D hydro was used to determine suitable grid size for evaluating various ramp shapes.

□ More simulation on ramp shapes, including validation.

Designing reservoir and convey system.

Running simulations for other scenarios.





(Wave-resolving model) has been used as a shallow water solver.

It is based on Boussinesq-type equations in which Reynolds equations are integrated over the water depth, so the vertical structure is not directly resolved but only modelled. (2DH models)

Bathymetry of the Port of Ancona for FUNWAVE: Combination of port bathymetry from port Authorities (CAD) and EMODnet data.

Volume conservation:

$$\eta_t + \nabla \cdot M = 0$$

 $\ensuremath{\mathsf{M}}$ is the horizontal volume flux.

The depth-averaged horizontal momentum:

$$u_{\alpha,t} + (u_{\alpha} \cdot \nabla)u_{\alpha} + g\nabla_{\eta} + V_1 + V_2 + V_3 + R = 0$$

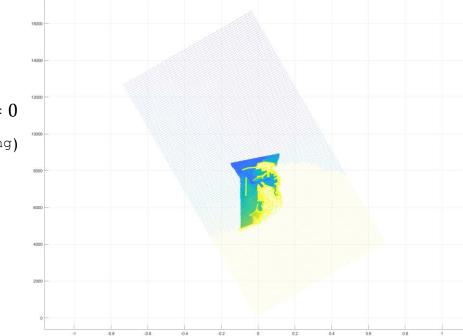
R: diffusive and dissipative terms (e.g. bottom friction, sub-grid lateral turbulent mixing)

 V_1 and V_2 are terms representing the dispersive Boussinesq terms (function of $z_{lpha})$.

 V_3 contribution of the order $O(\mu^2)$ (function of z_{α}). $\mu = \frac{h}{L}$

 $z_{\alpha} = \zeta h + \beta \eta$ that $\zeta and \beta$ are constants.

 u_{lpha} denotes the velocity at a reference elevation z = z_{lpha} .





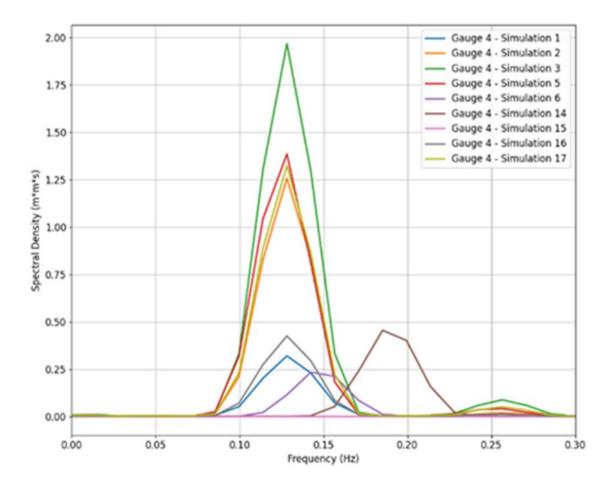


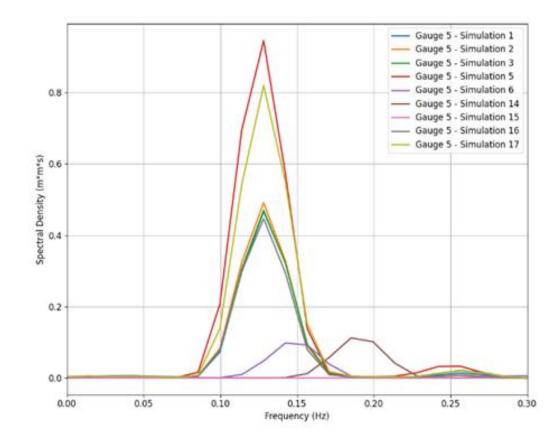
	al.,	al		Total	Stability		Smoothin	:		ndition	:	:	Mall or	:				ماسميرام	Total	Stability					Condition		:	Wall or	
	dx (m)	(m)	simulatio n time (s)		Condition	Smoothin	Smoothin g (Comman d)	height	WK and Spectru m	offshore		CFL and min depth	Wall or Periodic boundar y	VISCOSI TY_BRE AKING		dx (m)	dy (m)	simulati on time (s)		Conditio n	Smooth	Smoothi i ng (Comma nd)	hoight	: WK and	offshore platform	Froudec		Wall or Periodic boundar y	
1	2	2	1000	806,314	stable	7	smoothda ta2/sgola y/movme an	. 7	8 - monoch romatic	no	1	0.5/0.01	P.B	no	9	2	2	936	806,314	instable (it happens around 480 s)	7	smooth data/sg olay/sgo lay	5	12 - monoch romatic		1	0.5/0.01	P.B	n
2	2	2	1,000	806,314	stable	7	smoothda ta2/sgola y/movme an	7	12 - monoch romatic	no	1	0.5/0.01	P.B	no	10	1.5	1.5	399	1,432,629	around 150 s)	7	smooth data2/sg olay/mo vmean		12 - monoch romatic		1	0.5/0.01	P.B	r
3	2	2	1,000	806,314	stable	7	smoothda ta2/sgola y/movme an	7	16 - monoch romatic	no	1	0.5/0.01	P.B	no	11	1.5	1.5	185	1,432,629	instable (it happens around 122 s)	7	smooth data2/sg olay/mo vmean		16 - monoch romatic		1	0.5/0.01	P.B	n
4	2	2	1000	806,314	stabe but a instability sign could be seen in the domain around 810 s		smoothda ta2/sgola y/movme an	7	12 - monoch romatic	no	1	0.5/0.1	P.B	no	12	1.5	1.5	151	1,432,629	instable (it happens around 120 s)	7	smooth data2/sg olay/mo vmean		16 - monoch romatic		3	0.5/0.01	P.B	n
5	2	2	4292	806,314	stable	7	smoothda ta2/sgola y/movme an	7	12 - monoch romatic	no	1	0.5/0.01	P.B	yes/ Cbrk1 = 0.45 Cbrk2 = 0.35	13	2	2	7870	806,314	instable (it happens around 930 s)		smooth data2/sg olay/mo vmean		12 - monoch romatic	no	1	0.5/0.01	Wall boundar Y	n
6	2	2	7870	374,468	stable	7	smoothda ta2/sgola y/movme an	7	12 - monoch romatic	no	1	0.5/0.01	P.B	no- small domain	14	3	3	7870	358,771	stable	7	smooth data2/sg olay/mo vmean		12 - monoch romatic		1	0.5/0.01	P.B	r
7	2	2	467	806,314	instable (it happens around 330 s)	5	smoothda ta/sgolay /sgolay		12 - monoch romatic	no	1	0.5/0.01	P.B	no	15	5	5	7870	129,312	stable	7	smooth data2/sg olay/mo vmean		12 - monoch romatic	:	1	0.5/0.01	P.B	r
8	2	2	734	806,314	instable (it happens around 180 s)	5	smoothda ta/sgolay /sgolay		12 - monoch romatic	no	1	0.5/0.01	P.B	no	16	2	2	7870	806,314	stable	7	smooth data2/sg olay/mo vmean		12	yes	1	0.5/0.01	P.B	n
															17	2	2	7870	806,314	stable	7	smooth data2/sg olay/mo vmean		16	Yes	1	0.5/0.01	P.B	n





Near the breakwater (G5), simulation 5 (VISCOSITY BREAKING) provides much more energy.

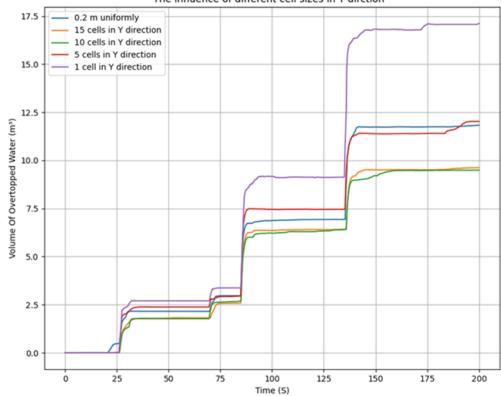




17







The influence of different cell sizes in Y dirction





19

References

Vicinanza, D., Contestabile, P., Quvang Harck Nørgaard, J., & Lykke Andersen, T. (2014). Innovative rubble mound breakwaters for overtopping wave energy conversion. Coastal Engineering, 88, 154–170. https://doi.org/10.1016/j.coastaleng.2014.02.004

Iuppa, C., Contestabile, P., Cavallaro, L., Foti, E., & Vicinanza, D. (2016). Hydraulic performance of an innovative breakwater for overtopping wave energy conversion. Sustainability (Switzerland), 8(12). <u>https://doi.org/10.3390/su8121226</u>

Iuppa, C., Cavallaro, L., Musumeci, R. E., Vicinanza, D., & Foti, E. (2019). Empirical overtopping volume statistics at an OBREC. *Coastal Engineering*, 152. https://doi.org/10.1016/j.coastaleng.2019.103524

Kralli, V. E., Theodossiou, N., & Karambas, T. (2019). Optimal Design of Overtopping Breakwater for Energy Conversion (OBREC) Systems Using the Harmony Search Algorithm. Frontiers in Energy Research, 7. https://doi.org/10.3389/fenrg.2019.00080

Cavallaro, L., Iuppa, C., Castiglione, F., Musumeci, R. E., & Foti, E. (2020). A simple model to assess the performance of an overtopping wave energy converter embedded in a port breakwater. *Journal of Marine Science and Engineering*, 8(11), 1-20. <u>https://doi.org/10.3390/jmse8110858</u>

P. Contestabile, E. Di Lauro, M. Buccino, and D. Vicinanza, "Economic assessment of Overtopping BReakwater for Energy Conversion (OBREC): A case study in Western Australia," Sustainability (Switzerland), vol. 9, no. 1, 2017, doi: 10.3390/su9010051.

P. Contestabile, G. Crispino, E. Di Lauro, V. Ferrante, C. Gisonni, and D. Vicinanza, "Overtopping breakwater for wave Energy Conversion: Review of state of art, recent advancements and what lies ahead," Renew Energy, vol. 147, pp. 705-718, Mar. 2020, doi: 10.1016/j.renene.2019.08.115.

Serinaldi, F., Briganti, R., Kilsby, C. G., & Dodd, N. (2022). Sailing synthetic seas: Stochastic simulation of benchmark sea state time series. Coastal Engineering, 176. https://doi.org/10.1016/j.coastaleng.2022.104164

M. Postacchini et al., "A model chain approach for coastal inundation: Application to the bay of Alghero,"



Laboratory of Hydraulic Engineering Department of Civil Engineering University of Patras



Sediment Transport and Morphodynamics in Marine and Coastal Waters with Engineering Solutions

Large-eddy simulations of turbulent oscillatory flow and sediment transport induced by waves Ioannis Gerasimos Tsipas Supervisor: Athanassios A. Dimas, Professor





This project has received funding from the European Union's Horizon Europe research and innovation programme under the Marie Skłodowska-Curie grant agreement No 101072443.



EUROPEAN COMMISSION

Outline

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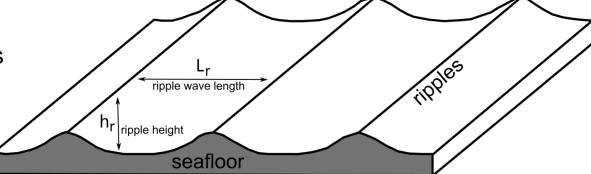
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- Introduction
- Objectives
- Methodology
 - Simulation Set-up
 - Results
- Conclusions/Future Work

- Surface waves induce oscillatory flow at seabed
- Generation of bed forms (ripples, dunes, bars)
- Significant impact on wave propagation and sediment transport by increasing bed roughness and promoting sediment suspension due to vortex shedding

The dynamics of turbulent oscillatory flow and sediment transport over sandy beds is critical for understanding various environmental and engineering processes, such as coastal erosion, sedimentation patterns, and habitat formation.



https://www.vhv.rs/viewpic/hbJbwRw_transparent-water-ripplespng-ripple-of-water-diagram/

Objectives

 Development of large-eddy simulation (LES) software to model turbulent oscillatory flow and mixed-grain sediment transport induced by waves.

• Methodology

Flow Equations (non-dimensional)

Continuity equation: $\frac{\partial u_i}{\partial x_i} = 0$

Navier-Stokes equations
$$\frac{\partial u_i}{\partial t} + \frac{\partial}{\partial x_j} (u_i u_j) = -\frac{\partial p}{\partial x_i} - \frac{\partial \tau_{ij}}{\partial x_j} + \frac{1}{\text{Re}} \frac{\partial^2 u_i}{\partial x_j \partial x_j} + f_i$$
 $Re = \frac{U_0 a_0}{v}$

 u_i is the resolved velocity field according to LES .

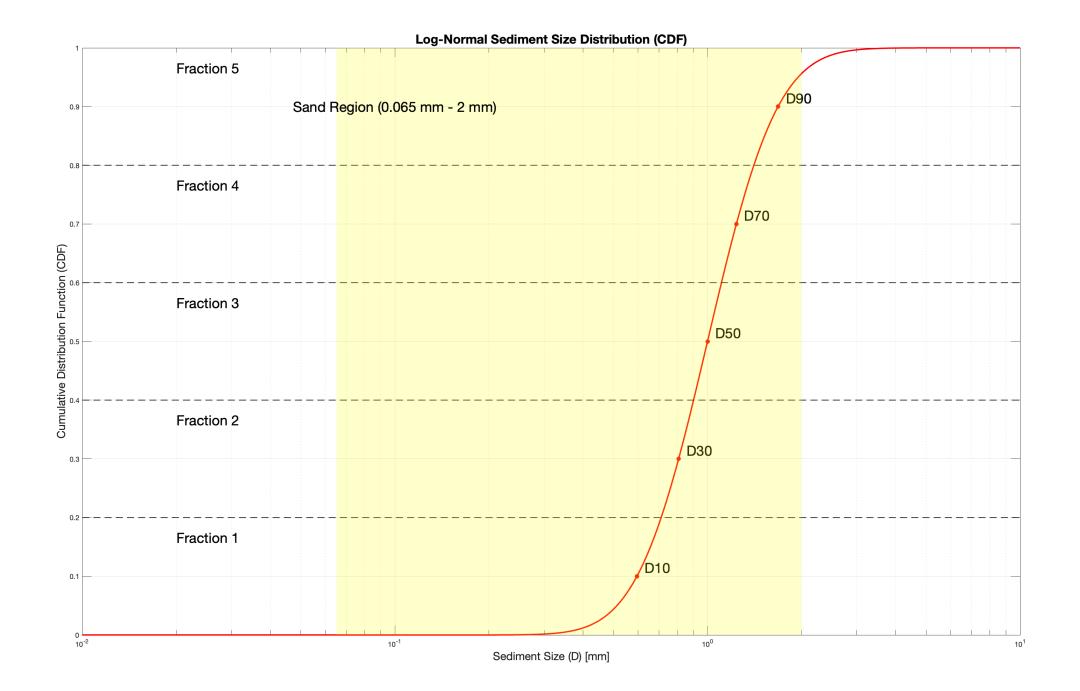
Dynamic pressure $p = P_o + P$ the externally imposed pressure field.

where P_o is $u_o(t) = U_o(\cos(\omega t) + B\cos(2\omega t))$

$$\tau_{ij} = -2D_{wall} v_{sgs} S_{ij} = -2D_{wall} \left(C_s \Delta\right)^2 \left|S\right| S_{ij} \qquad \Delta = \left(\Delta x_1 \Delta x_2 \Delta x_3\right)^{1/3} \qquad \left|S\right| = \left(2S_{ij} S_{ij}\right)^{1/2}$$

Subgrid-scale (SGS) stresses

$$S_{i,j} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$



Sediment Transport Equations

Bed load transport rate (Engelund and Fredsøe, 1976):

$$\frac{q_{b(nd)}}{\sqrt{(S-1)gDg_{(nd)}^3}} = \begin{cases} \operatorname{sgn}(\theta) \frac{5\pi}{3} \left[1 + \left(\frac{\pi}{6} \frac{\mu_d}{|\theta_{(nd)}| - \theta_{c(nd)}} \right)^4 \right]^{-\frac{1}{4}} \left(\sqrt{|\theta_{(nd)}|} - 0.7 \sqrt{|\theta_{c(nd)}|} \right), \left(\theta_{(nd)} > \theta_{c(nd)} \right) \\ 0, \left(\theta_{(nd)} < \theta_{c(nd)} \right), \left(\theta_{(nd)} < \theta_{c(nd)} \right) \end{cases}$$

S

Shields number : $\theta = \frac{\tau_b}{\rho_w(S-1)gDg_{(nd)}^3}$

Critical Shields number: $heta_c(D_{g(nd)},S)$

Grain diameters: $Dg_{(nd)}$

Sediment specific gravity:

Dynamic friction coefficient: $\mu_d \approx 0.5 \mu_s$

Advection-diffusion equation for the suspended sediment concentration:

$$\frac{\partial c_{(nd)}}{\partial t} + u_j \frac{\partial c_{(nd)}}{\partial x_j} - w_{s(nd)} \frac{\partial c_{(nd)}}{\partial x_3} = \frac{1}{\sigma Re} \frac{\partial^2 c_{(nd)}}{\partial x_j \partial x_j} - \frac{\partial \chi_j}{\partial x_i} + f_c$$

where
$$w_{s(nd)}$$
 is the sediment fall velocity (Hallermeier 1981) for each

$$\frac{w_{s(nd)}Dg_{(nd)}}{v} = \begin{cases}
D_{*(nd)}^{3}/18 & D_{*(nd)}^{3} < 39 \\
D_{*(nd)}^{2.1}/6 & for & 39 < D_{*(nd)}^{3} < 10^{4} \\
1.05D_{*(nd)}^{1.5} & 10^{4} < D_{*(nd)}^{3} \le 3 \cdot 10^{6}
\end{cases}$$
 $\chi_{j} = \frac{v_{sgs}}{\sigma_{l}} \frac{\partial c_{(nd)}}{\partial x_{j}}$

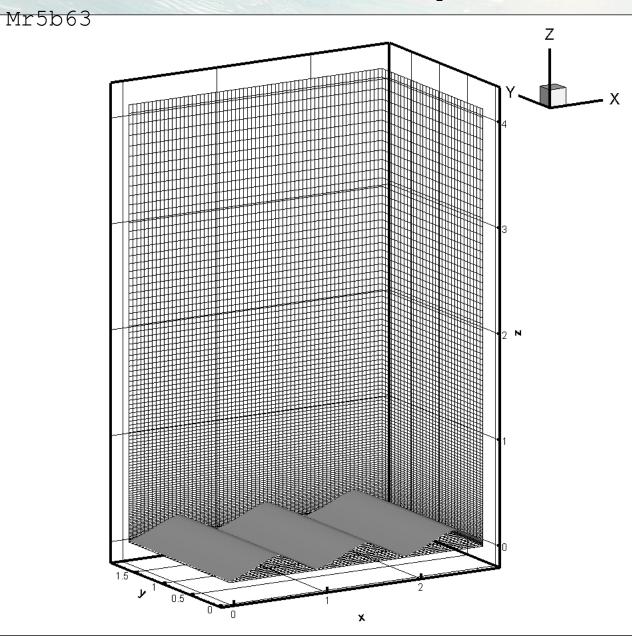
 σ is the Schmidt number, χ_j is the SGS turbulent term (Zedler and Street 2001):

and σ_t is the turbulent Schmidt number.

Suspended load transport rate: $q_{s1,2(nd)} = \int_{x_{3bed}}^{x_{3top}} u_{1,2}c_{(nd)}dx_3$

Simulation Set-up

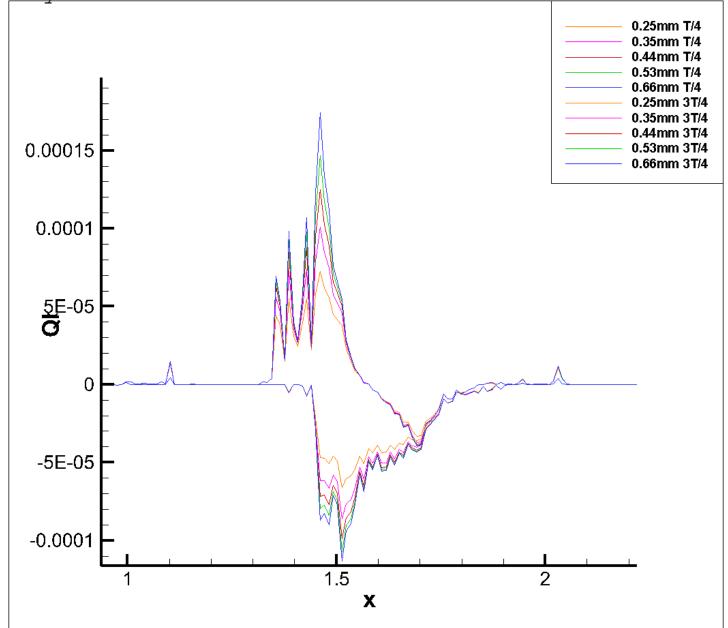
Van Der Werf et al. (2007) , experiment



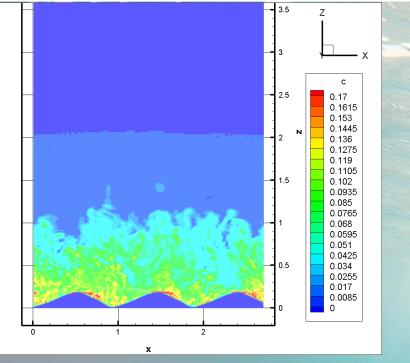
Re = 33000
Lr = 0.94
$h_r = 0.17$ $\Delta x = 0.01055 m$
$\Delta y = 0.01406 m$
$\Delta z = 0.0025 ->$ 0.02719 m
$Grid = 257 \times 129 \times 400$
= 13261200 cells

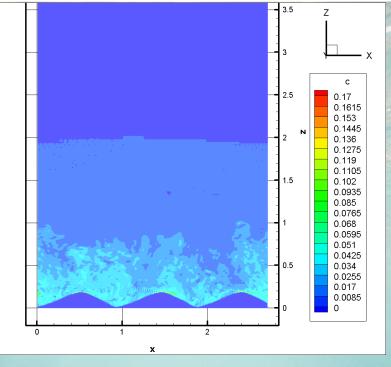
D _g (m)	Ψ	a_0/D_g
0.00025	104.4	2100
0.00035	74.5	1500
0.00044	59.3	1193
0.00053	49.2	990
0.00066	39.5	795

Bed Load, Spanwise & Phase(max positive/negative) averaged for every diameter fraction

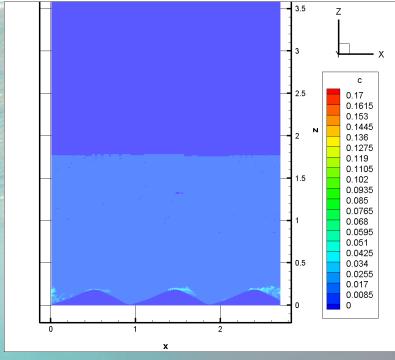


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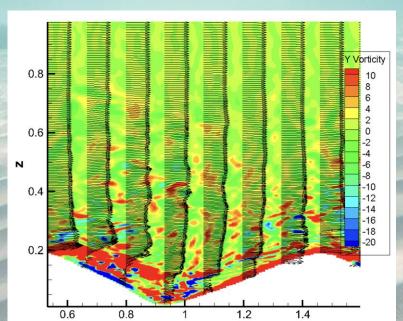


D50 at T/4



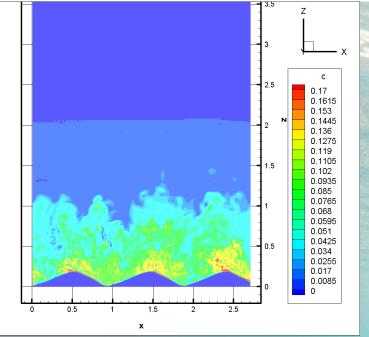
D90 at T/4

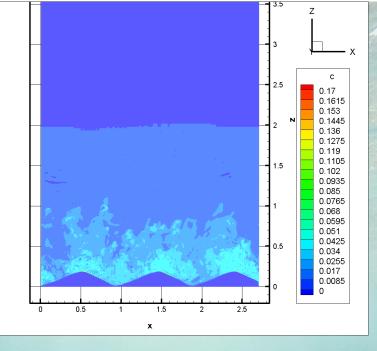
Vorticity Field T/4

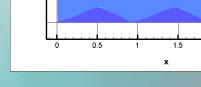


х

D10 at T/4







D10 at T/2

D50 at T/2

D90 at T/2

Ζ

0.17 0.1615 0.153 0.1445 0.136 0.1275 0.119 0.1105 0.102 0.0935 0.085

0.085 0.0765 0.068 0.0595 0.051 0.0425 0.034 0.0255 0.017 0.0085 0

2.5

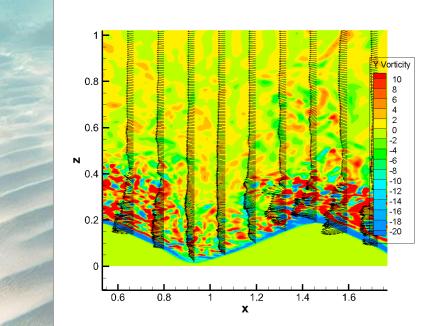
1.5

0.5

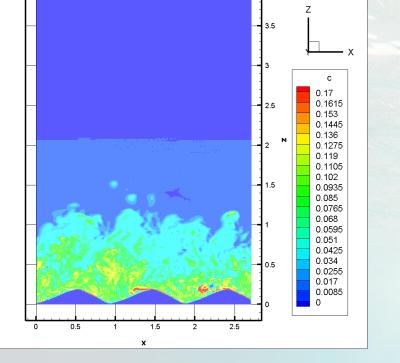
2.5

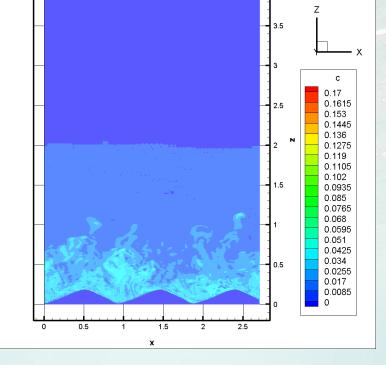
12

2

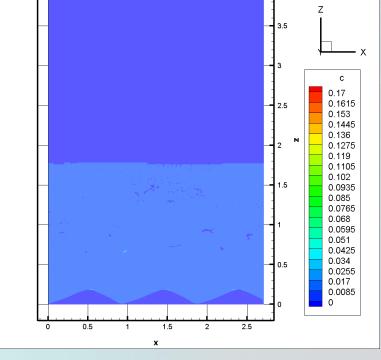


Vorticity Field T/2



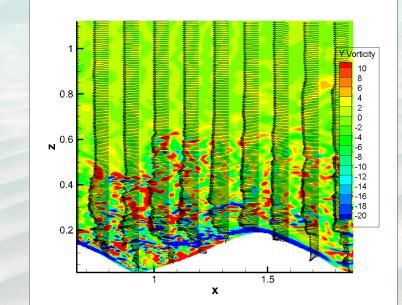


D50 at 3T/4

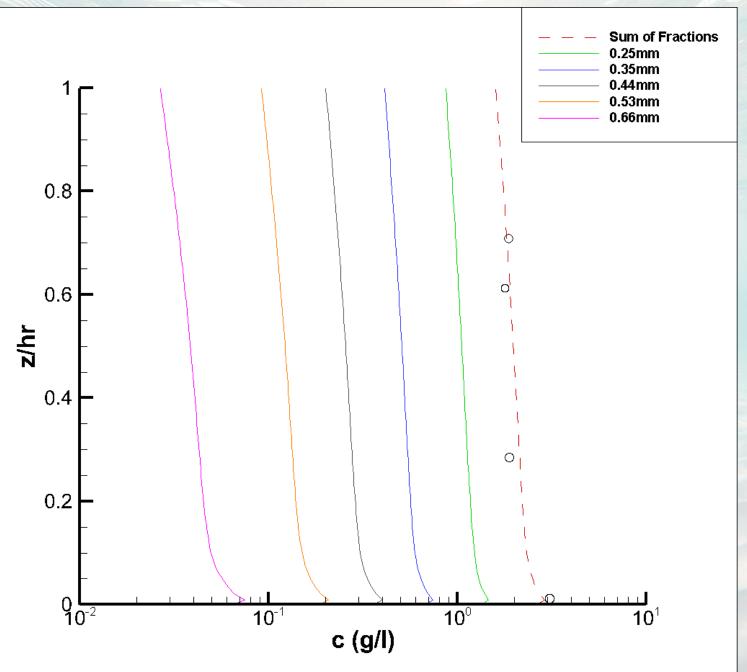


D90 at 3T/4

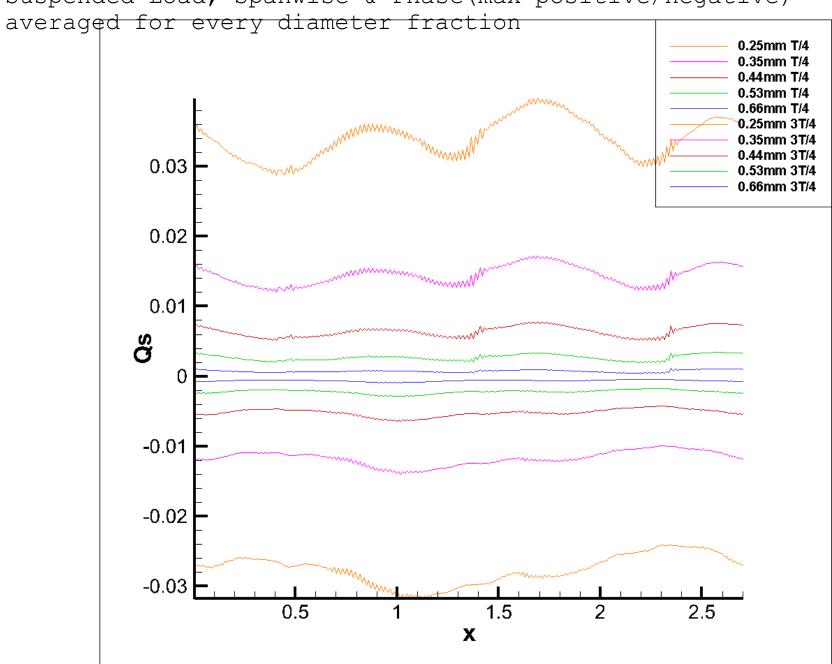
D10 at 3T/4



Vorticity Field 3T/4 Concentration profile for each diameter fraction and the



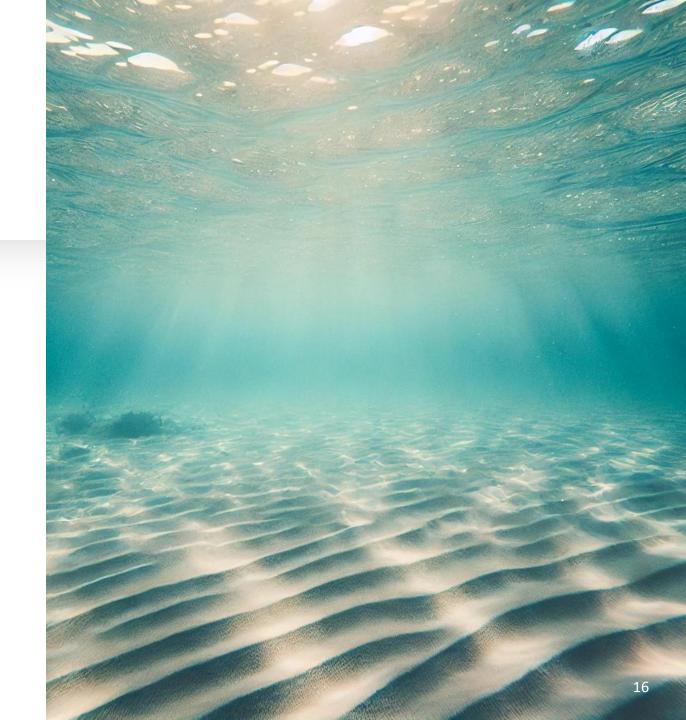
14



Suspended Load, Spanwise & Phase(max positive/negative)

Conclusions

- Bed Load : Diameter Size increasing -> Bed Load Increases
- Suspended Load : Diameter
 Size Decreases -> Suspended
 Load Increases
- Concentration : Diameter Size Decreases -> Concentration Increases









Marie Skłodowska-Curie Actions

SEDIMARE 2023-2027

Sediment Transport and Morphodynamics in Marine and Coastal Waters with Engineering Solutions

2nd Network Training School: Experimental and Practical Modelling of

Sediment Transport and Coastal Morphology

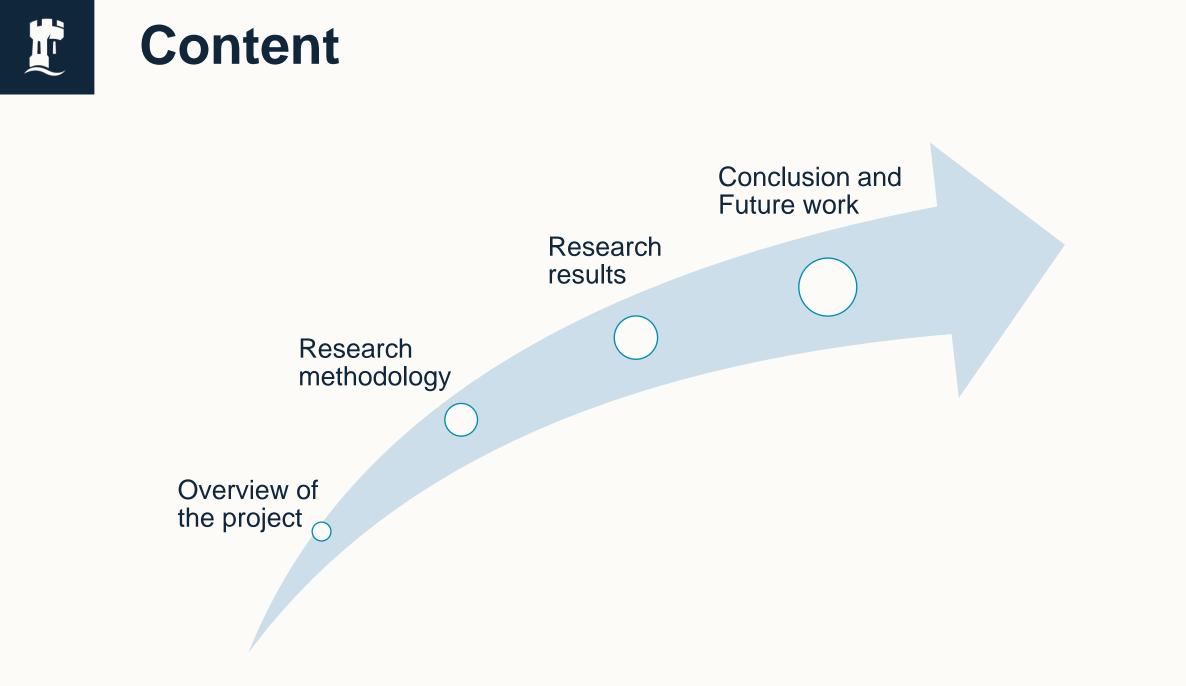
Deltares, The Netherlands

7 – 8 Nov 2024

Morphodynamic swash zone modelling

Doctoral Candidate: Quan NGUYEN

Supervising Scientists: Nicholas DODD and Riccardo BRIGANTI





Overview of the project

Overview of the project

The swash zone

- A dynamic region separating land from the inner surf zone
- Wave run-up/run-down impacts the nearshore morphology and sediment transport.
- Unsteady flow, high turbulence, rapid boundary layer growth and decay, large sediment transport rates, and rapid bed level changes

> Important to coastal management and conservation strategies.

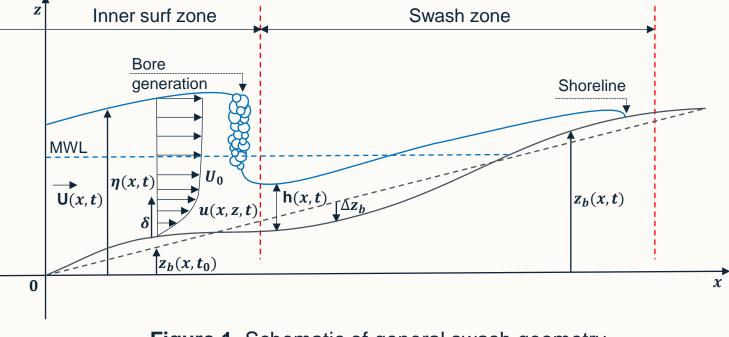
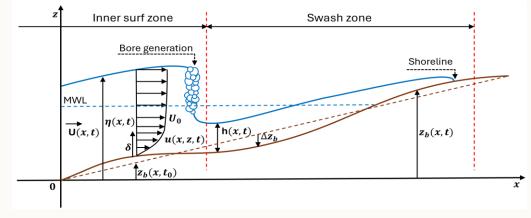


Figure 1. Schematic of general swash geometry.

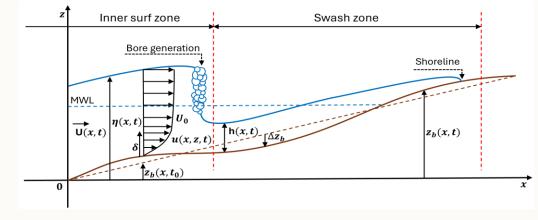
Overview of the project Motivation



- Difficulty in adequately represent the wave boundary layer in the swash zones.
- Difficulty in directly measure the bed shear stress within the bottom boundary layer (BBL)
- An existing BBL sub-model for a **fixed bed** in the swash zone

Further developed on <u>mobile bed</u> beaches, to improve our capacity to accurately described the bed shear stress within the BBL

Overview of the project Objectives



- To develop an <u>improved</u> boundary layer description (sub-model) for a <u>mobile bed</u> that is suitable for incorporation into a Nonlinear Shallow Water Wave Equation (NLSWE) morphodynamic solver.
- > To validate the resulting morphodynamic solver against laboratory and existing field data.

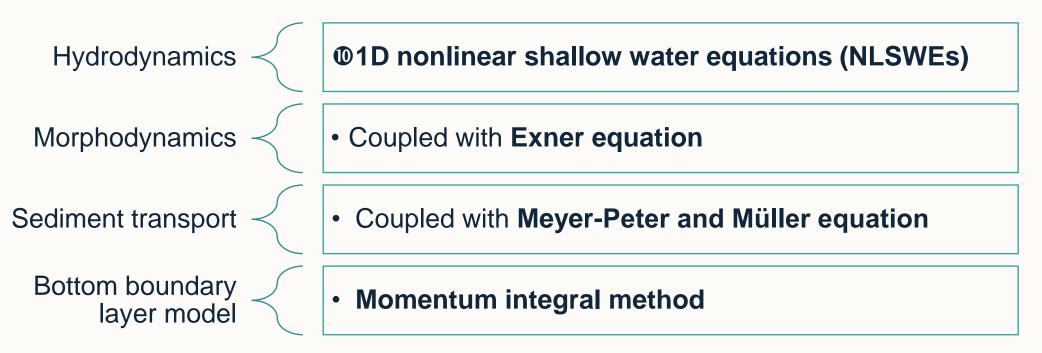


Research methodology

Research methodology

Numerical model

Depth-averaged, phase-resolving, fully-coupled model



Research methodology

Model validation

□ Single dam-break-driven swash event on the fixed bed (Kikkert et al., 2012)

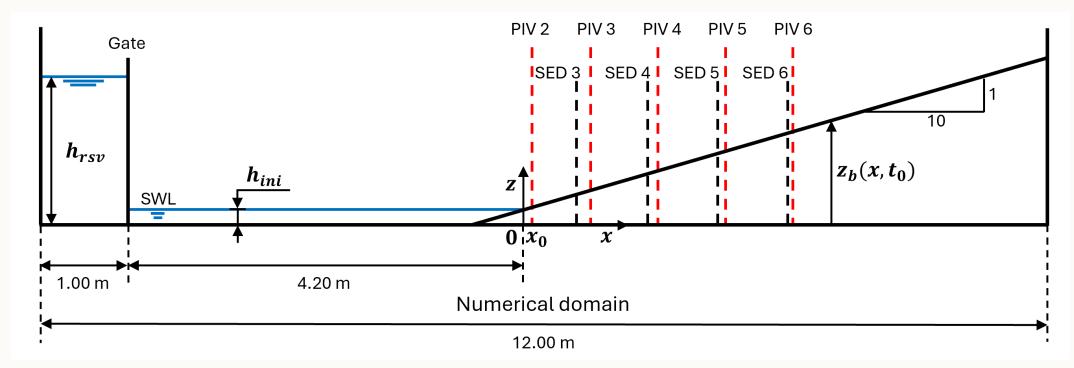


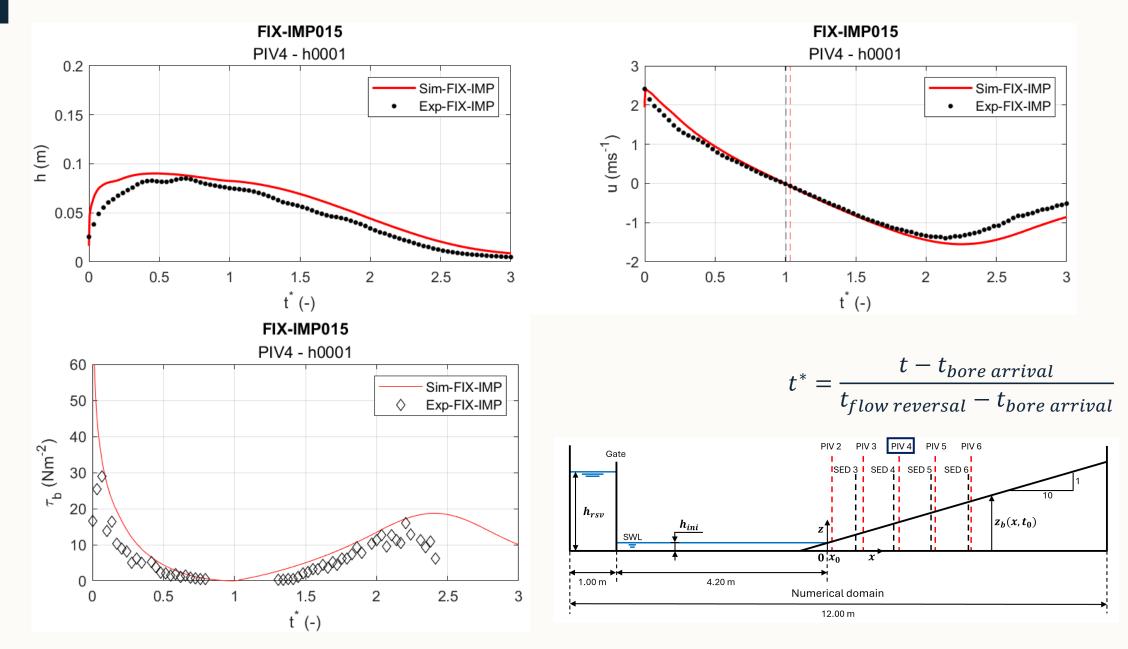
Figure 2. Schematic of the numerical setup based on the Aberdeen Swash facility.



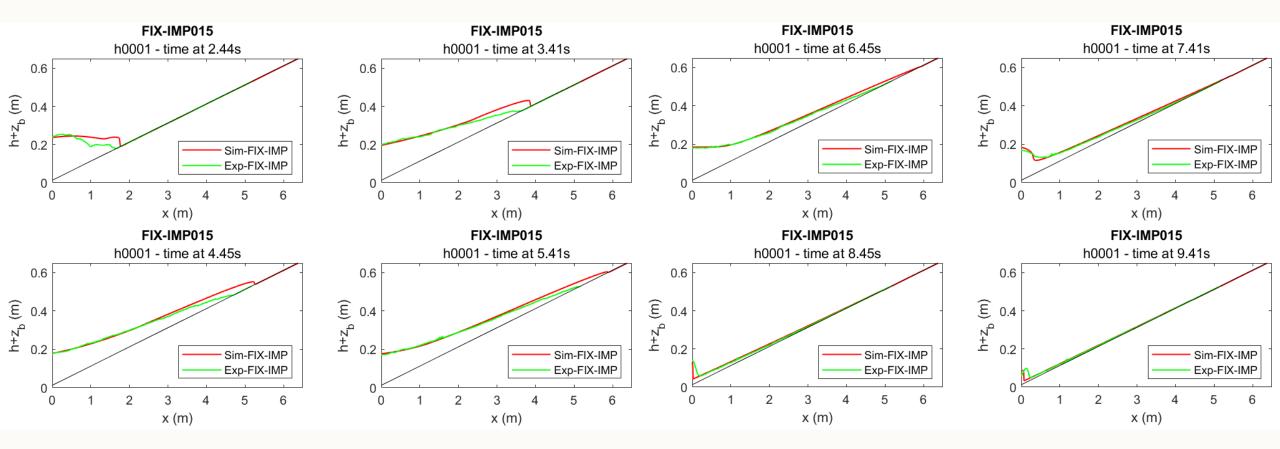
Research results

Model validation: Single swash on impermeable fixed beds

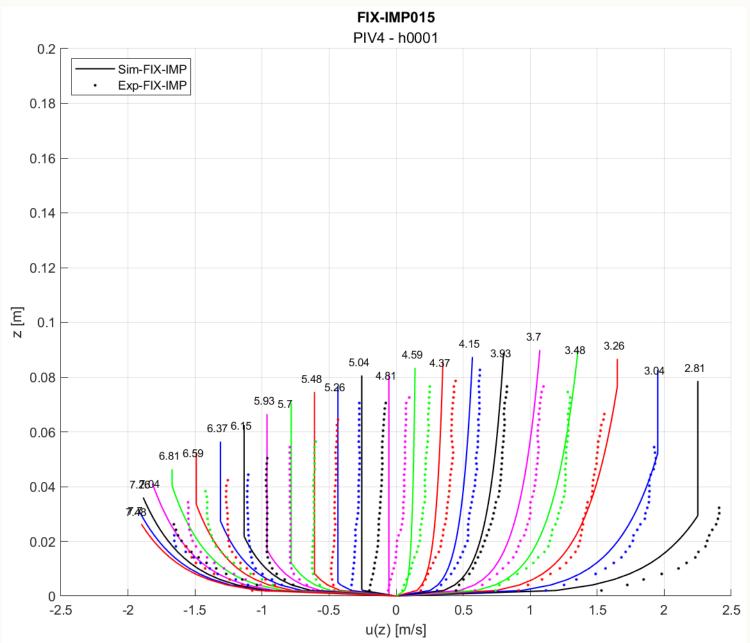
Water depth - Depth-averaged velocity – Bed shear stress



Swash lens



Vertical profile of horizontal velocity

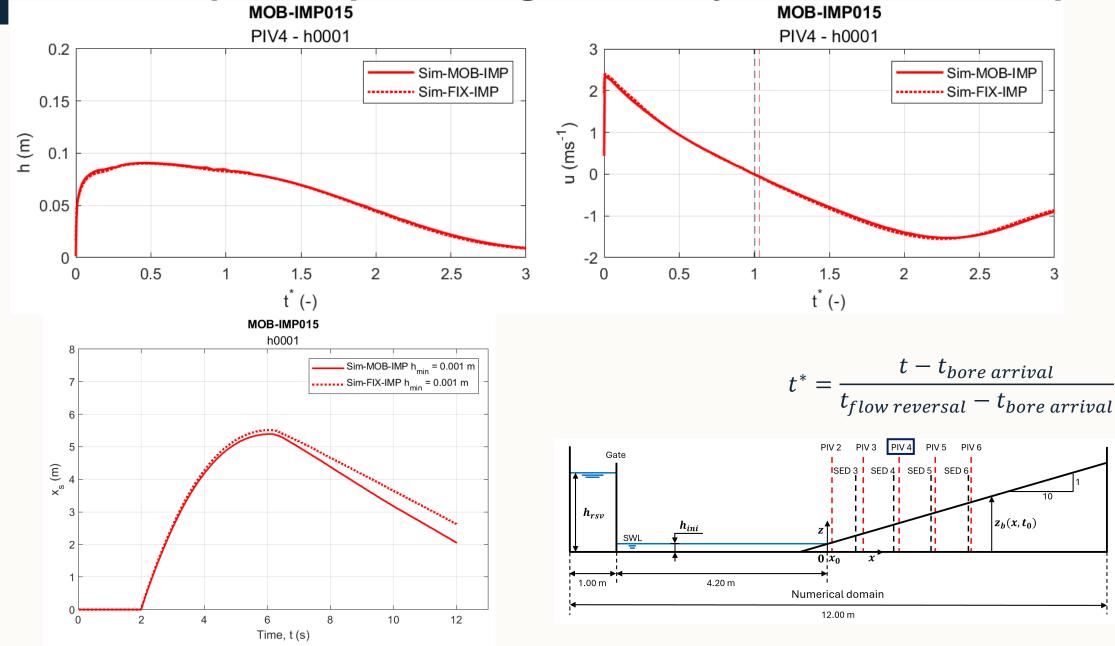




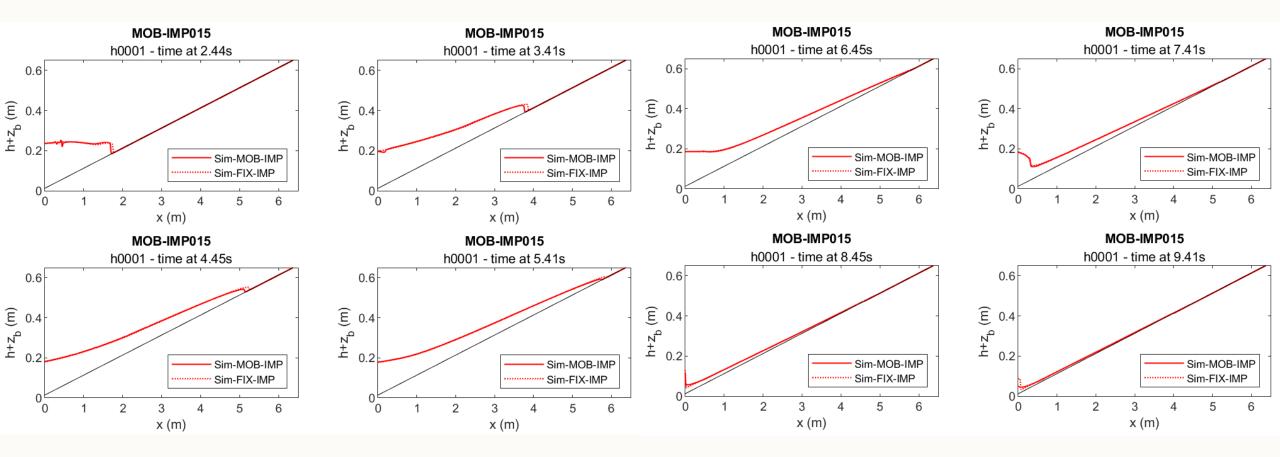
Research results

The impact of mobility on model performance

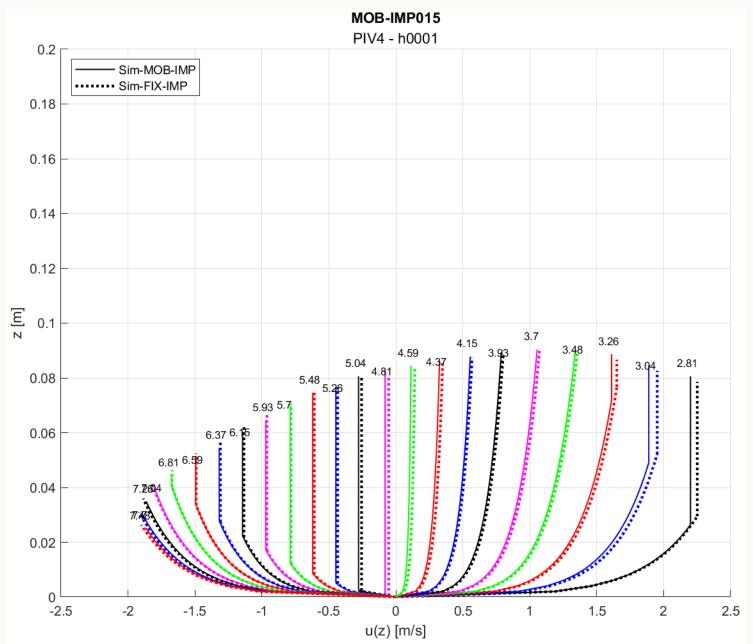
Water depth - Depth-averaged velocity – Maximum run-up



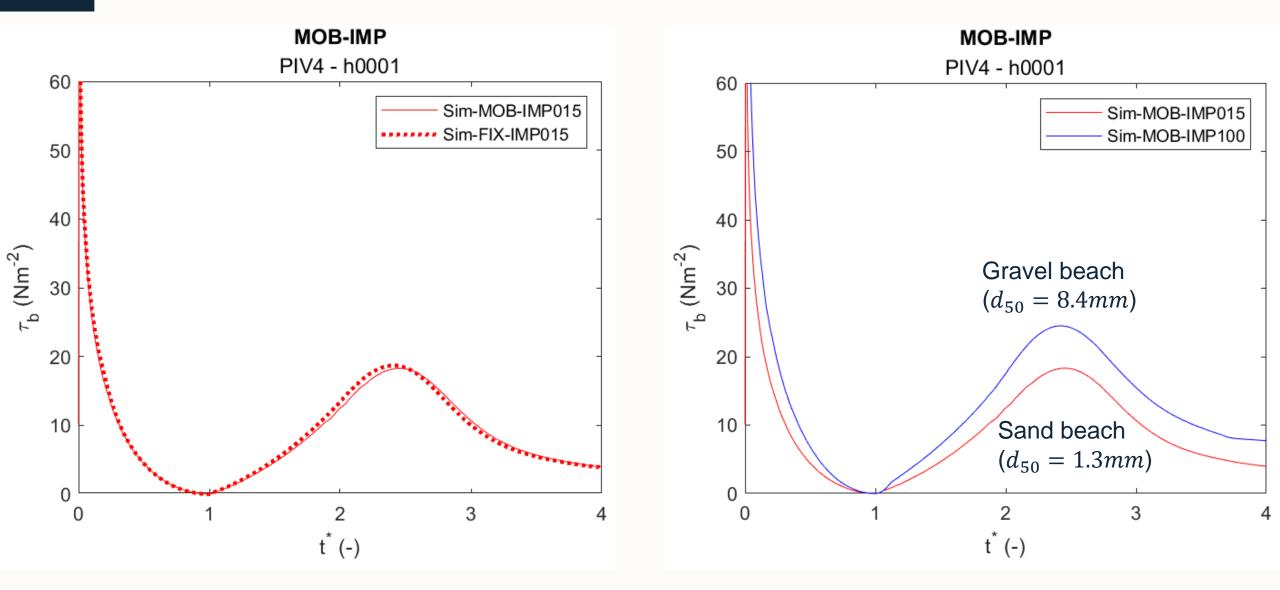
Swash lens



Vertical profile of horizontal velocity



Bed shear stress





Conclusion

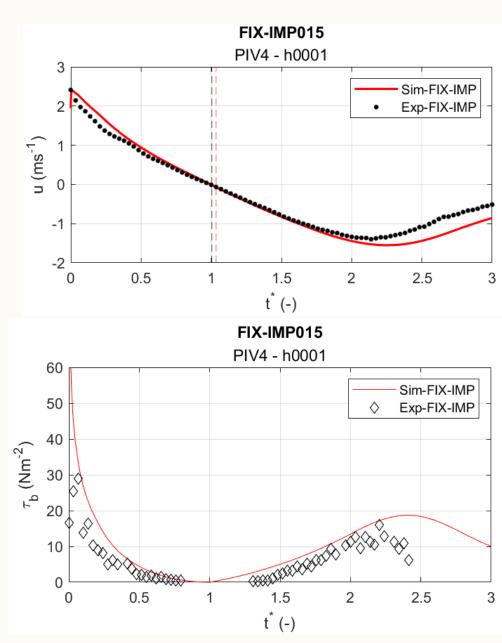
Model advancements

Numerical results of the flow variables for the test with fixed bed are close to the experimental results

The model accurately predicts the depth-averaged horizontal velocity in the run-up phase

The modelled bore arrival time almost coincides with the experimental one

The modelled bed shear stress is well predicted when compared with the log-law-derived shear stress



Model limitations

- Overestimation of boundary layer thickness and underestimation of velocity profile
- Affecting the model's ability to accurately simulate the interactions between the flow and the seabed, which are crucial for understanding the sediment transport processes.
- Uncertainties in bed shear stress modelling
- Affecting the accuracy in modelling of bed shear stress at critical moments such as the arrival of the bore and during flow reversal
- Overestimation of Flow Variables
- > Affecting the reliability of the model in practical applications.



Future works

- Improving the simulation of the vertical profile of horizontal velocity inside the bottom boundary layer, with a focus on the velocity profile at the flow reversal and the swash tip.
- Parameterising the physical parameters and model configuration on the simulation of the dam-break-driven swash event on an impermeable mobile bed.



01

Deliverables

D2.2. Technical report on evaluation of morphodynamic swash zone model with existing boundary layer sub-model.

02

D2.8. Technical report on evaluation of new boundary layer submodel.

D3.11. Technical report on application of morphodynamic swash zone model.









Marie

SEDIMARE 2023-2027 Skłodowska-

Curie Actions Sediment Transport and Morphodynamics in Marine and **Coastal Waters with Engineering Solutions**

2nd Network Training School

University of Twente/Deltares, The Netherlands

Nov 5th – 8th, 2024 Morphodynamic swash zone modelling

Doctoral Candidate (#11): Quan NGUYEN

Supervising Scientists: Nicholas DODD and Riccardo BRIGANTI

Structural stability of the Brinkman–Forchheimer equations for flow in porous media with variable porosity

Evangelos Petridis

UCL ouvain Institute of Mechanics, Materials and Civil Engineering



Funded by the European Union



This project has received funding from the European Union's Horizon Europe research and innovation programme under the Marie Skłodowska-Curie grant agreement No 101072443.



Introduction

Payne and Straughan (1999) established continuous dependence in the Brinkman-Forchheimer equations with constant porosity . When the porosity is space dependent :

- Velocity is not divergence free.
- The term describing viscous shear stresses is not the Laplacian .
- We have normal viscous stresses (bulk viscosity ζ) .
- \bullet The shear viscosity μ enters the expression for the interfacial drag (Darcy coefficient ${\it a})$.
- \bullet We are working in the weighted L^2 space with the porosity $\phi({\bf x})$ being the weight .

$$\|\mathbf{u}\| = \left(\int_{\Omega} \phi |\mathbf{u}|^2 d\mathbf{x}\right)^{1/2}$$

Inequalities

Arithmetic-geometric mean inequality

$$\mathsf{a} b \leq rac{1}{2c} \mathsf{a}^2 + rac{c}{2} b^2$$

Hölder's inequality

$$\int_{\Omega} |\mathbf{f}| |\mathbf{g}| d\mathbf{x} \leq \Bigl(\int_{\Omega} |\mathbf{f}|^p d\mathbf{x} \Bigr)^{1/p} \Bigl(\int_{\Omega} |\mathbf{g}|^q d\mathbf{x} \Bigr)^{1/q}, \ \frac{1}{p} + \frac{1}{q} = 1$$

Sobolev inequality

$$\int_{\Omega} |\mathbf{G}|^4 d\mathbf{x} \leq C \, \left(\int_{\Omega} |\mathbf{G}|^2 d\mathbf{x}\right)^{1/2} \left(\int_{\Omega} \boldsymbol{\nabla} \mathbf{G} : \boldsymbol{\nabla} \mathbf{G} d\mathbf{x}\right)^{3/2}$$

Poincare's inequality

$$\lambda_1 \int_{\Omega} |\mathbf{w}|^2 d\mathbf{x} \leq \int_{\Omega} \mathbf{
abla} \mathbf{w} : \mathbf{
abla} \mathbf{w} d\mathbf{x}$$

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Model

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The Brinkman–Forchheimer equations for flow in porous media with variable porosity are

$$\begin{split} \phi \frac{\partial \mathbf{u}}{\partial t} + \phi \nabla p \\ &= \nabla \cdot (\phi \zeta (\nabla \cdot \mathbf{u}) I) + \nabla \cdot (\phi \mu \mathbf{V}^d) - a^*(\phi) \mathbf{u} - b^*(\phi) |\mathbf{u}| \mathbf{u} + \phi \mathbf{f} \quad , \quad \zeta, \mu > 0 \\ &\nabla \cdot (\phi \mathbf{u}) = 0 \\ &\mathbf{V}^d(\mathbf{u}) = \frac{1}{2} (\nabla \mathbf{u} + (\nabla \mathbf{u})^T) - \frac{1}{3} \nabla \cdot \mathbf{u} \ I \\ &b^*(\phi) = b(1 - \phi) + d(1 - \phi)^2 \quad , \ a^*(\phi) = a\mu(1 - \phi) \ , \ a, b, d > 0 \end{split}$$

$$0 < \phi_{\min} \leq \phi \leq \phi_{\max} < 1$$

where **u** is the average fluid velocity in the porous medium, *a* is the Darcy coefficient, *b* is the Forchheimer coefficient, ζ is the bulk viscosity , μ is the shear viscosity, *p* is the pressure , **f** is the gravity and ϕ is the variable porosity.

Forchheimer coefficient b

To study continuous dependence on b, we let u and v solve the following boundary initial-value problems for different Forchheimer coefficients b_1 and b_2 :

$$\phi \frac{\partial \mathbf{u}}{\partial t} + \phi \nabla p = \nabla \cdot (\phi \zeta (\nabla \cdot \mathbf{u}) I) + \nabla \cdot (\phi \mu \mathbf{V}^d) - a^*(\phi) \mathbf{u} - b_1^*(\phi) |\mathbf{u}| \mathbf{u} + \phi \mathbf{f} ,$$

$$\phi \frac{\partial \mathbf{v}}{\partial t} + \phi \nabla q = \nabla \cdot (\phi \zeta (\nabla \cdot \mathbf{v})I) + \nabla \cdot (\phi \mu \mathbf{V}^d) - a^*(\phi)\mathbf{v} - b_2^*(\phi)|\mathbf{v}|\mathbf{v} + \phi \mathbf{f} ,$$

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abla} \cdot (\phi {oldsymbol u}) = {oldsymbol
abla} \cdot (\phi {oldsymbol v}) = 0 \ , \ \Omega imes \{t > 0\} \ ,$$

$$\mathbf{u} = \mathbf{v} = \mathbf{0} \ , \ \partial \Omega \times \{t > 0\} \ ,$$

$$\boldsymbol{\mathsf{u}}(\boldsymbol{\mathsf{x}},0)=\boldsymbol{\mathsf{v}}(\boldsymbol{\mathsf{x}},0)=\boldsymbol{\mathsf{g}}(\boldsymbol{\mathsf{x}})\;,\;\boldsymbol{\mathsf{x}}\in\Omega$$

 Ω in these problems is a bounded domain in \mathbb{R}^3 whose boundary is $\partial\Omega$ and g is the given initial data.

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Forchheimer coefficient b

We define the difference variables ${\bf w}$, p and b by

$$\mathbf{w} = \mathbf{u} - \mathbf{v}$$
, $\pi = p - q$, $b^* = b_1^* - b_2^* = b(1 - \phi)$, $b = b_1 - b_2$

and then \mathbf{w} satisfies the boundary initial-value problem

$$\begin{split} \phi \frac{\partial \mathbf{w}}{\partial t} &- \nabla \cdot (\phi \zeta (\nabla \cdot \mathbf{w}) I) - \nabla \cdot (\phi \mu \mathbf{V}^d) \\ &= -a^*(\phi) \mathbf{w} - (b_1^*(\phi) |\mathbf{u}| \mathbf{u} - b_2^*(\phi) |\mathbf{v}| \mathbf{v}) - \phi \nabla \pi , \\ \nabla \cdot (\phi \mathbf{w}) &= 0 , \ \Omega \times \{t > 0\} , \\ \mathbf{w} &= \mathbf{0} , \ \partial \Omega \times \{t > 0\} , \end{split}$$

$$\boldsymbol{\mathsf{w}}(\boldsymbol{\mathsf{x}},\boldsymbol{0}) = \boldsymbol{\mathsf{0}} \ , \ \boldsymbol{\mathsf{x}} \in \boldsymbol{\Omega}$$

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Forchheimer coefficient b

Using the relations

$$\mathbf{V}^d$$
 : $\mathbf{V}^d = \mathbf{V}^d$: $\mathbf{\nabla}\mathbf{w}$

$$\|\mathbf{V}^d\|^2 \ge rac{\phi_{\min}}{2\phi_{\max}}\|\mathbf{\nabla}\mathbf{w}\|^2$$

and calculations we can conclude

$$\|\mathbf{w}\|^2 \leq rac{c}{\lambda_1^{1/2}\mu^2} (rac{1}{\phi_{\min}} - 1)^2 rac{(\phi_{\max})^{7/2}}{(\phi_{\min})^{9/2}} \|\mathbf{g}\|^4 e^{-2a\mu(rac{1}{\phi_{\max}} - 1)t} b^2$$

This specific estimate establishes continuous dependence on the Forchheimer coefficient b in the weighted L^2 norm.

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Results

Also through calculations we can establish continuous dependence on

• the Darcy coefficient

$$\|\mathbf{w}\|^2 \leq (\frac{1}{\phi_{\min}} - 1)^2 \frac{1}{a_1(\frac{1}{\phi_{\max}} - 1)^2(a_1 - 2a_2)} \|\mathbf{g}\|^2 \left(e^{-2a_2\mu(\frac{1}{\phi_{\max}} - 1)t} - e^{-a_1\mu(\frac{1}{\phi_{\max}} - 1)t}\right)a^2$$

• the shear viscosity

$$\begin{split} \|\mathbf{w}\|^{2} &\leq \Big[(\frac{1}{\phi_{\min}} - 1)^{2} \frac{\|\mathbf{g}\|^{2}}{\mu_{1}(\frac{1}{\phi_{\max}} - 1)^{2}(\mu_{1} - 2\mu_{2})} \left(e^{-2a\mu_{2}(\frac{1}{\phi_{\max}} - 1)t} - e^{-a\mu_{1}(\frac{1}{\phi_{\max}} - 1)t} \right) \\ &+ \frac{\phi_{\max}}{2\mu_{1}\mu_{2}\phi_{\min}} \|\mathbf{g}\|^{2} e^{-a\mu_{2}(\frac{1}{\phi_{\max}} - 1)t} \Big] \mu^{2} \end{split}$$

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Now, let \mathbf{u} be a solution to the boundary initial-value problem for the Brinkman-Forchheimer equations with variable porosity.

$$\begin{split} &\frac{1}{2}\frac{d}{dt}\int_{\Omega}\phi|\mathbf{u}|^{2}d\mathbf{x}+\int_{\Omega}a^{\star}(\phi)|\mathbf{u}|^{2}d\mathbf{x}+\int_{\Omega}b^{\star}(\phi)|\mathbf{u}|^{3}d\mathbf{x}+\int_{\Omega}\phi\mu\mathbf{V}^{d}:\mathbf{\nabla}\mathbf{u}d\mathbf{x}\\ &=-\int_{\Omega}\phi\zeta(\mathbf{\nabla}\cdot\mathbf{u})^{2}d\mathbf{x} \end{split}$$

Through basic calculations and Poincare's inequality we can easily establish an upper bound

$$\|\mathbf{u}\|^{2} \leq \|\mathbf{g}\|^{2} e^{-(2a\mu(\frac{1}{\phi_{\max}}-1)+\frac{\lambda_{1}\mu\phi_{\min}^{2}}{\phi_{\max}^{2}})t}$$

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Energy bounds

Next, we derive a lower bound for $\|\boldsymbol{u}\|$. To do this, let the kinetic energy be

 $\Phi = \|\boldsymbol{u}\|^2$

and then by calculation,

$$\begin{split} \frac{d\Phi}{dt} &= \\ &-2\int_{\Omega}\phi\zeta(\boldsymbol{\nabla}\cdot\boldsymbol{\mathbf{u}})^{2}d\boldsymbol{\mathbf{x}} - 2\int_{\Omega}\phi\mu\boldsymbol{\mathbf{V}}^{d}:\boldsymbol{\mathbf{V}}^{d}d\boldsymbol{\mathbf{x}} - 2\int_{\Omega}a^{*}(\phi)|\boldsymbol{\mathbf{u}}|^{2}d\boldsymbol{\mathbf{x}} - 2\int_{\Omega}b^{*}(\phi)|\boldsymbol{\mathbf{u}}|^{3}d\boldsymbol{\mathbf{x}} \\ &\leq \chi \end{split}$$

where we have set

$$\chi(t) = -2\int_{\Omega}\phi\mu\mathbf{V}^{d}:\mathbf{V}^{d}d\mathbf{x}-2\int_{\Omega}\phi\zeta(\mathbf{\nabla\cdot u})^{2}d\mathbf{x}-2\int_{\Omega}a^{*}(\phi)|\mathbf{u}|^{2}d\mathbf{x}-\frac{4}{3}\int_{\Omega}b^{*}(\phi)|\mathbf{u}|^{3}d\mathbf{x}$$

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By further computation, we find that

$$\frac{d\chi}{dt} = 4 \int_{\Omega} \phi |\frac{\partial \mathbf{u}}{\partial t}|^2 d\mathbf{x} = 4 \|\frac{\partial \mathbf{u}}{\partial t}\|^2$$

and so with the aid of the Cauchy-Schwarz inequality and some basic calculations we can see that

$$\Phi(t) \geq \Phi(0) e^{\frac{3\chi(0)}{2\|\mathbf{g}\|^2}t}$$

so we have established a lower and an upper bound for $\|\boldsymbol{u}\|$

$$\|\mathbf{g}\|^2 e^{\frac{3\chi(0)}{2\|\mathbf{g}\|^2}t} \leq \|\mathbf{u}\|^2 \leq \|\mathbf{g}\|^2 e^{-(2a\mu(\frac{1}{\phi_{\max}}-1)+\frac{\lambda_1\mu\phi_{\min}^2}{\phi_{\max}^2})t}$$

This, of course, shows that u cannot vanish identically in a finite time

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Energy bounds

Now that we know $\|\mathbf{u}\|$ is strictly positive, we can improve the upper bound. If we use Holder's inequality and with the aid of Poincare's inequality, we may derive

$$\frac{d}{dt} \|\mathbf{u}\|^2 + (2a\mu(\frac{1}{\phi_{\max}} - 1) + \frac{\lambda_1 \mu \phi_{\min}^2}{\phi_{\max}^2}) \|\mathbf{u}\|^2 + \frac{2b}{m^{1/2}(\phi_{\max})^{1/2}} (\frac{1}{\phi_{\max}} - 1) \|\mathbf{u}\|^3 \le 0$$

where $m = m(\Omega)$ is the measure of Ω . This inequality may be rearranged and

integrated into

$$\leq \frac{\|\mathbf{g}\|(2a\phi_{\max}(1-\phi_{\max})+\lambda_{1}\phi_{\min}^{2})A(t)}{2a\phi_{\max}(1-\phi_{\max})+\lambda_{1}\phi_{\min}^{2}+2\|\mathbf{g}\|b\ m^{-1/2}\mu^{-1}\phi_{\max}^{1/2}(1-\phi_{\max})(1-A(t))}$$

where

||...||

$$A(t)=e^{-(a\mu(rac{1}{\phi_{\max}}-1)+rac{\lambda_1\mu\phi_{\min}^2}{2\phi_{\max}^2})t}$$

So finally we established a lower and an upper bound for $\|\boldsymbol{u}\|$

 $\|\mathbf{g}\|e^{\frac{3\chi(0)}{4\|\mathbf{g}\|^2}t}$

 $\leq \| \mathbf{u} \|$

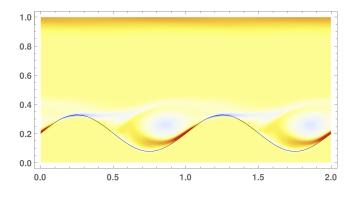
$$\leq \frac{\|\mathbf{g}\|(2a\phi_{\max}(1-\phi_{\max})+\lambda_1\phi_{\min}^2)A(t)}{2a\phi_{\max}(1-\phi_{\max})+\lambda_1\phi_{\min}^2+2\|\mathbf{g}\|b\ m^{-1/2}\mu^{-1}\phi_{\max}^{1/2}(1-\phi_{\max})(1-A(t))}$$

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We will work on a problem include the flow of currents over deformable dunes in a channel.

Detailed numerical simulations of shear-driven (Couette) of water-sand mixtures. (Emphasis on sediment mobilization and resuspension)



L. E. Payne, B. Straughan *Convergence and Continuous Dependence for the Brinkman–Forchheimer Equations*. Stud. Appl. Math. **102** (1999), 419–439.